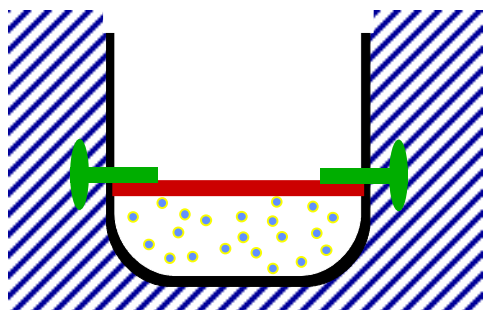
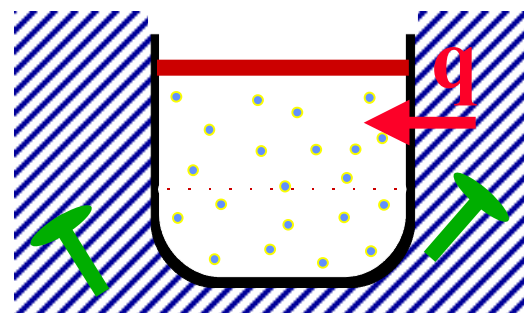


Isothermal expansion of a perfect gas



n, R, T constant
 $\Rightarrow p_1 V_1 = p_2 V_2$
 (for a perfect gas)



constant temperature (isothermal)

$n = 1 \text{ mol}; p_1 = 15 \text{ atm}; V_1 = 1 \text{ L}$

$\Rightarrow p_2 = 1 \text{ atm}; V_2 = 15 \text{ L}$

If the pins are removed what happens?

What is the change in internal energy?

$$\Delta U = U_2 - U_1 \propto \Delta T \quad \Rightarrow \Delta U = 0$$

What is the change in enthalpy?

$$\Delta H = H_2 - H_1 = (U_2 + \overset{0}{p_2 V_2}) - (U_1 + \overset{0}{p_1 V_1})$$

$$\Rightarrow \Delta H = 0$$

Note: $\Delta U = 0$ does not mean that
 q or w is 0, only that $q = -w$;
 in fact, $w = -p_2 \Delta V$

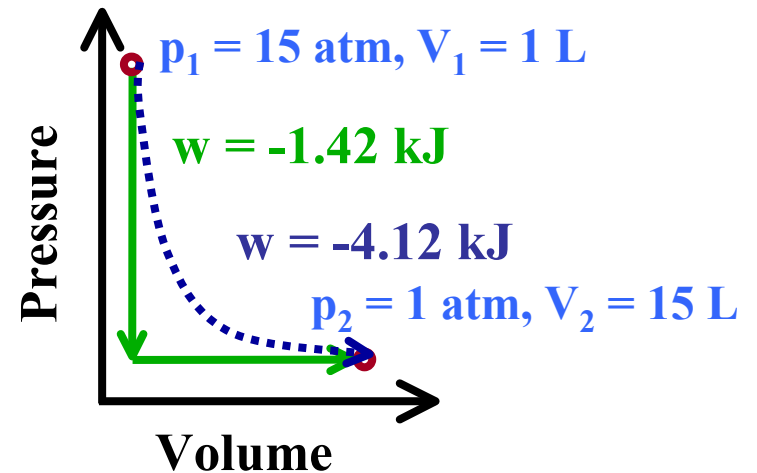
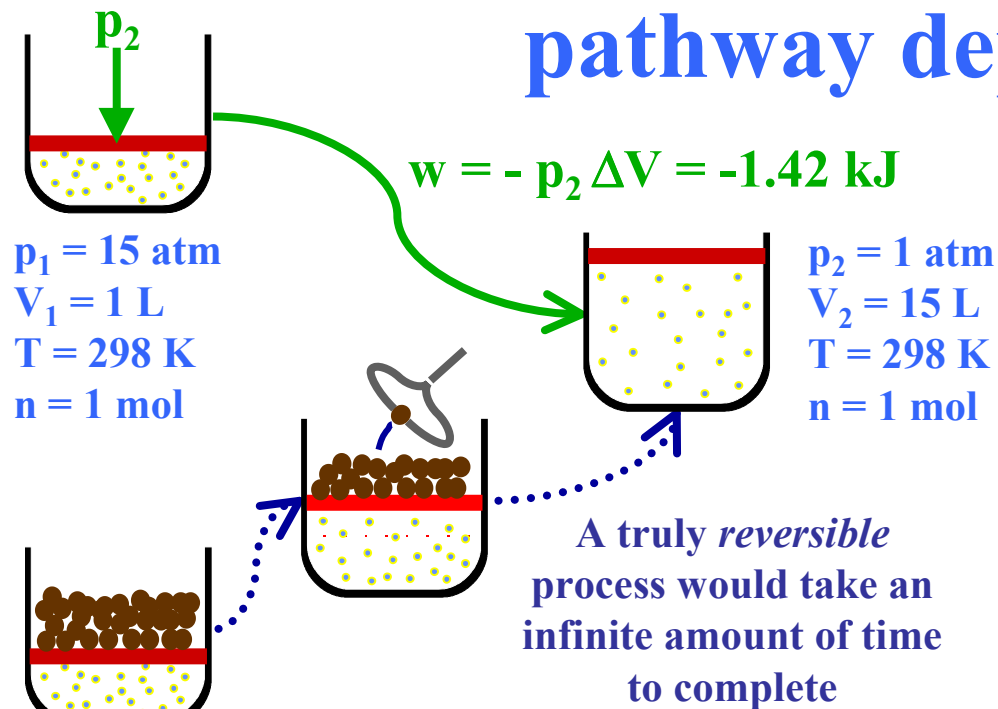
e.g., $p_2 = 1 \text{ atm} = 101,325 \text{ Pa}$

$w = -101,325 \text{ Pa (14 L) (10}^{-3} \text{ m}^3/\text{L)}$

$\Rightarrow w = -1.42 \text{ kJ} \quad q = 1.42 \text{ kJ}$



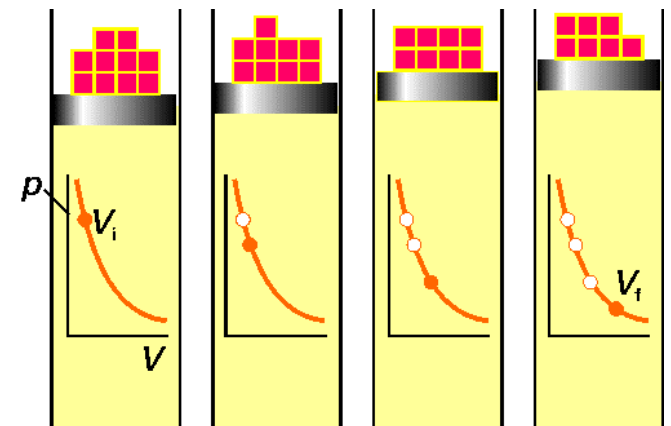
Work in an isothermal expansion is pathway dependent



A reversible process is in equilibrium at all times

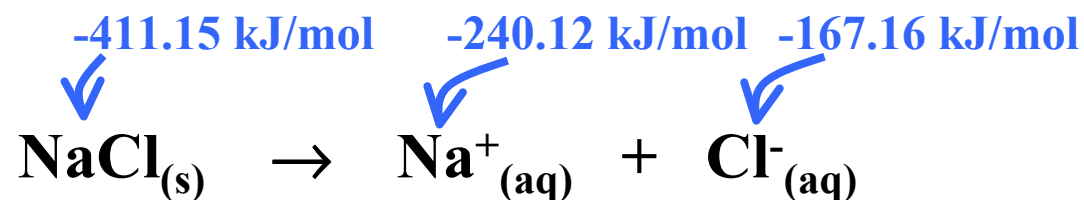
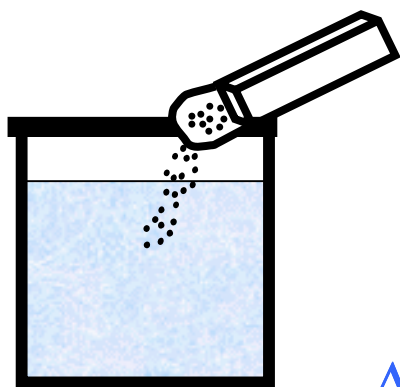
$$w = - \int_{V_1}^{V_2} p dV \stackrel{\text{perfect gas}}{\downarrow} = -nRT \int_{V_1}^{V_2} \frac{dV}{V}$$

$$w = -nRT \ln \left(\frac{V_2}{V_1} \right) = -4.12 \text{ kJ}$$



A reversible process represents the maximum possible pV work

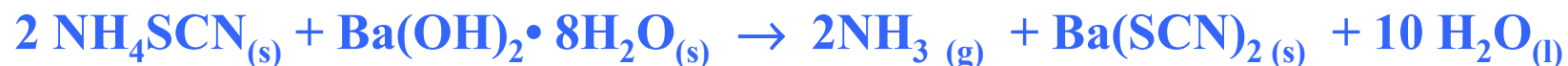
But there's more: Disorder



Hess: $\Delta_r H^\ominus = \sum v \Delta_f H^\ominus (\text{Prod.}) - \sum v \Delta_f H^\ominus (\text{React.})$

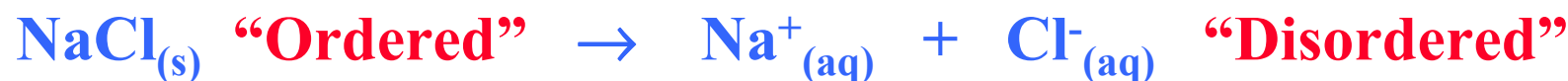
$$\Delta_r H = [-167.16 - 240.12 - (-411.15)] \text{ kJ/mol} = +3.87 \text{ kJ/mol}$$

This is an **endothermic** reaction - but clearly spontaneous, as was the endothermic reaction demonstrated in class:



The reverse reactions are not spontaneous

Qualitatively: “Nature Prefers Disorder”



Measuring disorder in thermodynamics: Statistical thermodynamics

Consider a system with 3 balls



--- and 3 bowls:



A possible state of the system is
1 ball in each bowl

1	2	3
A	B	C
A	C	B
B	A	C
B	C	A
C	A	B
C	B	A

6 Possibilities

$$\# = \frac{N!}{N_1! N_2! N_3!}$$

$$\# = \frac{3!}{1! 1! 1!}$$

Disordered State

Another possible state is 2 balls in
bowl 1, 1 ball in bowl 2

1	2	3
AB	C	
AC	B	
BC	A	

3 Possibilities

$$\# = \frac{3!}{2! 1! 0!}$$

More ordered
State

Or 3 balls in bowl 1

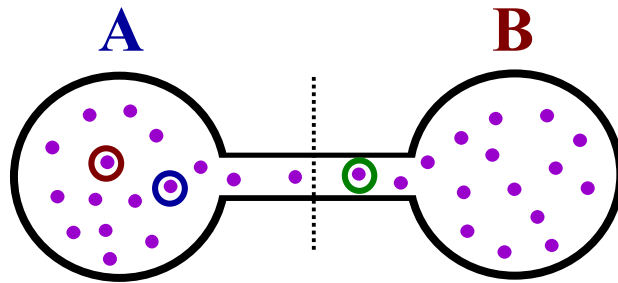
1	2	3
ABC		

1 Possibility

$$\# = \frac{3!}{3! 0! 0!}$$

Ordered State

What the h... does this have to do with a gas?



Gas molecules in two gas bulbs of equal size

The probability of

any one molecule being in A is $P = \frac{1}{2}$

any two molecules in A is $P = \frac{1}{2} \cdot \frac{1}{2} = \frac{1}{4}$

any three molecules in A is $P = \frac{1}{2} \cdot \frac{1}{2} \cdot \frac{1}{2} = \frac{1}{8}$



Probability of N molecules in A

$$P = \left(\frac{1}{2}\right)^N$$

Coin Flips



$$\left(\frac{1}{2}\right)^1$$



$$\left(\frac{1}{2}\right)^2$$



$$\left(\frac{1}{2}\right)^3$$



N consecutive heads

$$P = \left(\frac{1}{2}\right)^N$$

Nils W

Measuring disorder in thermodynamics: Statistical thermodynamics

Consider a system with 3 balls



--- and 3 bowls:



A possible state of the system is
1 ball in each bowl

1	2	3
A	B	C
A	C	B
B	A	C
B	C	A
C	A	B
C	B	A

6 Possibilities

$$W = 6$$

$$P = \frac{6}{18} = \frac{1}{3}$$

Another possible state is 2 balls in
bowl 1, 1 ball in bowl 2

1	2	3
AB	C	
AC	B	
BC	A	

3 Possibilities

$$W = 3$$

$$P = \frac{3}{18} = \frac{1}{6}$$

Or 3 balls in bowl 1

1	2	3
ABC		

1 Possibility

$$W = 1$$

$$P = \frac{1}{18}$$



The statistical definition of disorder = entropy S

**Boltzmann
Equation**

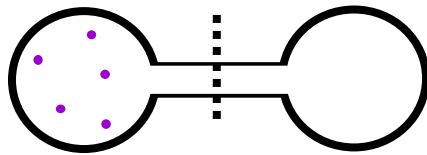


$$S = k \ln W$$

k = Boltzmann Constant
 $1.3805 \times 10^{-23} \text{ J K}^{-1}$

W = number of microstates or
degeneracy of a macrostate

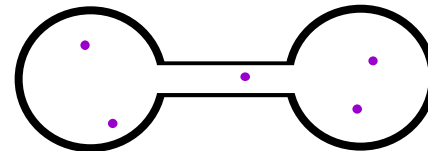
⇒ S has units of J K^{-1}



$$W = 5$$

$$S = 1.61 k$$

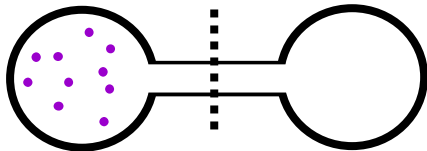
S accounts for the **disorder**



$$W = 10$$

$$S = 2.30 k$$

S is an **extensive** function



$$W = 10$$

$$S = 2.30 k$$

S is a **state function**, i.e., the value of S depends on the state, not on how the state was reached