

Numerical Noise: The Pros and Cons of Filters, Diffusion and Damping Mechanisms

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Diffusion, Filters and Fixers

- Equations of motion: diabatic effects
- Diffusion
 - Explicit horizontal diffusion (neglecting vertical diffusion)
 - Implicit numerical diffusion
 - Divergence damping
 - Decentering mechanism
- Spatial filters:
 - Polar filters / Fourier filters
 - Digital filters: e.g. Shapiro filters
- Time filters: Asselin-filter
- *a posteriori* Fixers:
 - Mass
 - Energy

The 3D Primitive Equations: diabatic effects

Horizontal momentum equation with $\vec{v}_h = (u, v)$

$$\frac{\partial \vec{v}_h}{\partial t} + \left(\vec{v}_h \vec{\nabla}_z \right) \vec{v}_h + w \frac{\partial \vec{v}_h}{\partial z} + f \vec{k} \times \vec{v}_h = -\frac{1}{\rho} \vec{\nabla}_z p + \vec{F}_r$$

temporal change horizontal & vertical advection Coriolis force pressure gradient friction

Hydrostatic equation:

$$\frac{\partial p}{\partial z} = -g\rho$$

Equation of state:

$$p = \rho RT$$

The 3D Primitive Equations: diabatic effects

Continuity equation:

$$\frac{\partial \rho}{\partial t} + \vec{\nabla}_z \cdot (\rho \vec{v}) = 0$$

Thermodynamic equation:

$$\frac{D\Theta}{Dt} = \frac{\partial \Theta}{\partial t} + (\vec{v} \vec{\nabla}) \Theta = \frac{1}{c_p} \left(\frac{p_0}{p} \right)^{R_d / c_p} Q$$

Q: diabatic heating

Conservation of water vapor mixing ratio q:

$$\frac{\partial q}{\partial t} + \vec{v}_h \cdot \nabla_z q + w \frac{\partial q}{\partial z} = S_q$$

S_q : sources/sinks

+ Conservation laws for liquid water + ice

Explicit Horizontal Diffusion

- Diffusion applied to the prognostic variables
 - Regular diffusion ∇^2 - operator
 - Hyper-diffusion ∇^4 , ∇^6 , ∇^8 - operators: more scale-selective
 - Example: Temperature diffusion, $i = 1, 2, 3, \dots$

$$\frac{\partial T}{\partial t} = \dots - (-1)^i K^{(2i)} \left(\nabla^{(2i)} T \right)$$

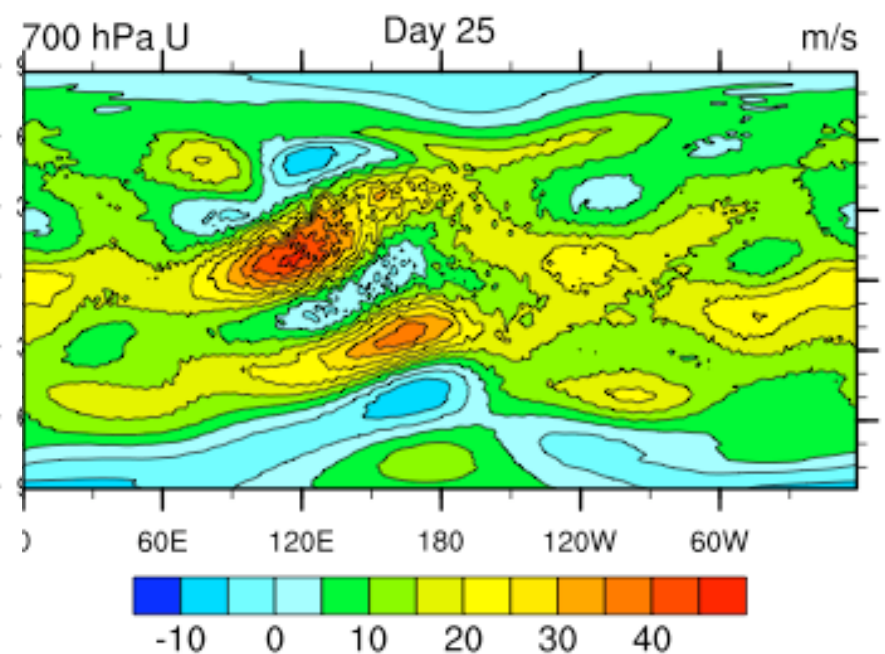
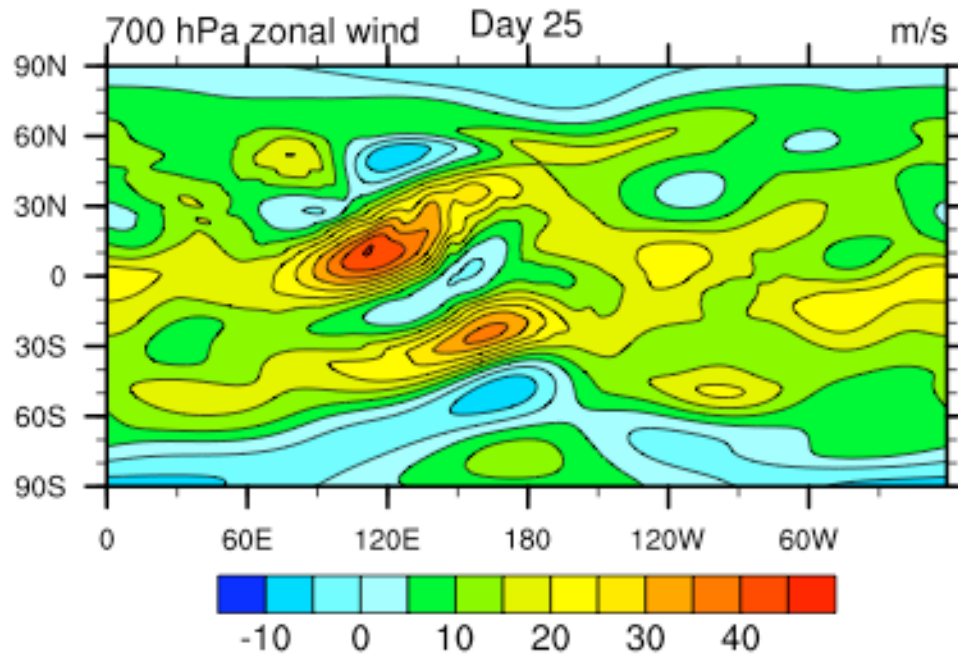
- K: diffusion coefficients, e-folding time dependent on the resolution
 - Choice of the prognostic variables and levels
- Divergence damping

Effects of Horizontal Diffusion

- Comparison of the 700 hPa zonal wind at day 25 in CAM FV and CAM EUL with test 5-0-0

CAM FV 1°x1°L26

CAM EUL T106L26



with monotonicity constraint,
divergence damping

with standard horizontal diffusion

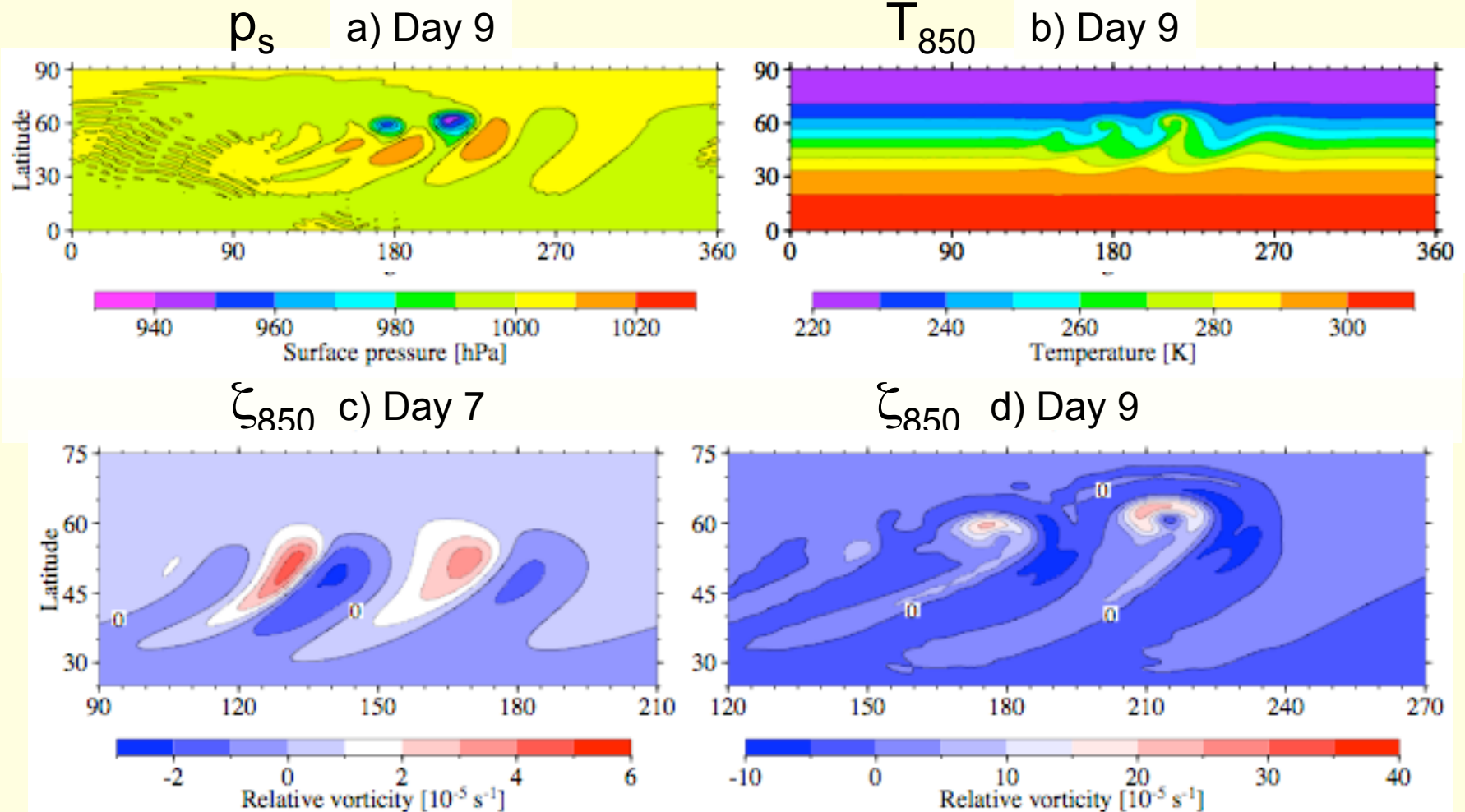
Horizontal Diffusion Coefficients

- Diffusion coefficients are scale-dependent
- Are guided by the so-called e-folding time: How quickly are the fastest waves damped so that their amplitude decrease by a factor of 'e'?
- Typical 4th-order diffusion coefficients K_4 for CAM EUL

Eulerian spectral transform dynamical core(EUL)				
Spectral Resolution	# Grid points lat × lon	Grid distance at the equator	Time step Δt	Diffusion coefficient K_4 ($\text{m}^4 \text{s}^{-1}$)
T21	32 × 64	625 km	2400 s	2.0×10^{16}
T42	64 × 128	313 km	1200 s	1.0×10^{16}
T85	128 × 256	156 km	600 s	1.0×10^{15}
T106	160 × 320	125 km	450 s	0.5×10^{15}
T170	256 × 512	78 km	300 s	1.5×10^{14}
T340	512 × 1024	39 km	150 s	1.5×10^{13}

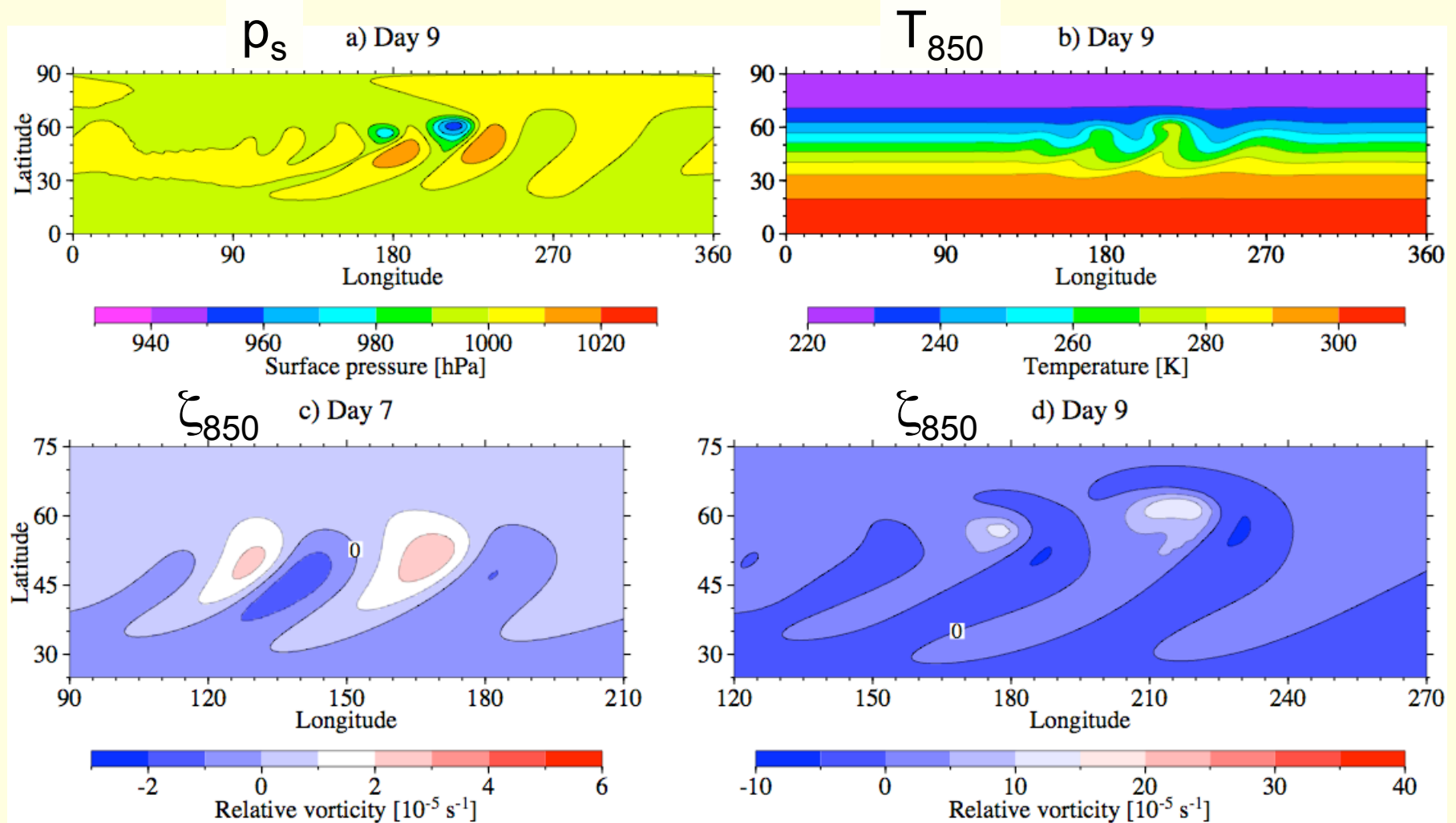
Impact of Explicit Diffusion: Baroclinic Waves

- EUL T85L26 with **standard** $K_4 = 10^{15} \text{ m}^4/\text{s}$ diffusion coefficient
- Spectral noise (Gibb's oscillations), test 2-0-0



Impact of Explicit Diffusion: Baroclinic Waves

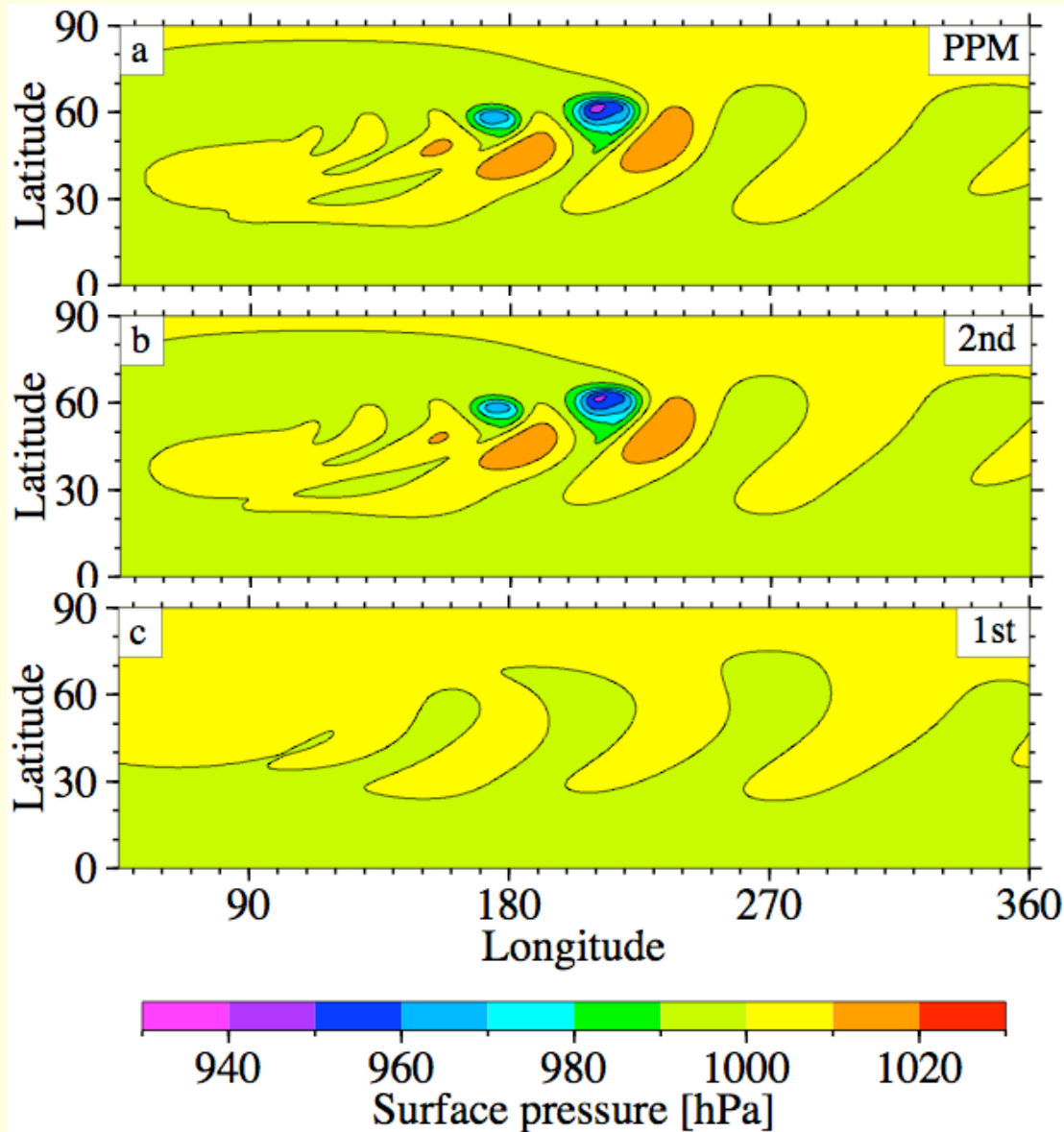
- EUL T85L26 with K_4 increased by a factor of 10 ($10^{16} \text{ m}^4/\text{s}$)
- No spectral noise, but **severe damping** of the circulation



Implicit / Numerical Diffusion

- Implicit diffusion: diffusion that is inherent in the numerical scheme
- Sources of implicit / numerical diffusion:
 - Order of accuracy: 1st order, 2nd order, 3rd order, ..., higher order schemes
 - The higher the order, the less diffusive
 - Monotonicity constraints
 - Decentering parameters in semi-implicit time-stepping schemes

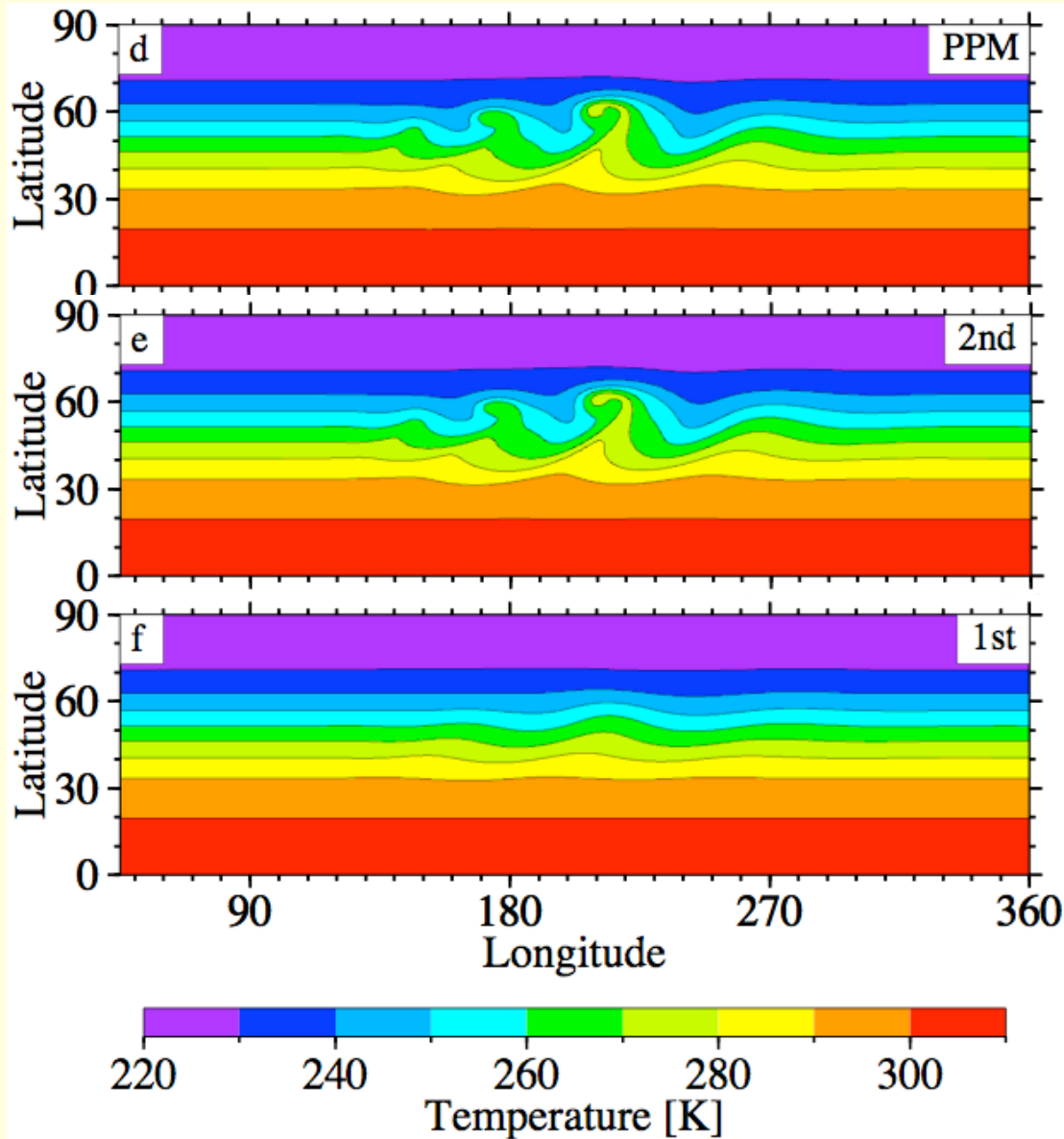
Implicit diffusion: Order of accuracy



- Third order (PPM)
- Second order (van Leer)
- First order upwind scheme

Test 2-0-0
CAM FV $1^\circ \times 1.25^\circ$ L26
 p_s at day 9

Implicit diffusion: Order of accuracy



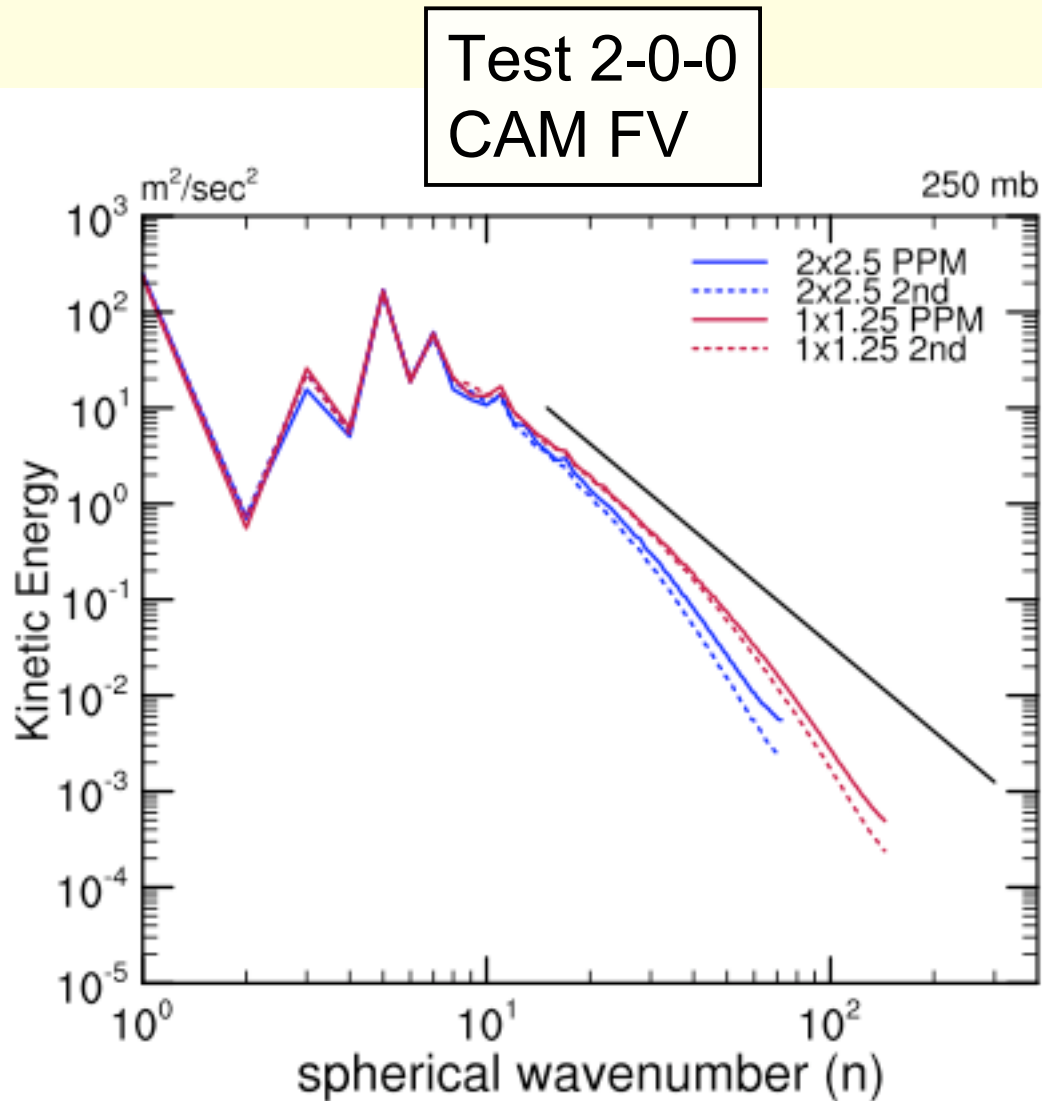
- Third order (PPM)
- Second order (van Leer)
- First order upwind scheme

Test 2-0-0

CAM FV $1^\circ \times 1.25^\circ$ L26

$T_{850 \text{ hPa}}$ at day 9

Implicit diffusion: Order of accuracy



- Time-averaged kinetic energy spectrum at two different horizontal resolutions
- Third order (PPM)
- Second order (van Leer scheme)
- Tail of 2nd order scheme drops faster

provided by
D. Williamson (NCAR)

Implicit diffusion: Monotonicity constraints in Finite Volume Methods

- Linear subgrid distribution (van Leer scheme)

Reconstruction: $h(x, y) = \bar{h} + \Delta a^x x + \Delta a^y y$

Slopes: $\Delta a^x = \frac{1}{2} (h_{i+1,j} - h_{i-1,j})$

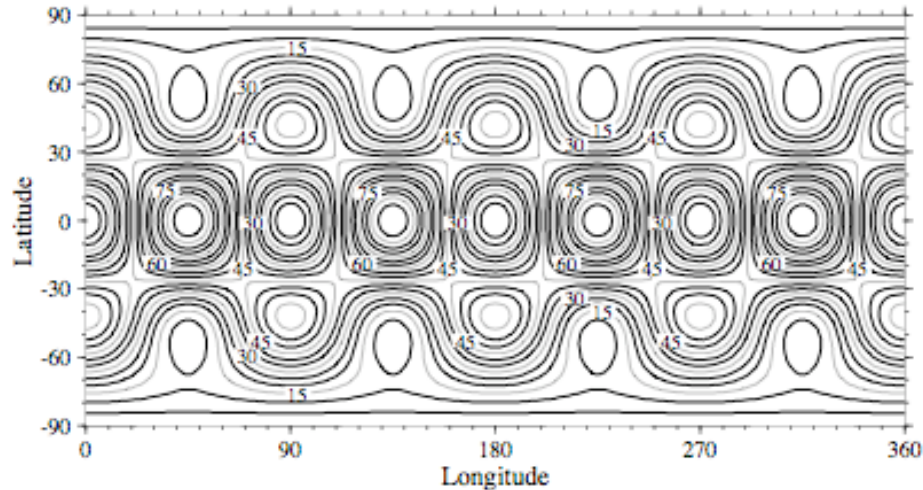
$$\Delta a^y = \frac{1}{2} (h_{i,j+1} - h_{i,j-1})$$

Slope limiter: $\Delta a^x = \min(|\Delta a^x|, 2|h_{i+1,j} - h_{i,j}|, 2|h_{i,j} - h_{i-1,j}|) \operatorname{sgn}(\Delta a^x)$
if $(h_{i+1,j} - h_{i,j})(h_{i,j} - h_{i-1,j}) > 0$
= 0 otherwise

- Parabolic subgrid distribution (PPM) with cross terms

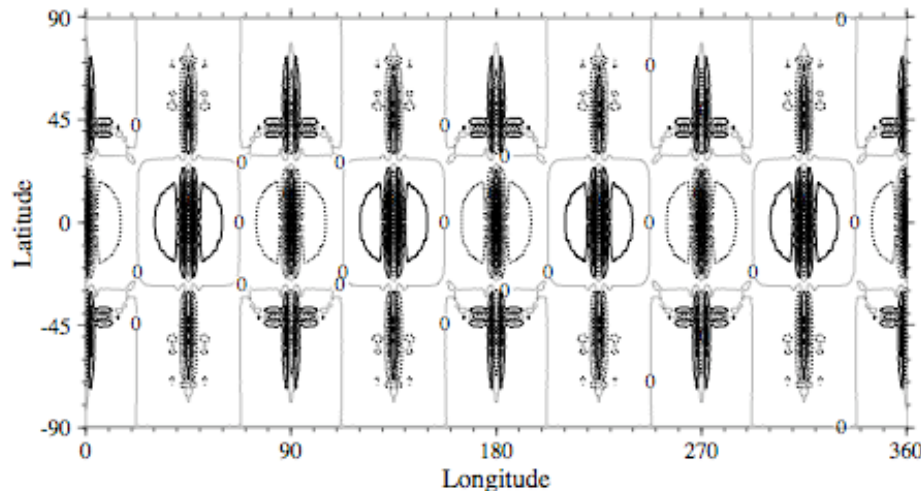
$$h(x, y) = \bar{h} + \delta a^x x + b^x \left(\frac{1}{12} - x^2 \right) + \delta a^y y + b^y \left(\frac{1}{12} - y^2 \right) + \frac{1}{2} (c^{xy} + c^{yx}) x y$$

Implicit diffusion: Monotonicity constraint



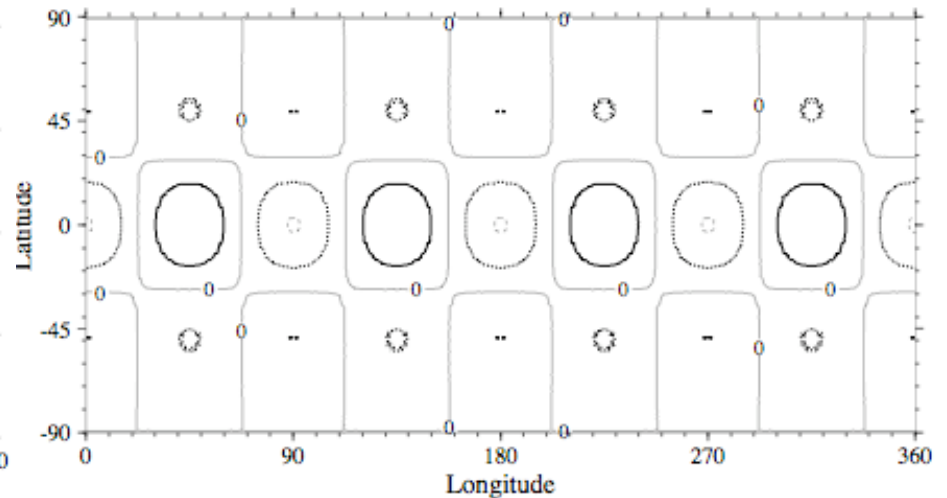
- SW Rossby-Haurwitz wave
- Initial u field at $2^\circ \times 2.5^\circ$
- Split cells to $1^\circ \times 1.25^\circ$ grid and interpolate via a PPM reconstruction, compare to analytical solution (error)

Error PPM constrained



Errors cluster near the extrema where the monotonicity constraint is strongest

Error PPM unconstrained



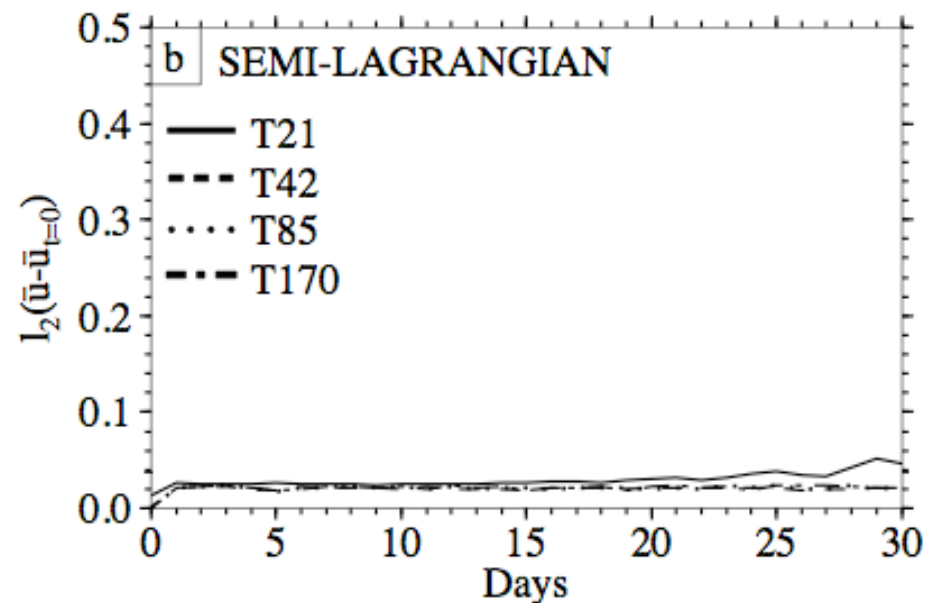
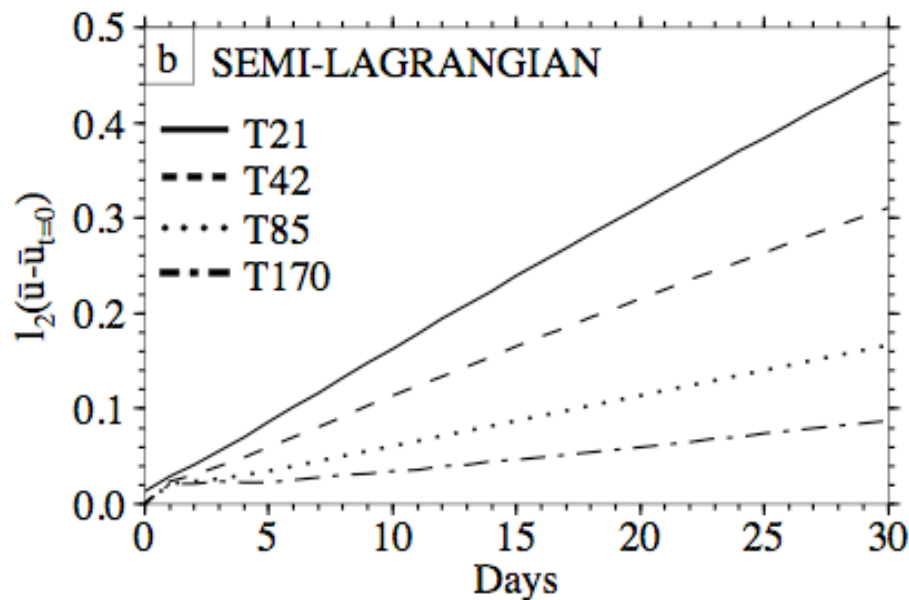
Errors are reduced, but over- or undershoots are possible

Decentering mechanism (semi-implicit)

- Decentering mechanism is used in the semi-implicit semi-Lagrangian model CAM SLD, parameter ε
- Decentering technique damps noise induced by orographic resonance, ε needed in real simulations
- Damping clearly shown in test 1-0-0, I_2 error (Eqn. 18)

$\varepsilon = 0.2$

$\varepsilon = 0$



Divergence damping

- Example: 2D shallow water momentum equation

Momentum equation:

$$\frac{\partial \vec{v}}{\partial t} = -\Omega_a \vec{k} \times \vec{v} - \nabla \left(\Phi + \mathcal{K} - \overset{\text{coefficient}}{\downarrow} c D \right)$$

Horizontal divergence:

$$D = \frac{1}{a \cos \varphi} \left[\frac{\partial u}{\partial \lambda} + \frac{\partial}{\partial \varphi} (v \cos \varphi) \right]$$
$$\approx \frac{1}{a \cos \varphi} \left[\frac{\Delta u}{\Delta \lambda} + \frac{\Delta (v \cos \varphi)}{\Delta \varphi} \right]$$

Semi-discretized:

$$\frac{\partial \vec{v}}{\partial t} = -\Omega_a \vec{k} \times \vec{v} - \nabla \left(\Phi + \mathcal{K} - [c_u \Delta u + c_v \Delta (v \cos \varphi)] \right)$$

Divergence damping coefficients divided by metric terms, different in both directions

Divergence damping

- Divergence damping diffuses the divergent part of the flow

$$\frac{\partial \vec{v}_h}{\partial t} = -\Omega_a \vec{k} \times \vec{v}_h - \nabla(\Phi + K - cD) \quad (\text{SW equation})$$

$$\Rightarrow \frac{\partial \vec{v}_h}{\partial t} = \dots + \nabla(cD)$$

$$\Rightarrow \nabla \cdot \frac{\partial \vec{v}_h}{\partial t} = \dots + \nabla \cdot \nabla(cD) \quad \text{Apply divergence operator}$$

$$\Leftrightarrow \frac{\partial D}{\partial t} = \dots + \nabla^2(cD) \quad \text{with } D : \text{horizontal divergence}$$



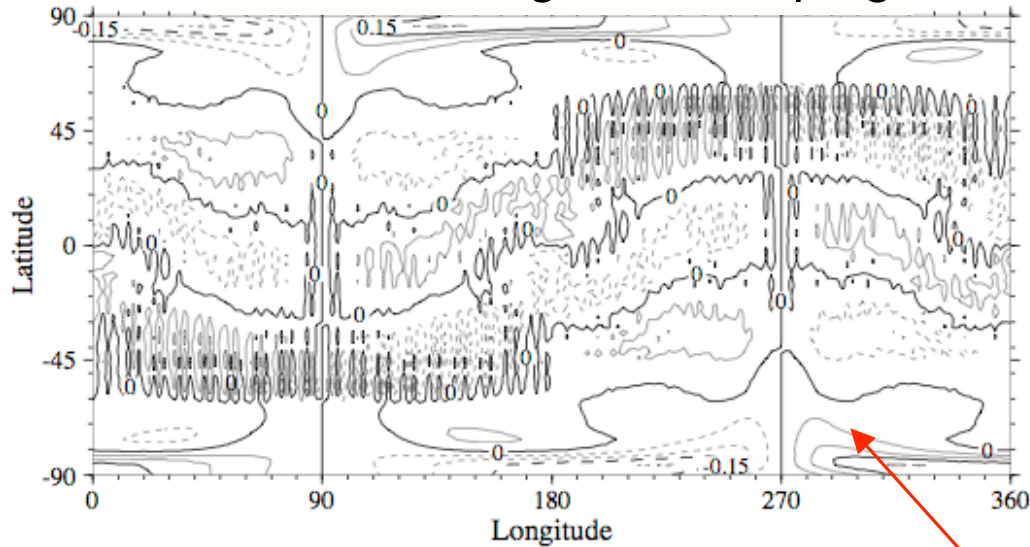
2nd order
diffusion

Spatially variant divergence
damping coefficient, units m^2/s

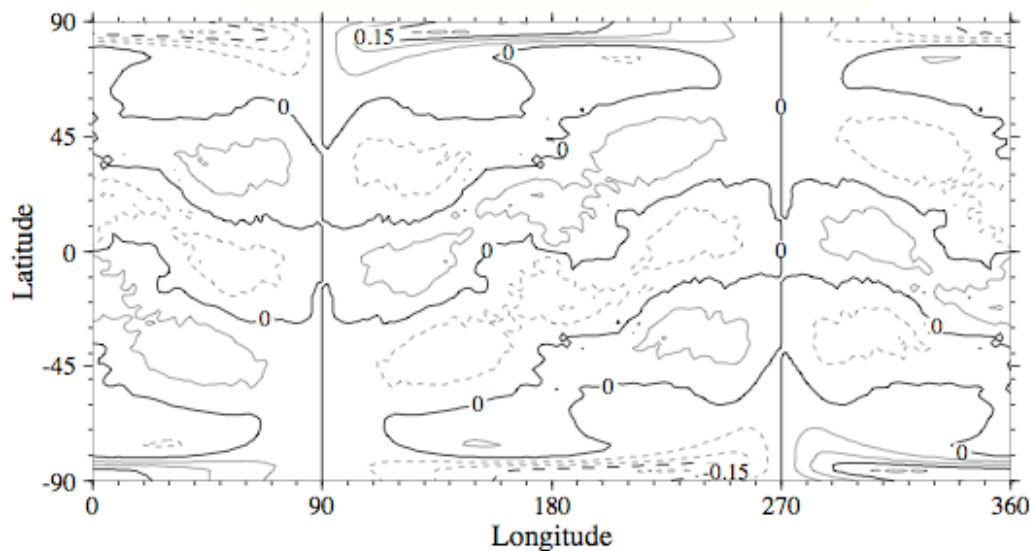
- Can you select any coefficient c ? Selection criterion?

Divergence damping

without divergence damping

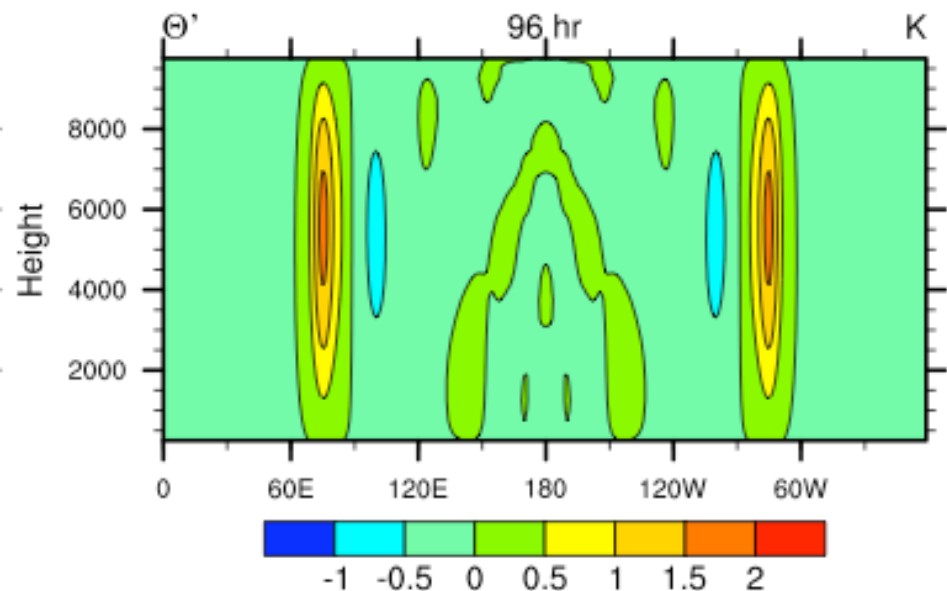
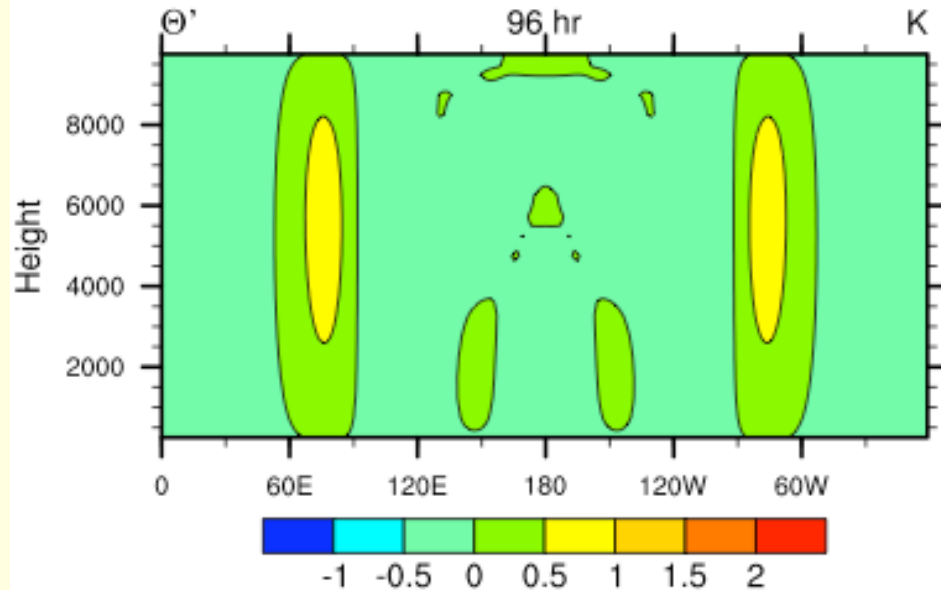


with divergence damping



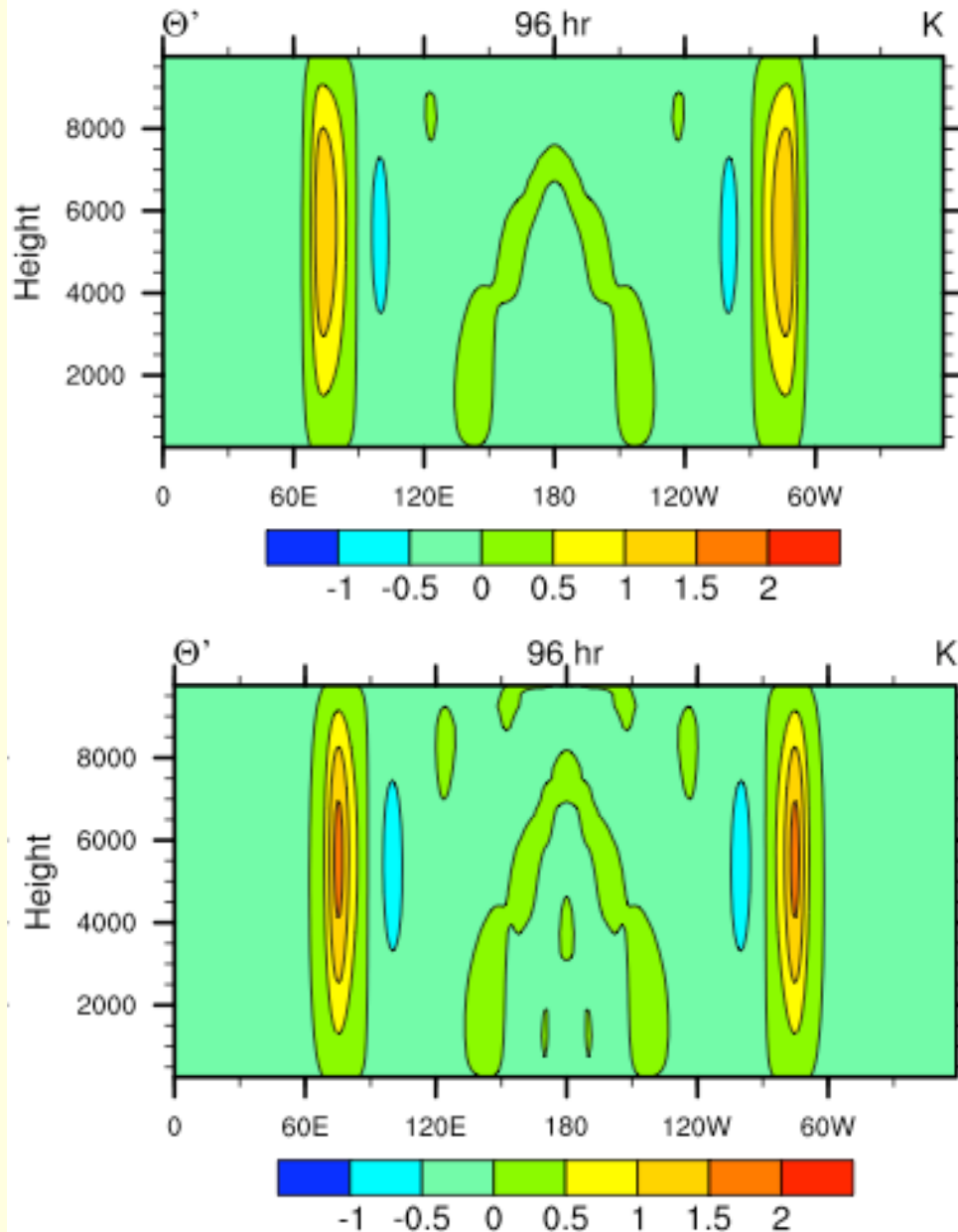
- Example: 2D SW steady state test case with $\alpha=90^\circ$, model FV
- Difference field at day 10 compared to analytical solution
- Contour interval is 0.05 m/s
- Why is the polar region always smooth? Because of other filters (here polar Fourier filter)

Divergence damping: Effects



- Example: 3D gravity wave test 6-0-0
- Model CAM FV 1°x 1° L20 at day 4, cross section at equator
- with standard divergence damping coefficients (top)
- without divergence damping (bottom)
- Clear difference in the amplitudes of the gravity wave

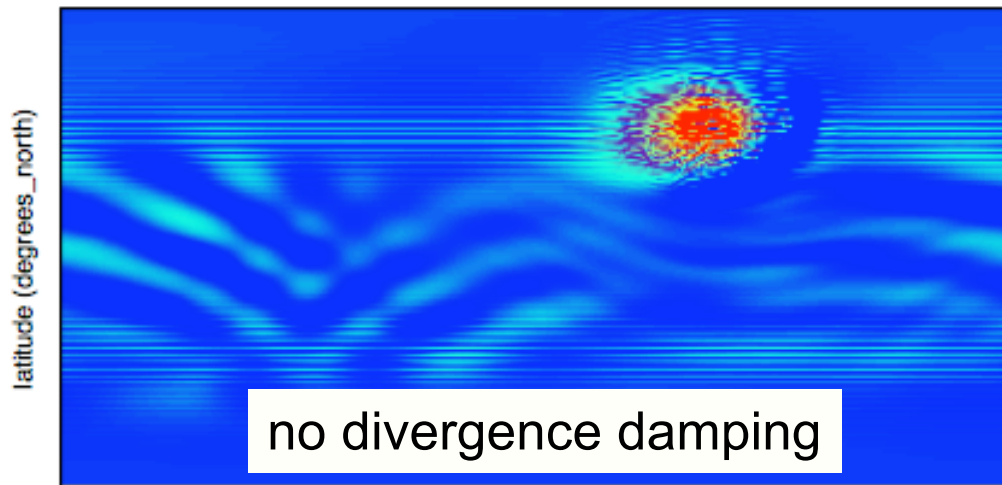
All types of diffusion change the solution



- Example: 3D gravity wave test 6-0-0, cross section at the equator at day 4
- Model CAM EUL T106L20 with explicit ∇^4 diffusion (top)
- Model CAM FV $1^\circ \times 1^\circ$ L20, no divergence damping (bottom)
- Clear difference in the shape of the potential temperature perturbation
- Check sharpness of leading edge

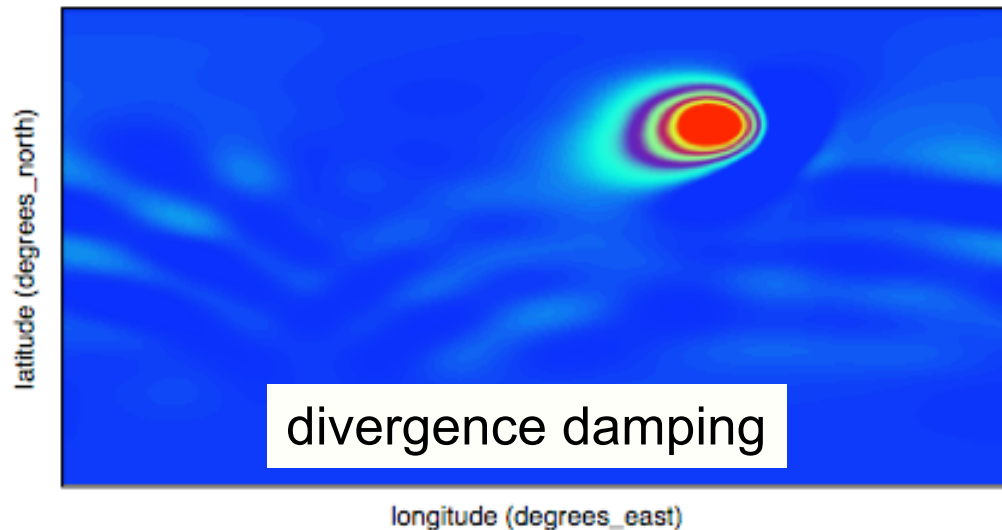
Divergence damping: Needed for stability ?

Temperature at 850 mbar pressure surface (K)



cjablano Fri Jun 6 09:27:08 2008

Temperature at 850 mbar pressure surface (K)

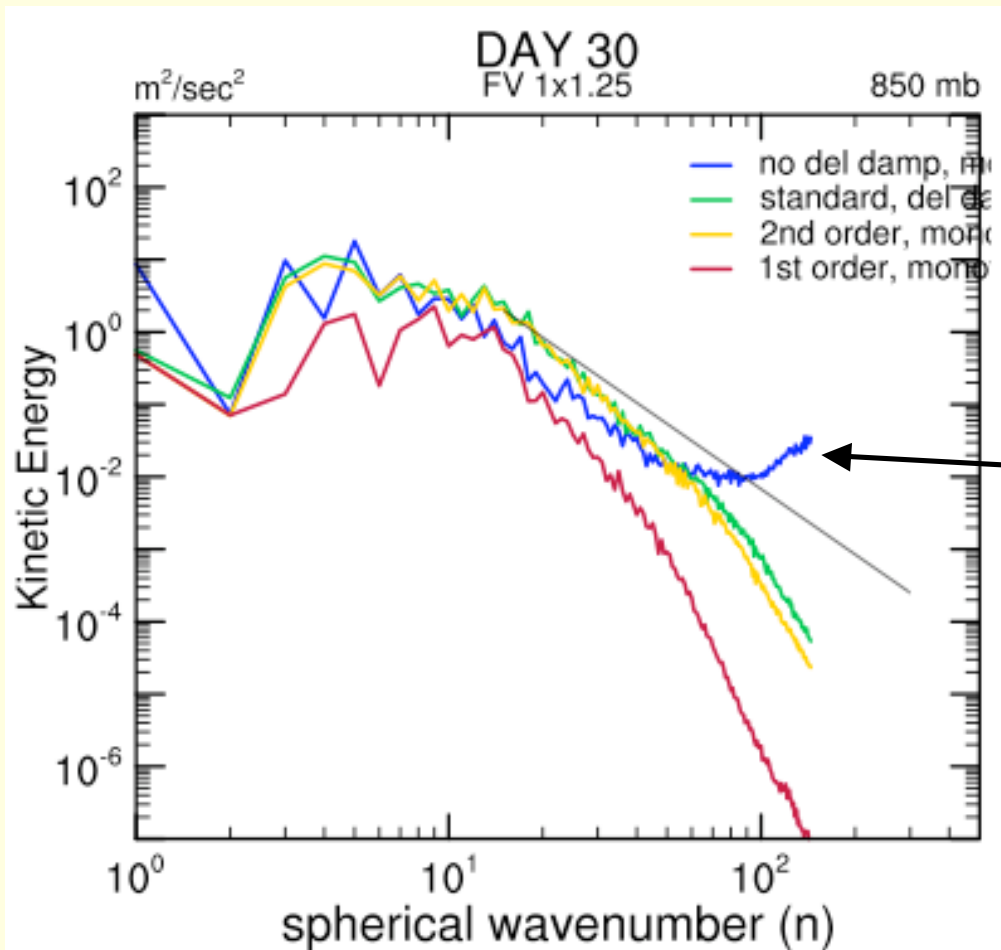


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- Example: alternative 3D inertio-gravity wave test with background flow
- Model CAM FV $1^\circ \times 1^\circ$ L20 at day 5.5, lat-lon cross section at 850 hPa
- Numerical stability of CAM FV depends on the resolution- and time step dependent choice of the divergence damping coefficient c

Divergence Damping

- Effects of the divergence damping and order of accuracy on the Kinetic Energy spectrum (test 2-0-0)



Blue: PPM, no divergence damping

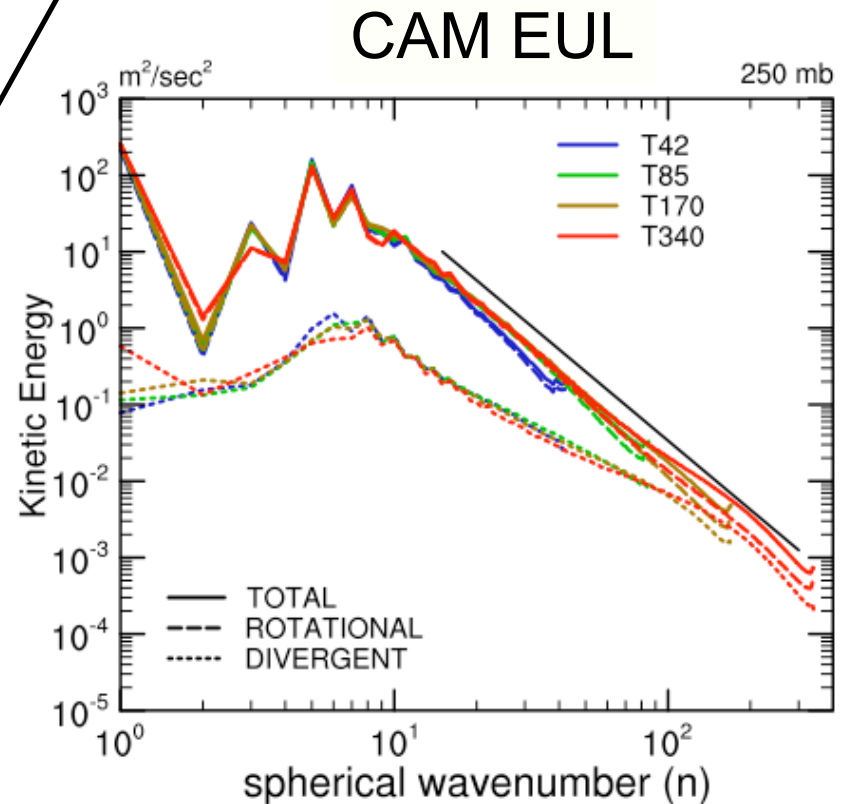
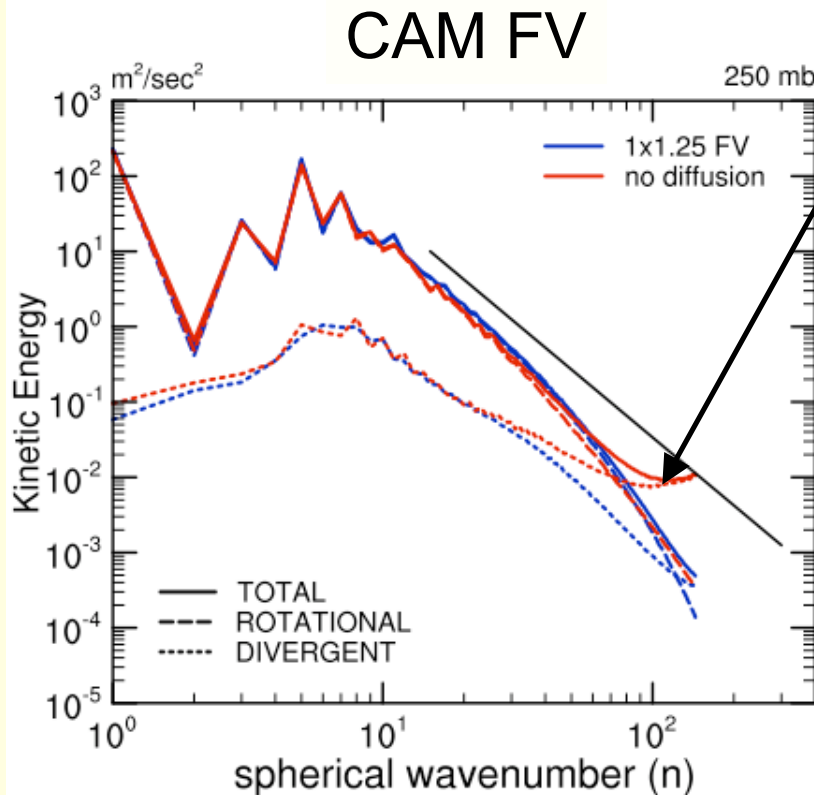
Green: PPM, standard divergence damping

Accumulation of energy at small scales without divergence damping

Model: CAM FV,
plot provided by
D. Williamson (NCAR)

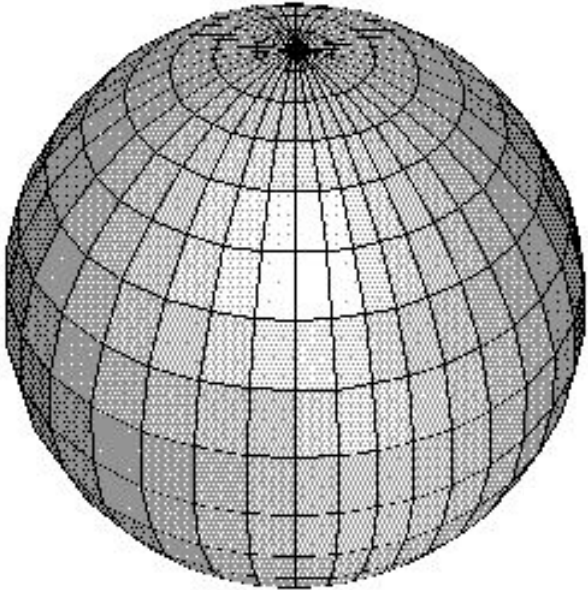
Divergence Damping

- Without diffusion (here divergence damping):
divergent part of the flow responsible for the hook



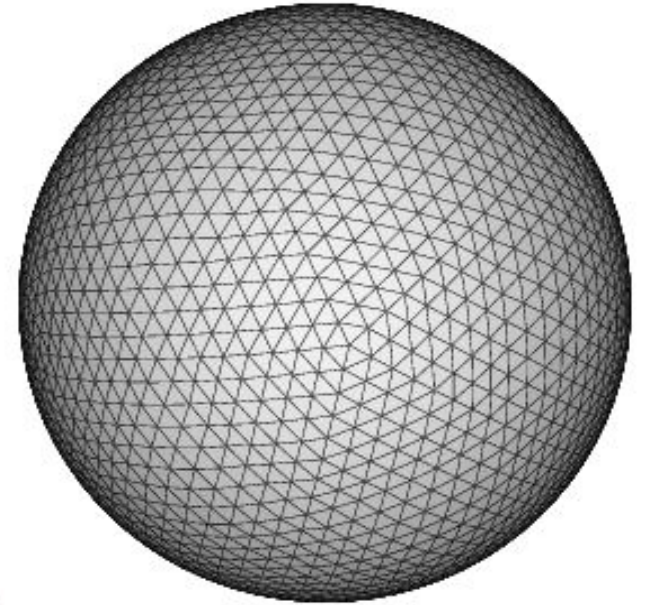
plots provided by D. Williamson (NCAR)

Computational grids (horizontal)



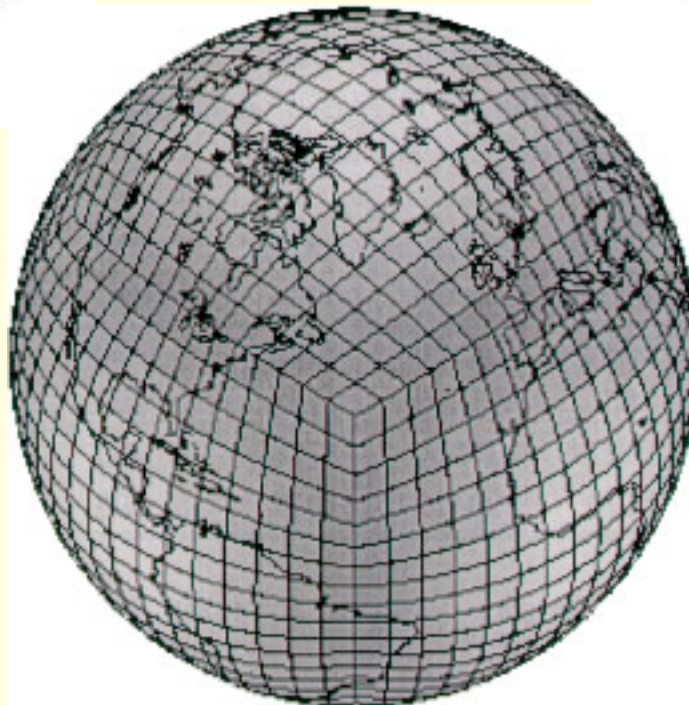
Latitude-longitude grid: **needs polar filtering due to convergence of meridians**

Cubed sphere



Icosahedral grid

No polar filter requires

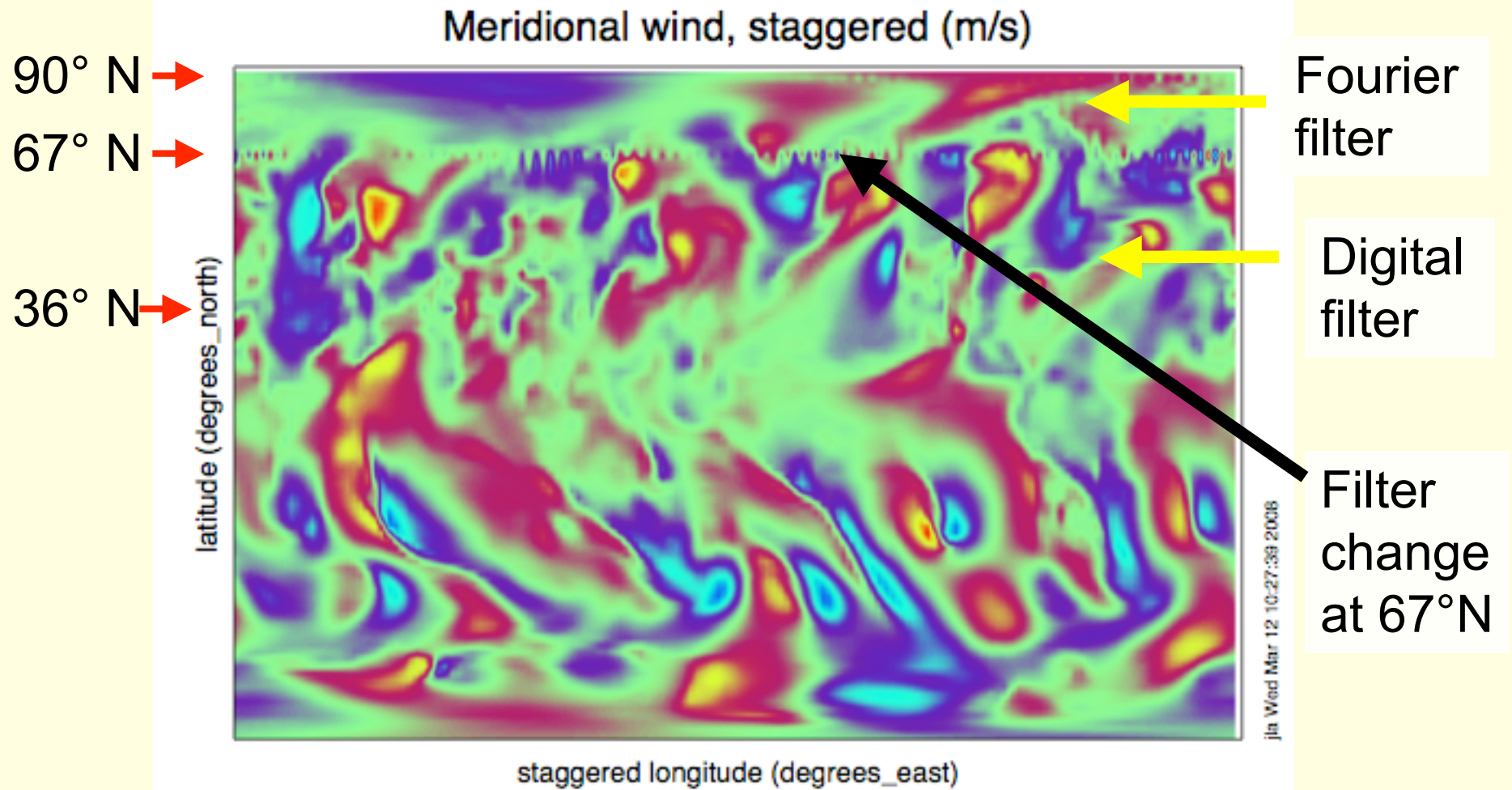


Spatial filters

- Most popular and most effective polar filter: 1D Fourier filter (spectral filter), used in the zonal (x) direction
- **Basic idea:**
 - Transform the grid point data into spectral space via Fourier transformations
 - Eliminate or damp high wave numbers (noise) by either setting the spectral coefficients to 0 or multiplying them with a damping coefficient $\in [0, 1]$
 - Transform the field back from spectral space into grid point space: result is a filtered data set
- Filter strength is determined by the spectral damping coefficients, can be made very scale-selective and dependent on the latitude (e.g. less strong towards equator)
- Drawback: needs all data along latitude ring (poor scaling)

Spatial filters: Fourier & Digital Filter

- Data assimilation run with CAM FV, D-grid v field at 266 hPa



provided by Jeff Anderson, NCAR

Digital filters

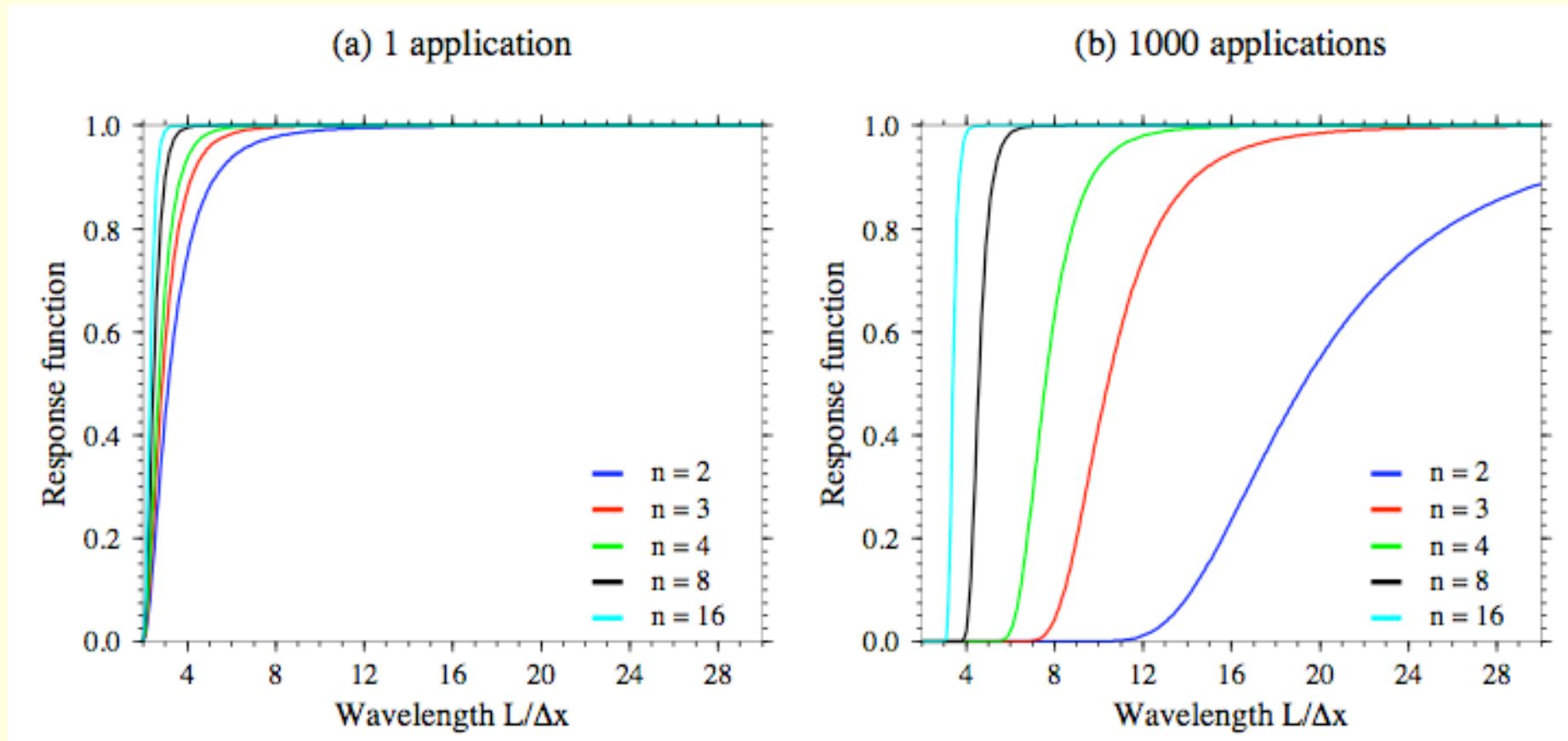
- Digital or algebraic filters are local grid-point filters that only take neighboring grid points into account
- Examples are the Shapiro filters (Shapiro, 1975)
- 2nd order Shapiro filter (i is the grid point index):

$$\bar{f}_i = \frac{1}{16} (-f_{i-2} + 4f_{i-1} + 10f_i + 4f_{i+1} - f_{i+2})$$

- The filter response/damping function is (Shapiro, 1971)

$$\begin{aligned} \rho_n(k) &= 1 - \sin^{2n} \left(k \frac{\Delta x}{2} \right) && 2n: \text{ order} \\ &= 1 - \sin^{2n} \left(\pi \frac{\Delta x}{L} \right) \end{aligned}$$

Digital filters: Response function

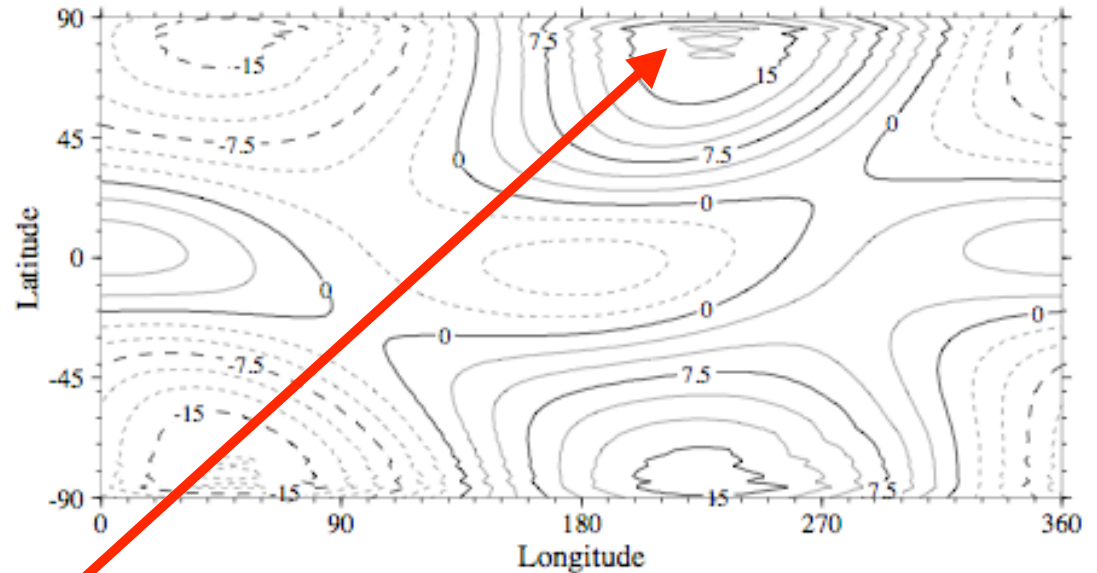


- Response function of different Shapiro filters after (a) 1 application and (b) 1000 applications. n indicates the order of the Shapiro filter. Higher orders need more data points.

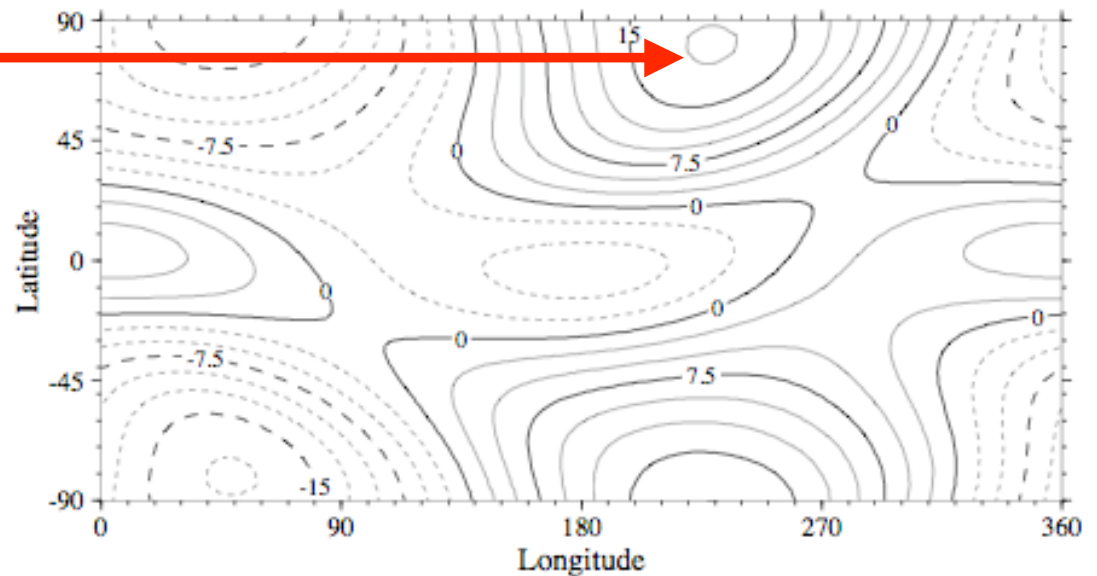
Digital filters

- Can provide a strong damping effect
- Use very selectively
- Example: SW simulation, digital filtering in y-direction applied near the pole points

(a) No Shapiro filtering in meridional direction

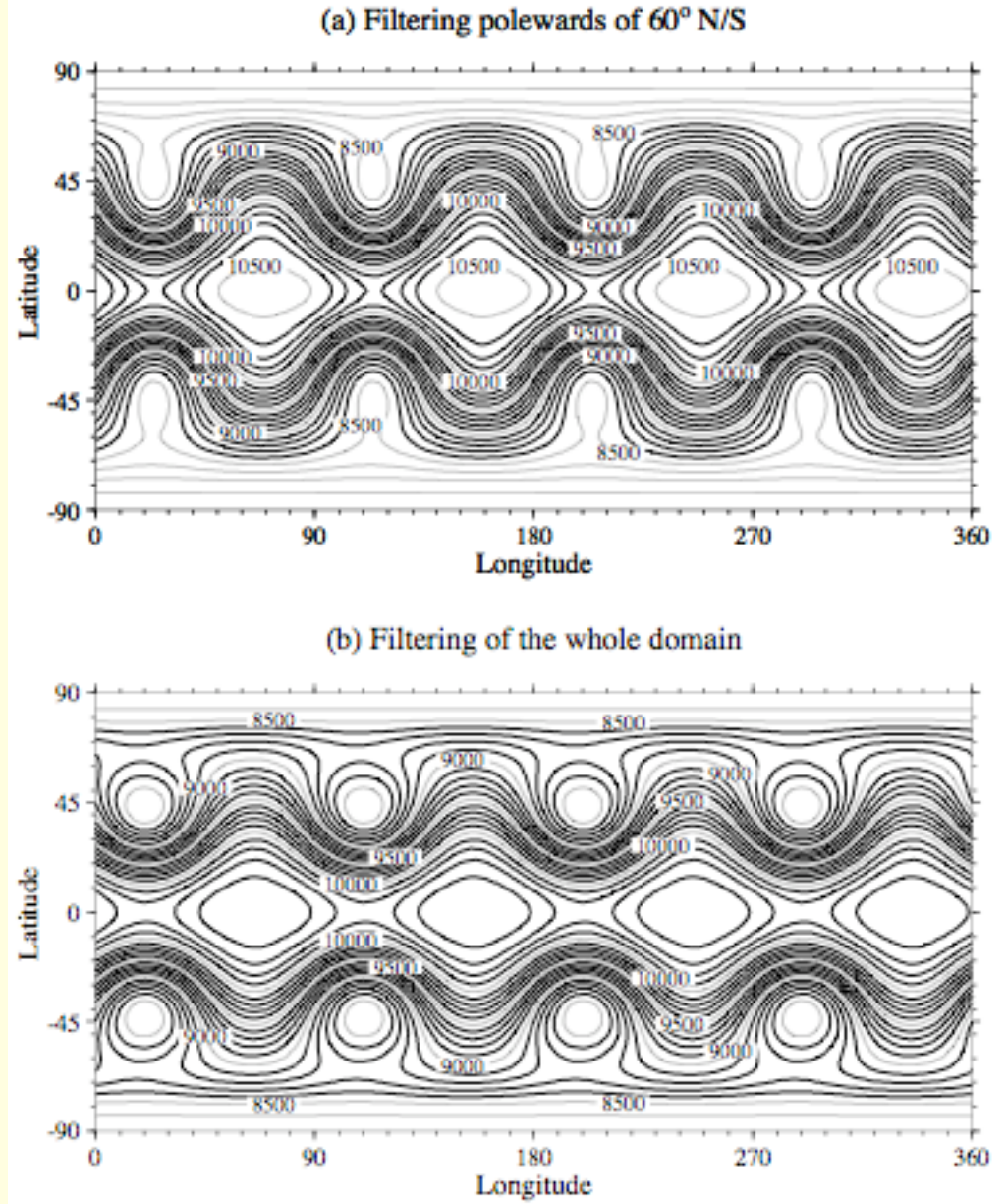


(b) Shapiro filtering in meridional direction



Spatial Filters

- Can provide a strong damping effect
- Example: Rossby-Haurwitz wave in SW FV model, height at day 14
- (a) Fourier (90° - 75° N/S) and digital Shapiro filtering (75° - 60° N/S)
- (b) Digital Shapiro filter also applied between 60° N - 60° S, very diffusive, not suitable



Time filters

- Used in models with 3-time level schemes
- Most often used: Asselin filter (Asselin, 1972)
- Avoids that the even and odd time steps separate
- Basic idea: Second-order diffusion in time
- Example with time levels $n-1$, n , $n+1$:

$$\overline{\psi}^n = \psi^n + \alpha \left(\overline{\psi}^{n-1} - 2\psi^n + \psi^{n+1} \right)$$

- Filter strength is determined by the coefficient α
- Often used $\alpha \approx 0.05$

Conservation of Mass

- Conservation of the (dry) air mass is only guaranteed if the continuity equation is written in conservative form:

$$\frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot (\rho \vec{v}) = 0$$

→ requires the density ρ to be a prognostic variable

- Alternative form for Lagrangian vertical coordinates:

$$\frac{\partial}{\partial t} \int_{\xi_l}^{\xi_u} \frac{\partial p}{\partial \xi} d\xi + \nabla_{\xi} \cdot \left(\vec{v} \int_{\xi_l}^{\xi_u} \frac{\partial p}{\partial \xi} d\xi \right) + \frac{\partial}{\partial \xi} \int_{\xi_l}^{\xi_u} \left(\dot{\xi} \frac{\partial p}{\partial \xi} \right) d\xi = 0 \quad \text{Integrate}$$

→
$$\frac{\partial}{\partial t} \delta p + \nabla_{\xi} \cdot (\vec{v} \delta p) = 0$$

Pressure thickness δp
is prognostic variable

Conservation of Mass: Mass fixers

- Evaluate mass conservation properties of some models in the colloquium:
ICON, CAM FV, CAM EUL
- Be careful what you see: Some models, especially climate models, apply *a posteriori* mass fixers
- Conservation of mass is needed in long-term climate simulations, less important in short weather prediction runs
- Basic idea behind the mass fixer: adjust the mean value of p_s after each time step, adjustment modifies all grid points at the surface
- This technique does not alter the pressure gradients
- Ask your modeling mentor!

Conservation of Total Energy

- There are many forms of the Total Energy (TE or E) Equation that depend on the choice of the fluid dynamics equations and the vertical coordinate (see Appendix F)
- An example for hydrostatic models with Cartesian coordinates is

$$E = \int_A \int_{z_{top}}^{z_s} \left(\frac{\mathbf{v}^2}{2} + c_v T + gz \right) \rho dz dA$$
$$\approx \int_A \left[\sum_{k=1}^{K_{max}} \left(\frac{u_k^2 + v_k^2}{2} + c_v T_k + gz_k \right) \rho_k \Delta z_k \right] dA.$$

- In general: The TE equation is a global integral of the kinetic, thermal and potential energy in the model.
- The global integral is conserved in the continuous equations.

Conservation of Total Energy

- The question is whether TE is a conserved quantity in a dynamical core with numerical discretizations.
- Should we care?
 - in Weather Prediction Models
 - The answer is ‘not necessarily’
 - in Climate Models
 - The answer is ‘yes’
- When running for long times the violation of the total energy conservation leads to artificial drifts in the climate system (e.g. ocean heat fluxes)

Total Energy Fixer

- In nature:
 - conservation of total energy
 - energy lost by molecular diffusion provides heat
- In atmospheric models:
 - Energy is lost due to explicit or implicit (numerical) diffusion processes
 - Molecular diffusion is not represented on the model grid (spatial scale in models is way too big)
 - Numerical scheme might also lead to increase in total energy
- Therefore: some models provide an *a posteriori* energy fixer that restores the conservation of total energy by modifying the temperature

A posteriori Total Energy Fixer

- Goal: Total energy at each time step should be constant
- Compute the residual: $RES = \hat{E}^+ - E^-$
- Compute the total energy before (-) and after (+) each time step

$$\hat{E}^+ = \int_A \left\{ \left[\sum_{k=1}^K \left(\frac{(\hat{\mathbf{v}}_k^+)^2}{2} + c_p \hat{T}_k^+ \right) (p_0 \Delta A_k + \hat{p}_s^+ \Delta B_k) \right] + \Phi_s \hat{p}_s^+ \right\} dA$$
$$E^- = \int_A \left\{ \left[\sum_{k=1}^K \left(\frac{(\mathbf{v}_k^-)^2}{2} + c_p T_k^- \right) (p_0 \Delta A_k + p_s^- \Delta B_k) \right] + \Phi_s p_s^- \right\} dA$$

A posteriori Total Energy Fixer

- **Idea:** Correct the temperature field to achieve the conservation of total energy (mimics heating by molecular diffusion)
- Option: **Fixer 1**, correction proportional to the magnitude of the local change in T at that time step

$$T^+(\lambda, \varphi, \eta) = \hat{T}^+(\lambda, \varphi, \eta) + \beta_1 |\hat{T}^+(\lambda, \varphi, \eta) - T^-(\lambda, \varphi, \eta)|$$

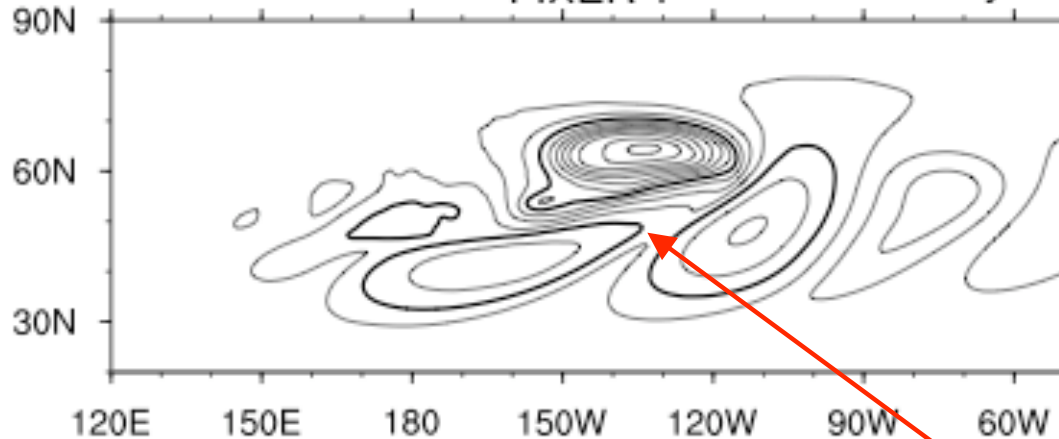
- Option: **Fixer 2**, correction is constant everywhere

$$T^+(\lambda, \varphi, \eta) = \hat{T}^+(\lambda, \varphi, \eta) + \beta_2$$

- Fixer 1 looks physical, but leads to wrong results

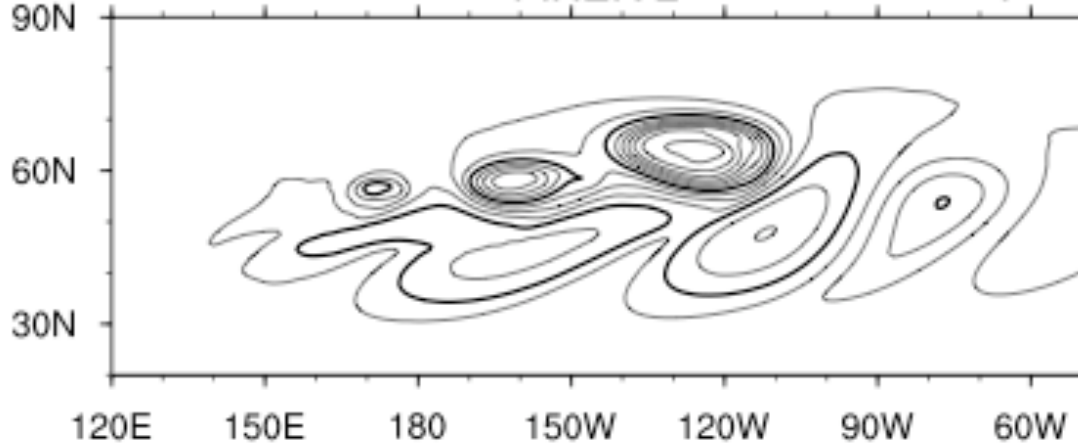
Energy Fixer: Surprising Consequences

wrong energy fixer FIXER 1 day 10



corrected energy fixer

FIXER 2 day 10

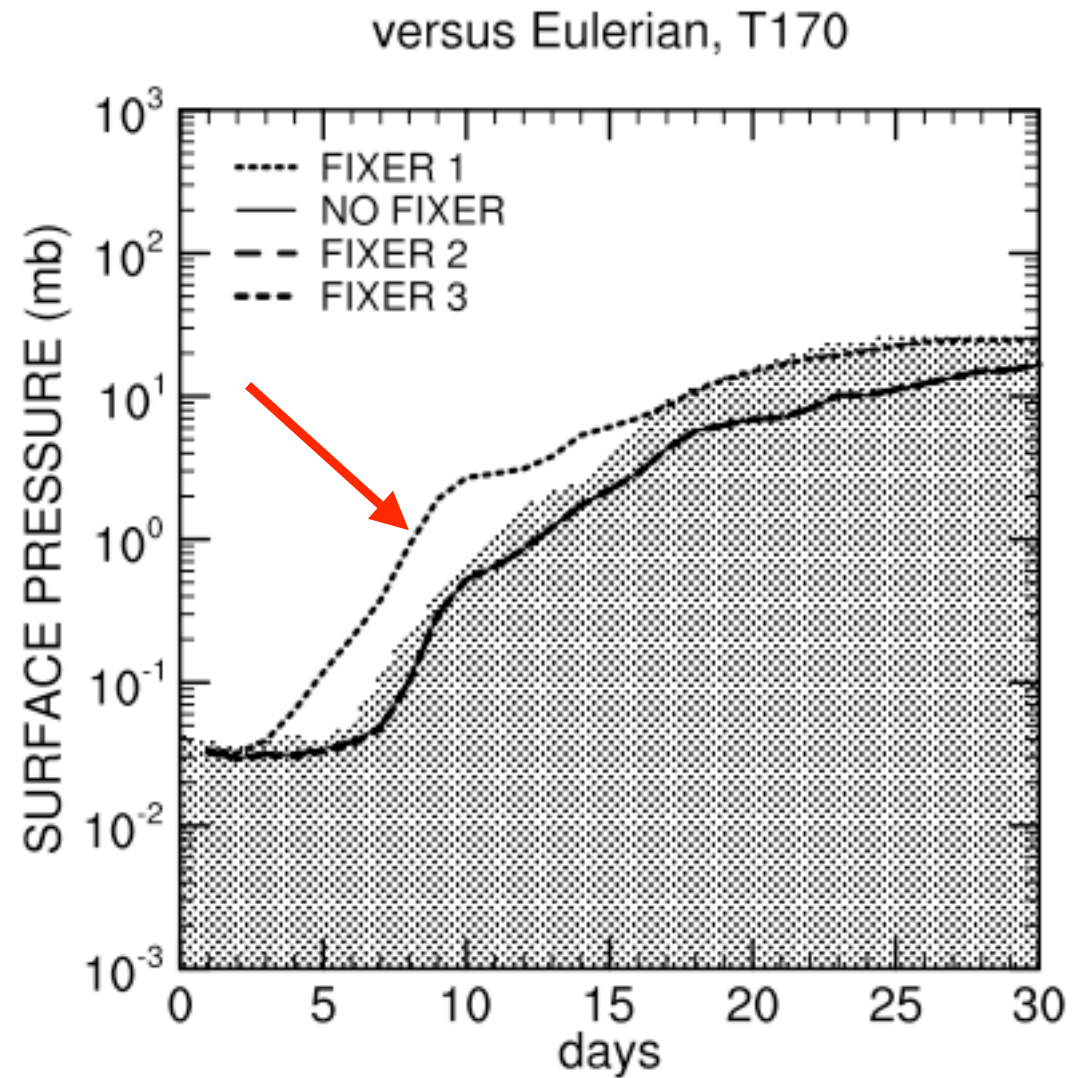


- Baroclinic wave test 2-0-0, p_s at day 10
- CAM SLD with a 'wrong' and 'corrected' choice of an energy fixer
- Wrong choice leads to wrong circulation pattern

Williamson, Olson & Jablonowski, MWR, in review

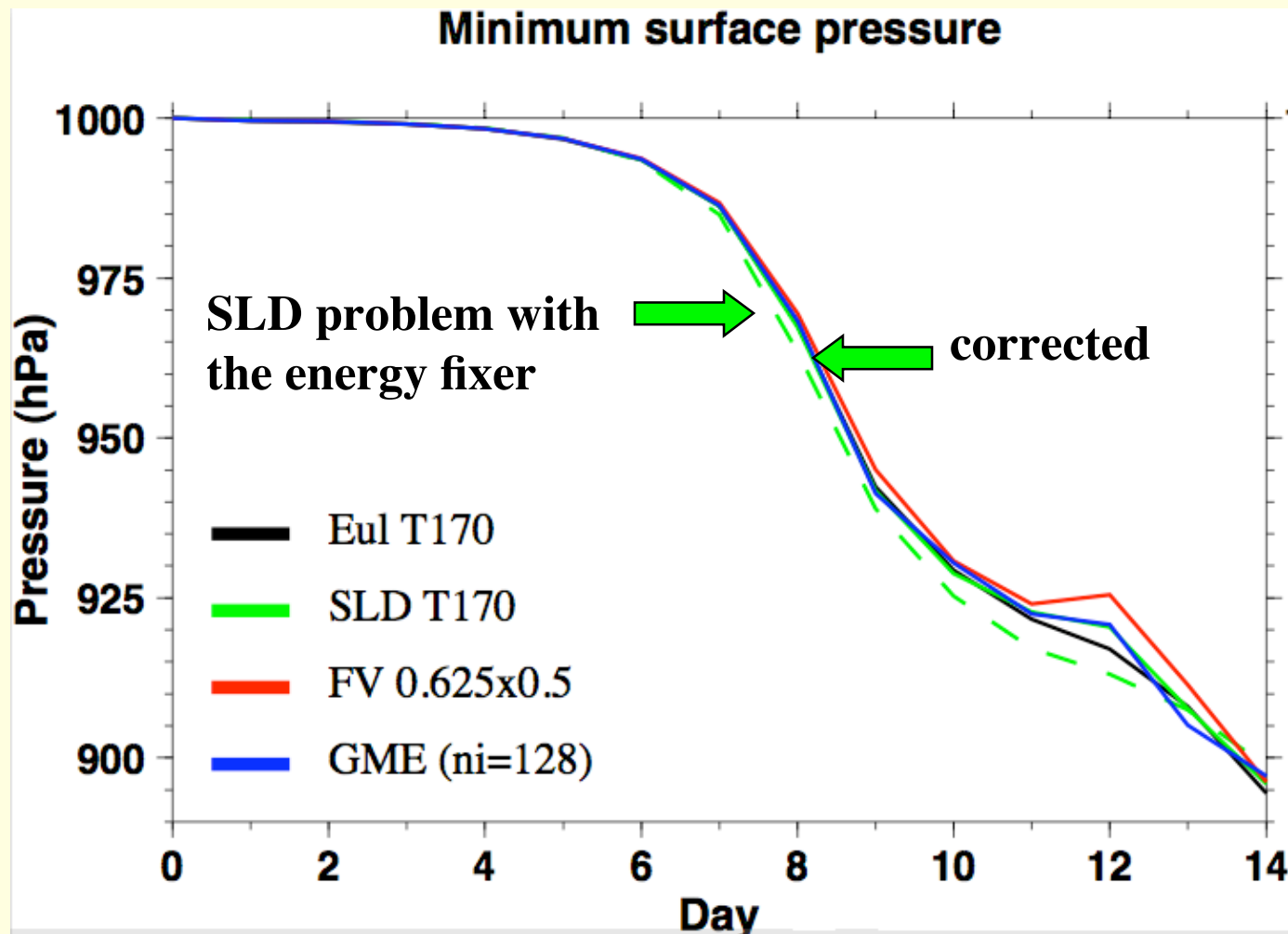
Energy Fixer: CAM SLD simulations

- Wrong choice (Fixer 1) is a clear outlier in the $I_2(p_s)$ error norm plot
- Lies above the uncertainty of the reference solutions (gray shaded)



Energy Fixer: SLD Dynamical Core

- Fixer 1 in the SLD simulation is also an outlier in the time series of the minimum surface pressure



Conclusions

- These are the modeling aspects that nobody will tell you unless you ask.
- Ask your modeling mentor lots of questions !!
- Diffusion and filters help maintain the numerical stability
- Some diffusion (either explicit or implicit) is always needed to prevent a build-up of energy at the smallest scale (due to truncated energy cascade)
- But: Use the techniques selectively and know their consequences.
- It is very easy to compute nice-looking smooth, highly diffusive, but very inaccurate solutions to the equations of motion.