

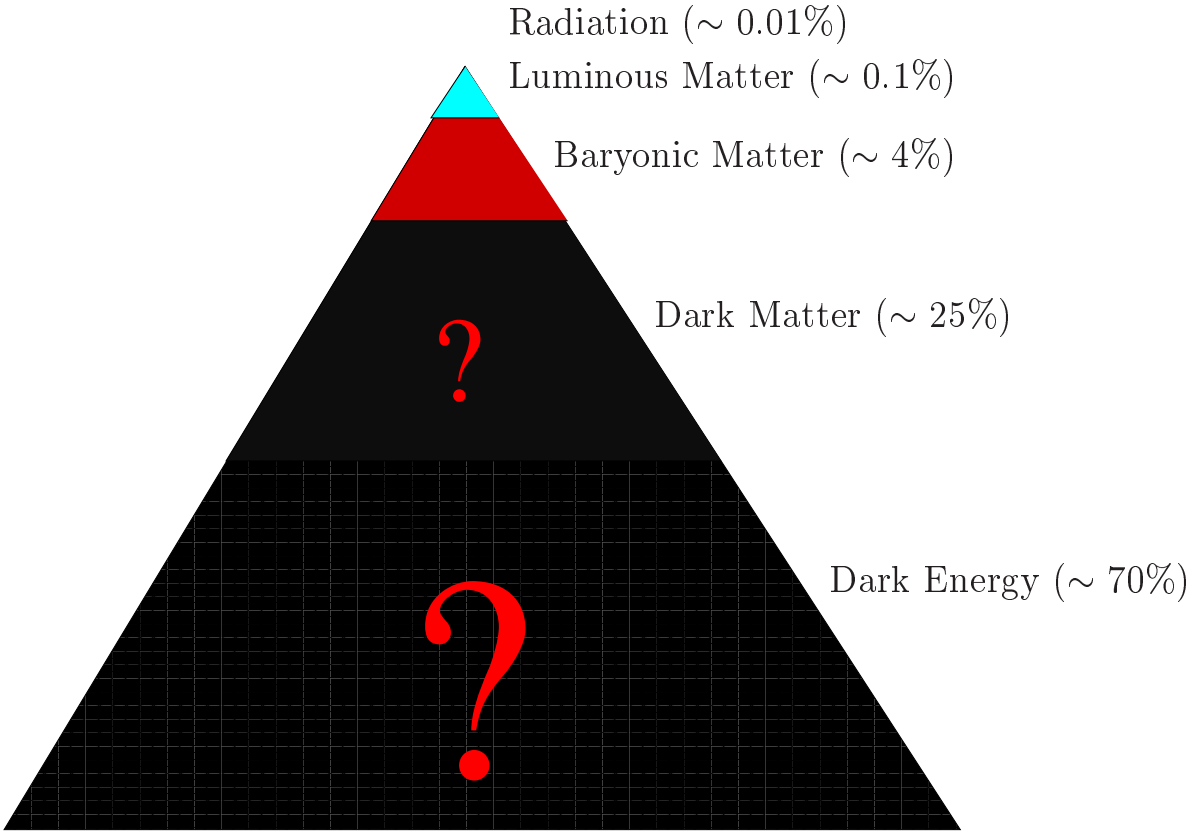
# CMB Windows on Dark Energy

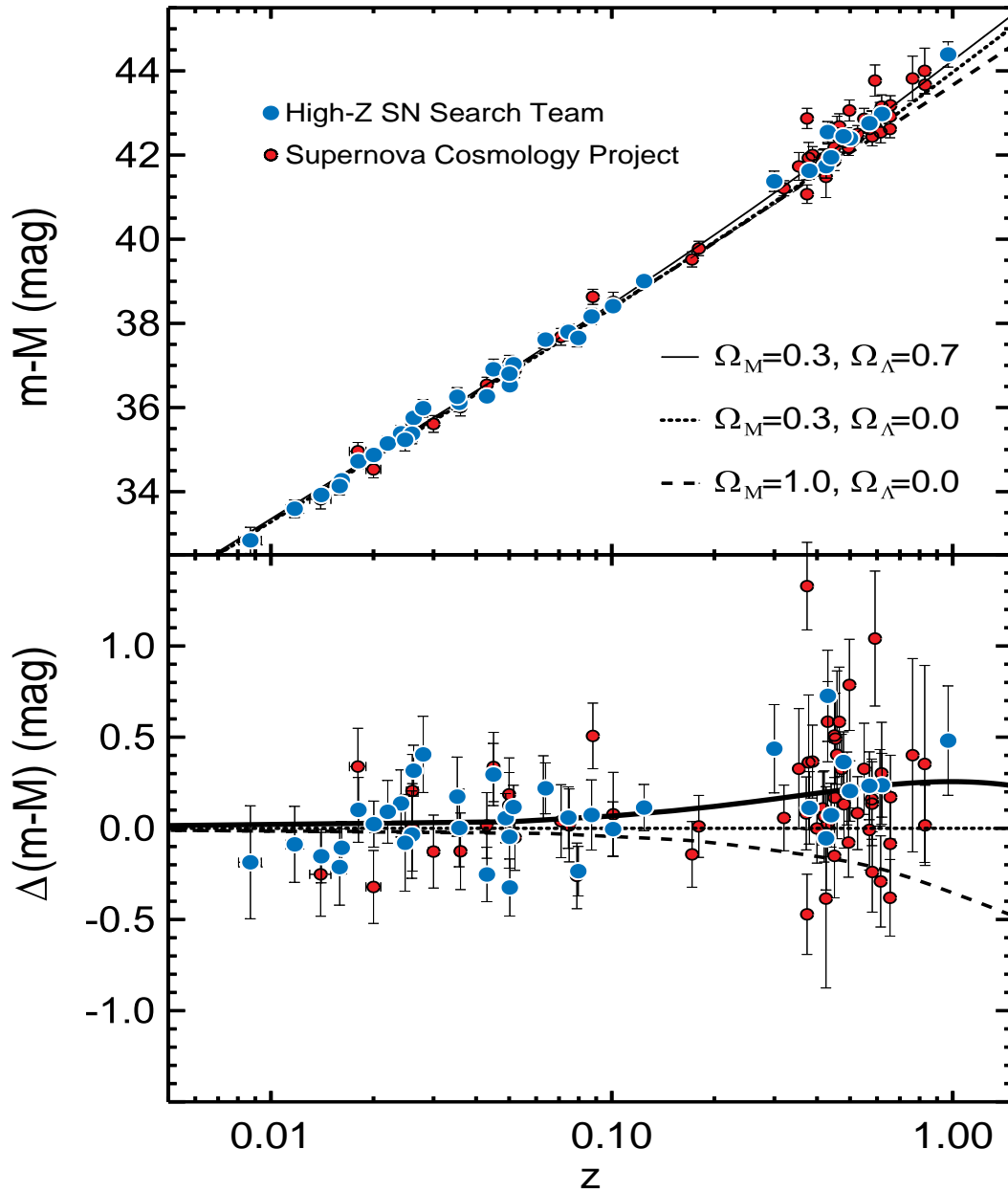
**Dragan Huterer**

(Case Western Reserve University)

LeBron James





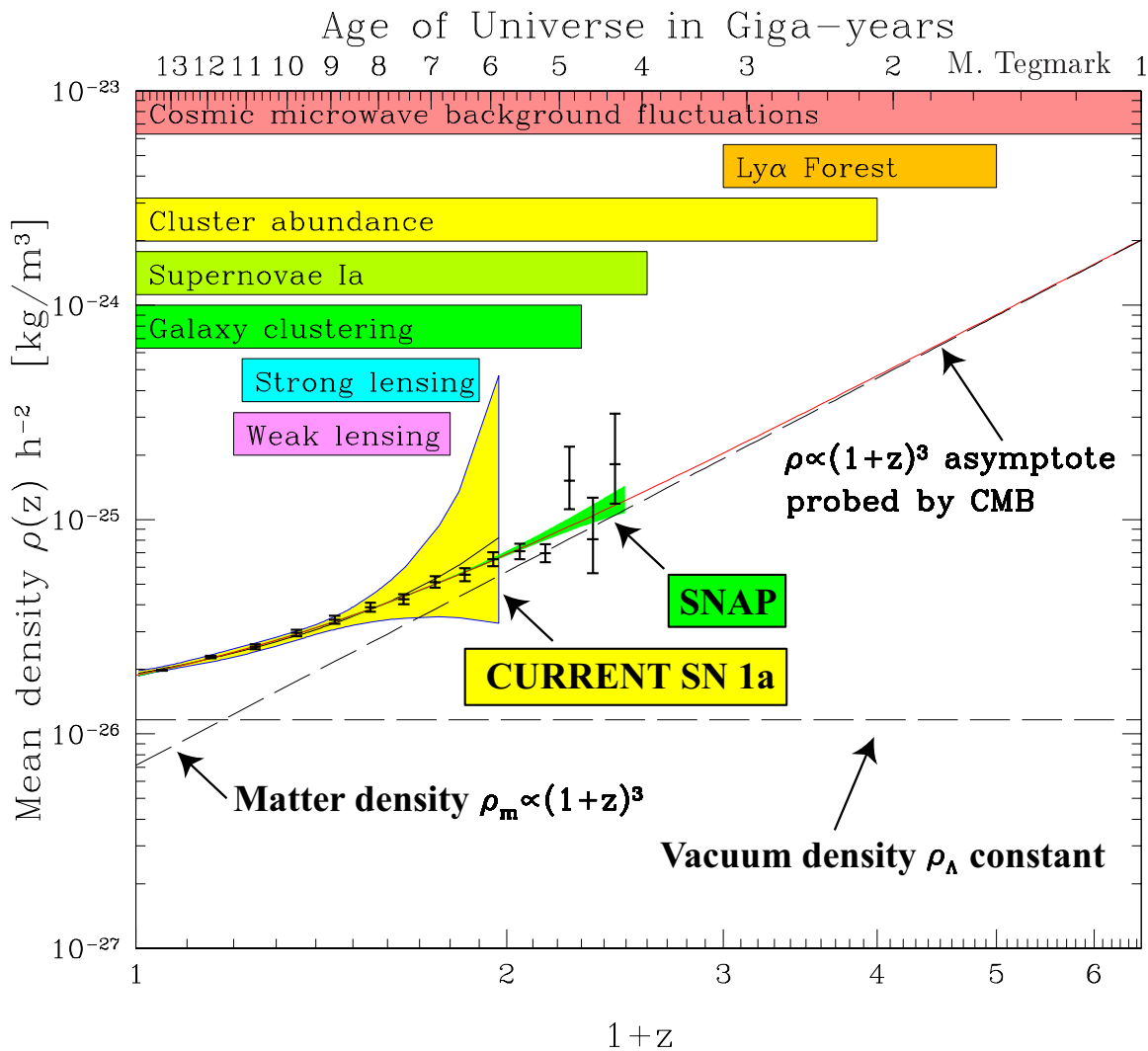


Describe dark energy by (Turner & White 1999)

$$\Omega_X = \frac{\rho_X}{\rho_{\text{crit}}}, \quad w = \frac{p_X}{\rho_X}$$

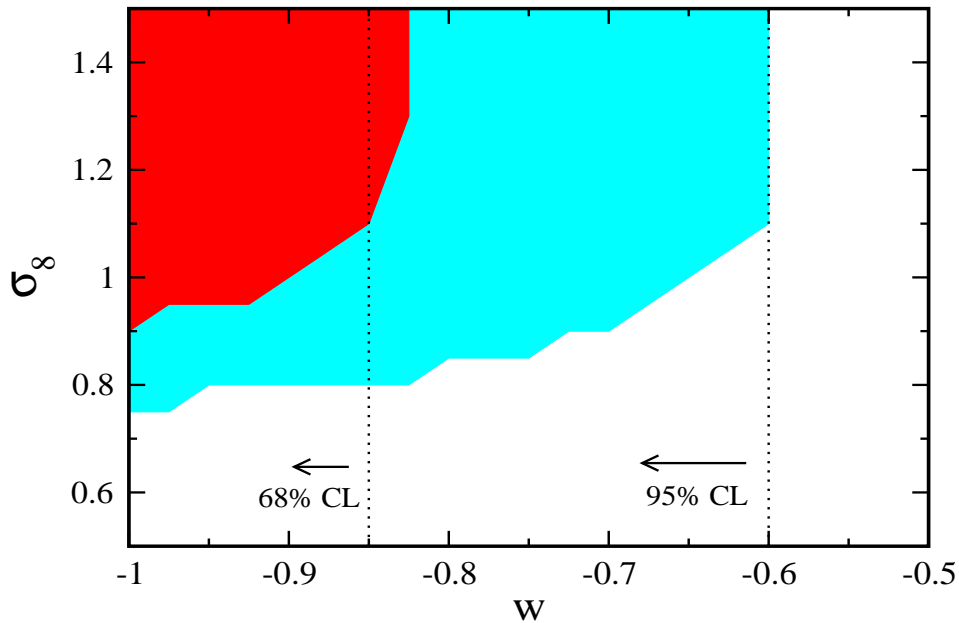
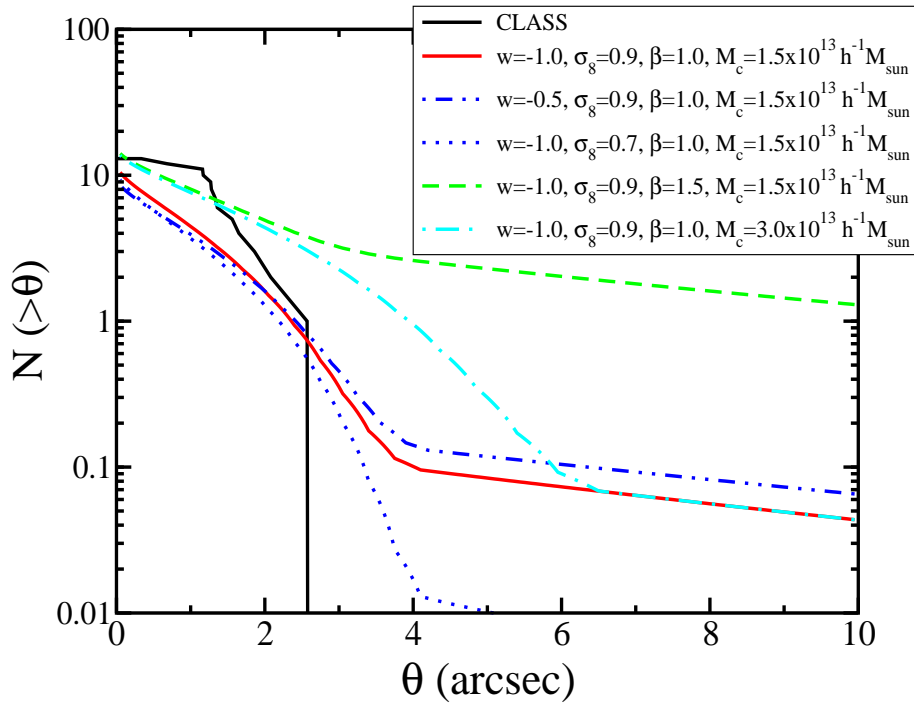
# Cosmological tests of Dark Energy

$$H^2(z) = H_0^2 \left[ \Omega_M(1+z)^3 + \Omega_K(1+z)^2 + \Omega_{DE}(1+z)^{3(1+w)} \right]$$



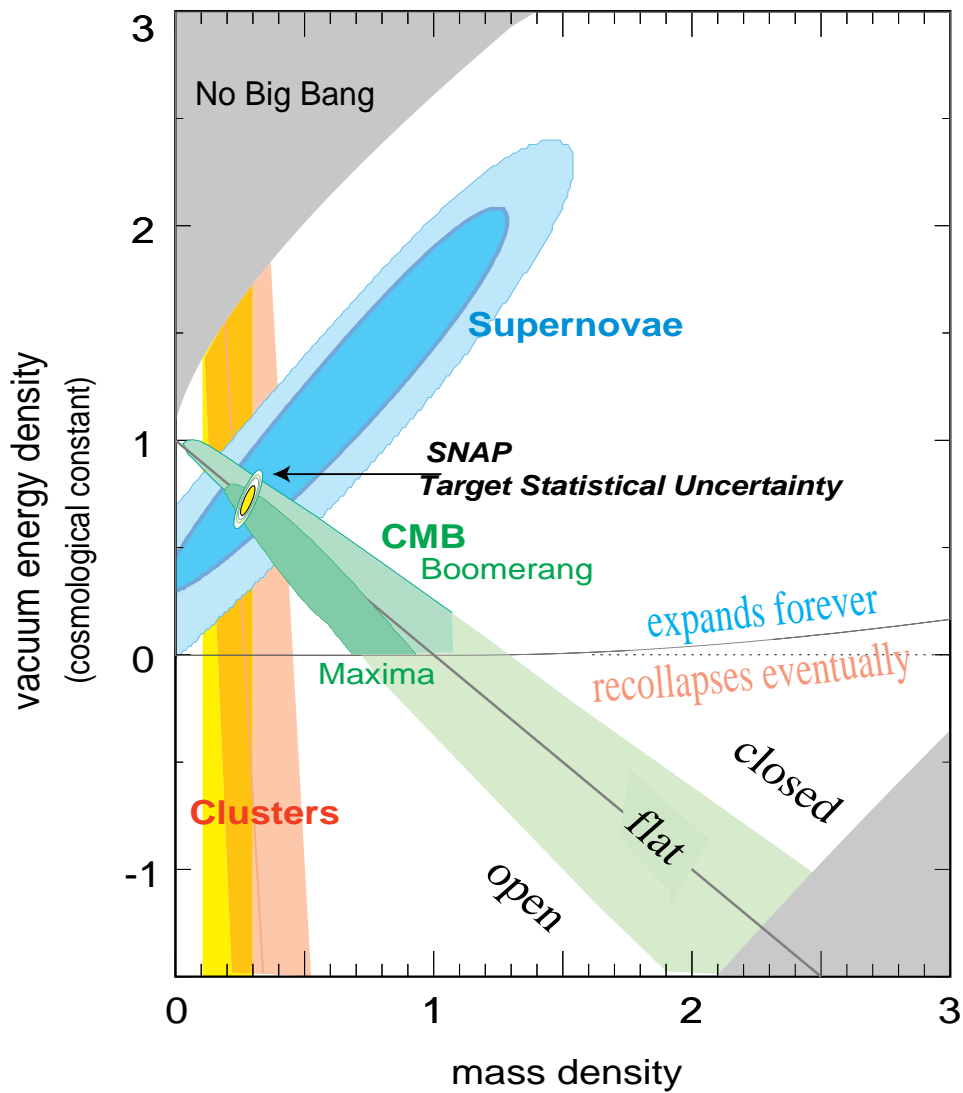
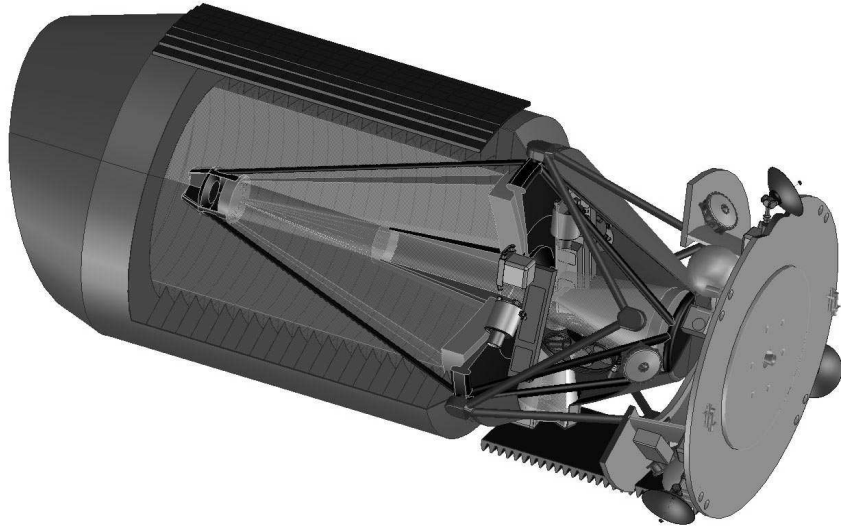
(Huterer & Turner, Weller & Albrecht, Kujat et al.,  
Linder, Hu, Tegmark, etc. etc.)

# Strong Gravitational Lensing Statistics

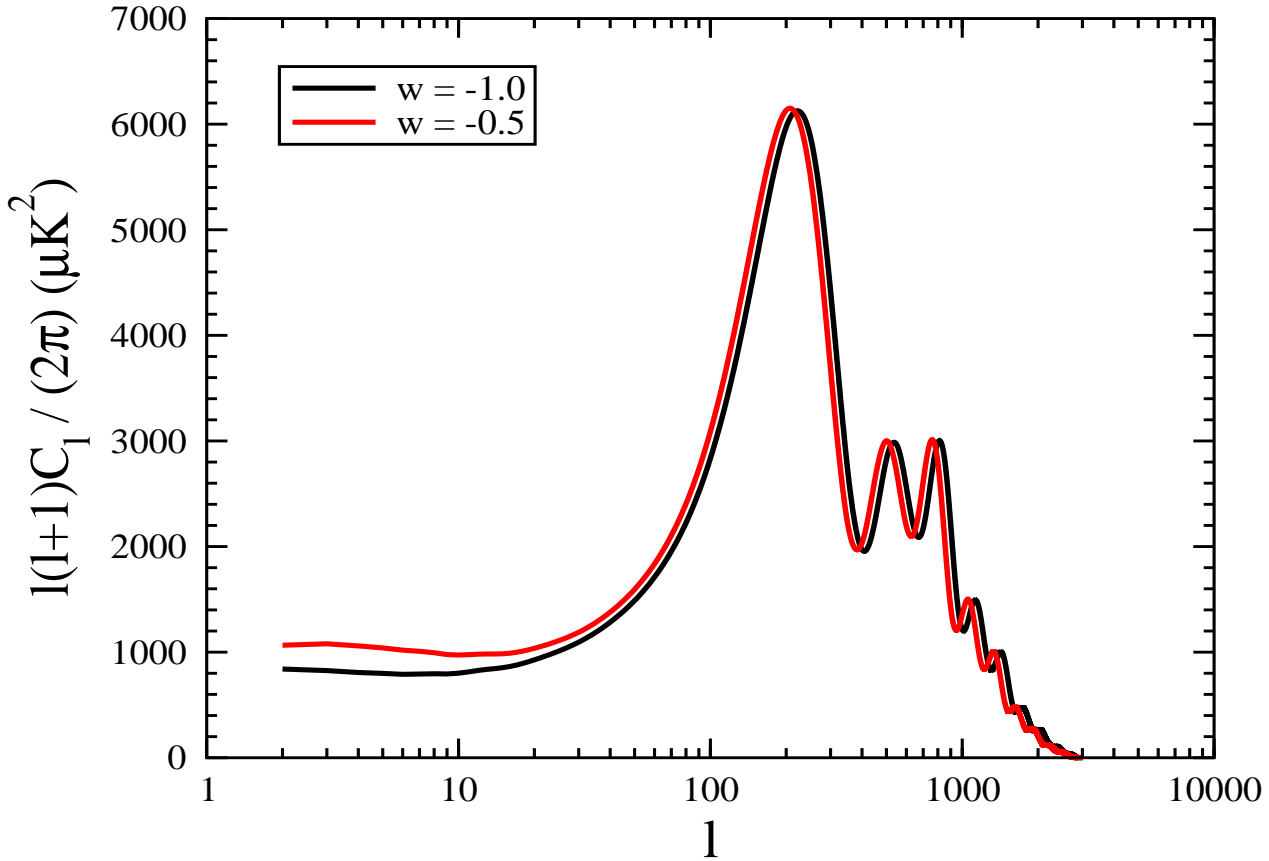


(Huterer & Ma 2003; also Chae et al. 2002)

SNAP



# CMB and Dark Energy



- **peak locations** are sensitive to dark energy (but not much):

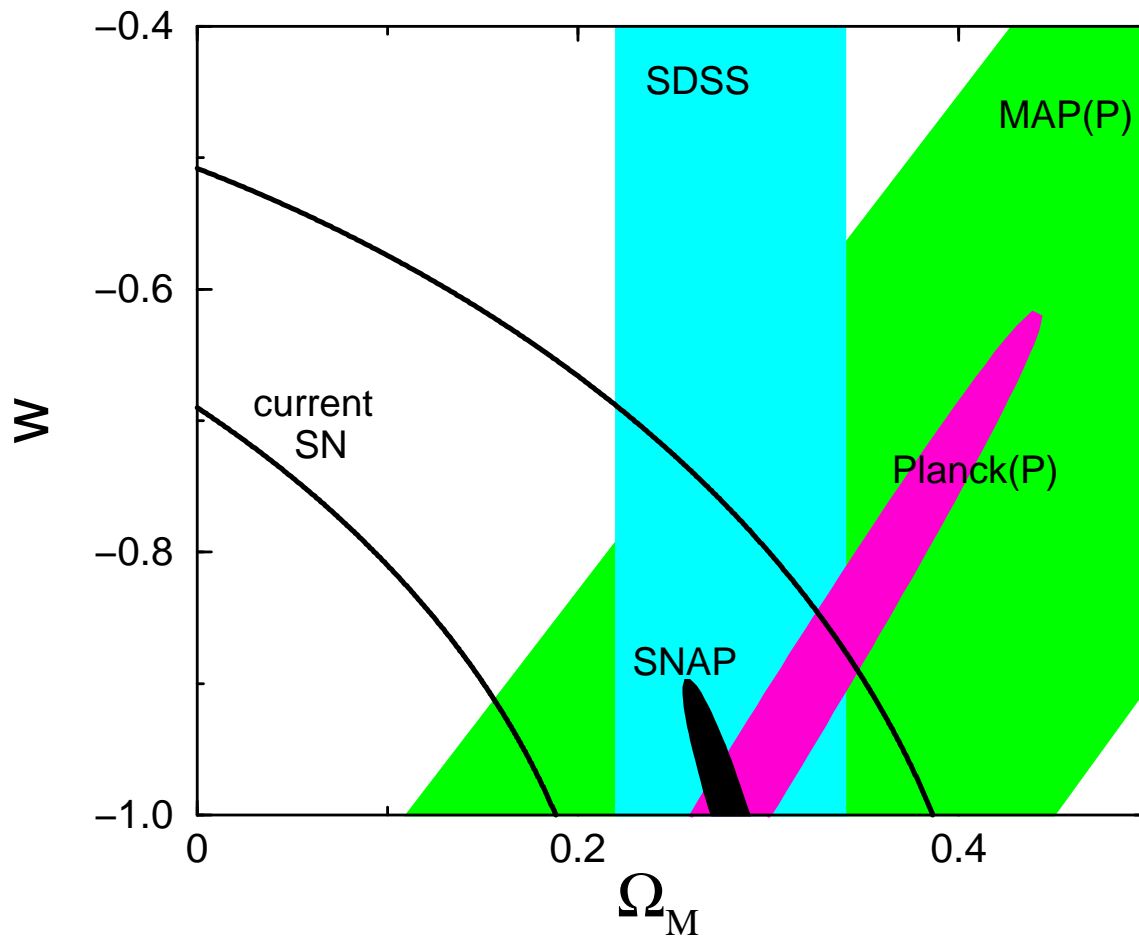
$$\frac{\Delta l_1}{l_1} = -0.084\Delta w - 0.23\frac{\Delta\Omega_M h^2}{\Omega_M h^2} + 0.09\frac{\Delta\Omega_B h^2}{\Omega_B h^2} + 0.089\frac{\Delta\Omega_M}{\Omega_M} - 1.25\frac{\Delta\Omega_{\text{TOT}}}{\Omega_{\text{TOT}}}$$

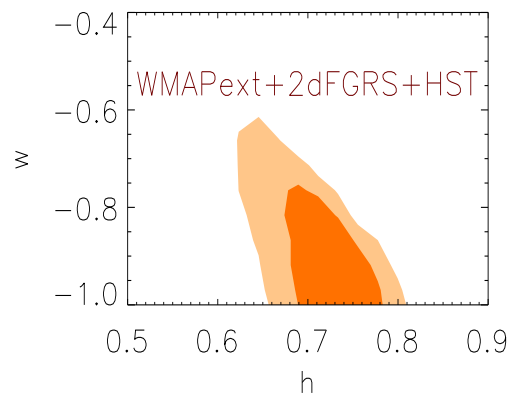
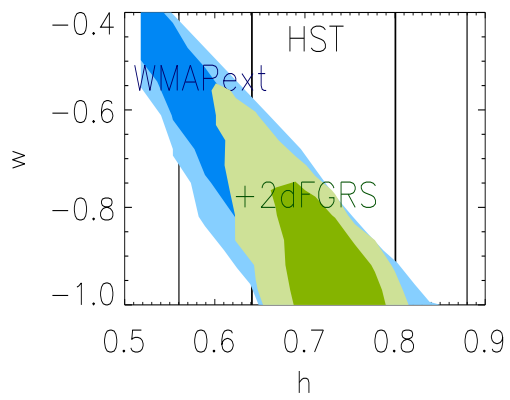
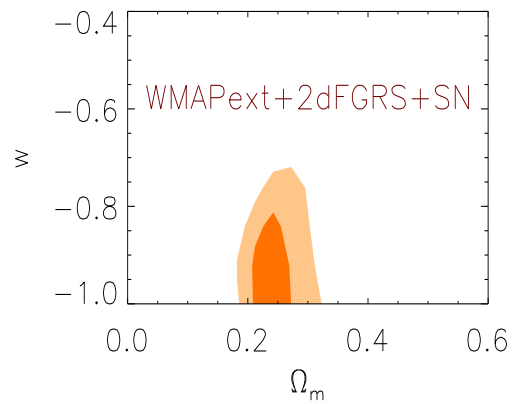
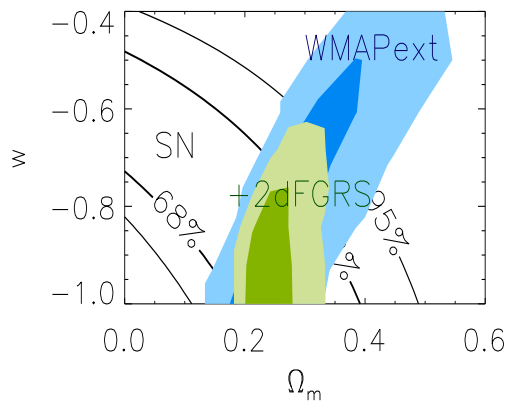
- Same as a measurement of angular diameter distance to  $z \sim 1000$   
*with  $\Omega_M h^2$  fixed*
- End up constraining:  $\mathcal{D} \equiv \Omega_M - 0.28(1 + w) \approx 0.3$   
(Planck:  $\mathcal{D}$  to  $\sim 10\%$ ) (Frieman et al. 2003)



CMB provides a *single* measurement of the distance to LSS, therefore

- Degeneracy in parameter estimation
- Only  $w_{\text{eff}}$  is probed

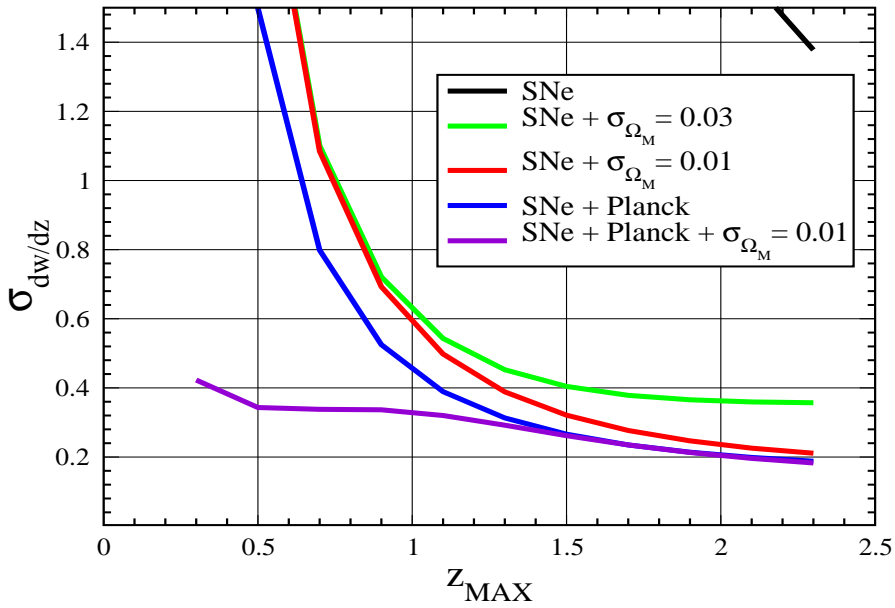
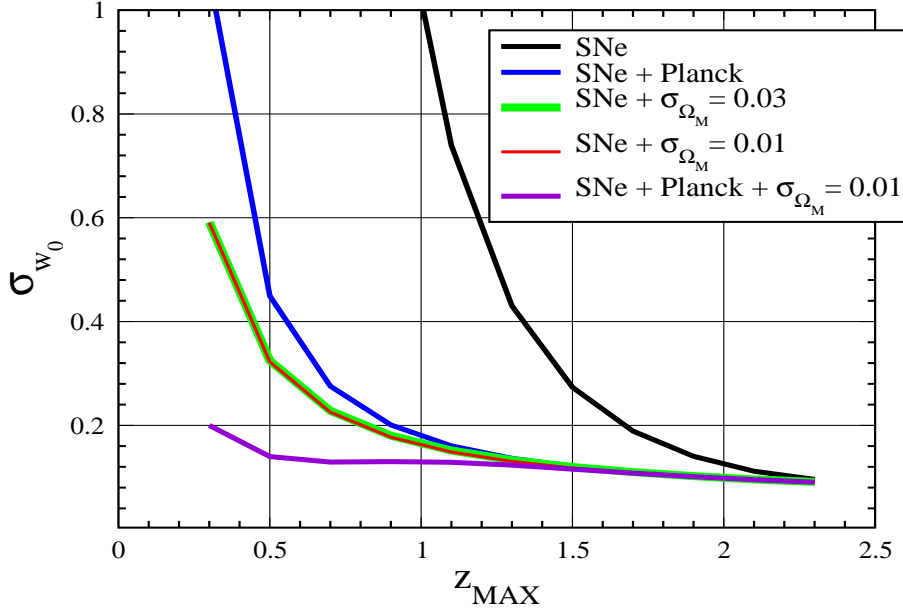




(Spergel et al. 2003)

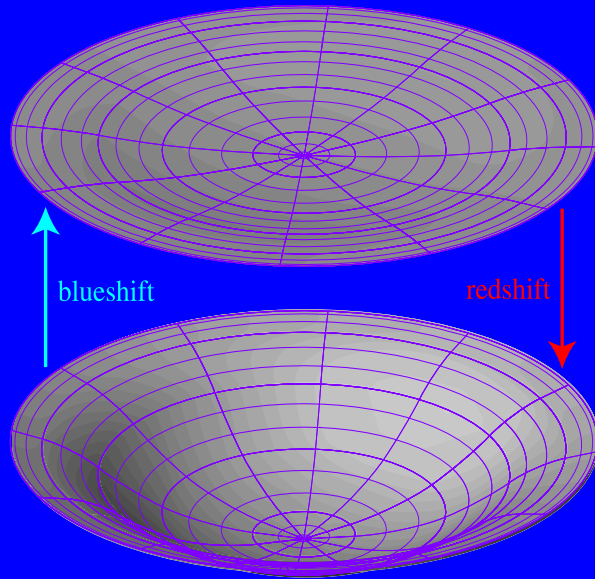
Assuming  $w(z) = w_0 + w_1 z$ :

Frieman et al. 2003



# Integrated Sachs–Wolfe Effect

- Potential redshift:  $g_{00} = -(1+\Psi)^2 \delta_{ij}$



Kofman & Starobinskii (1985)

Hu & Sugiyama (1994)

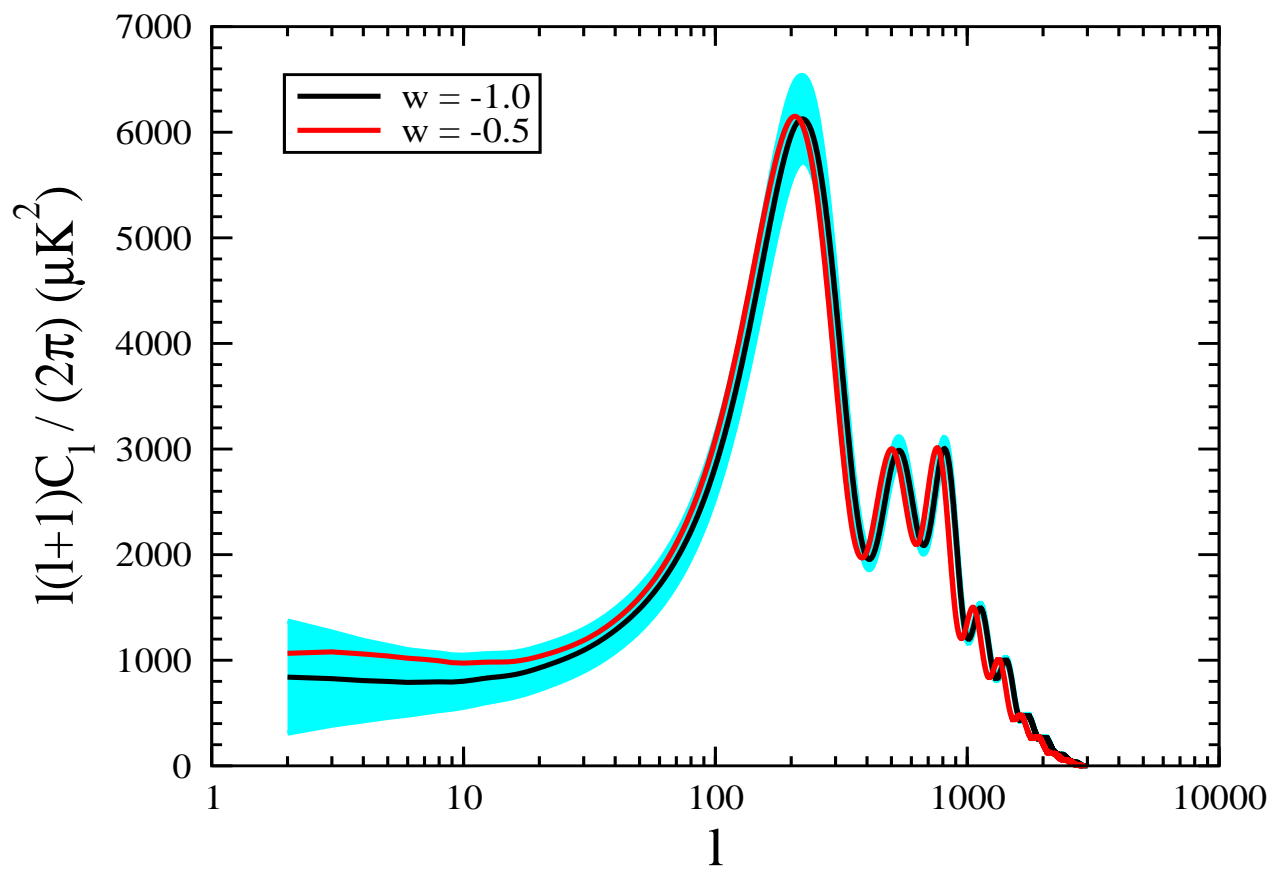
$$\Delta T^{\text{ISW}}(\hat{\mathbf{n}}) = -2 \int_0^{\eta_{\text{rec}}} d\eta' \frac{d\Phi(\eta')}{d\eta'}$$

Recall, Poisson eq.  $\nabla^2 \Phi = 3/2 H_0^2 \Omega_M (\delta/\mathbf{a})$

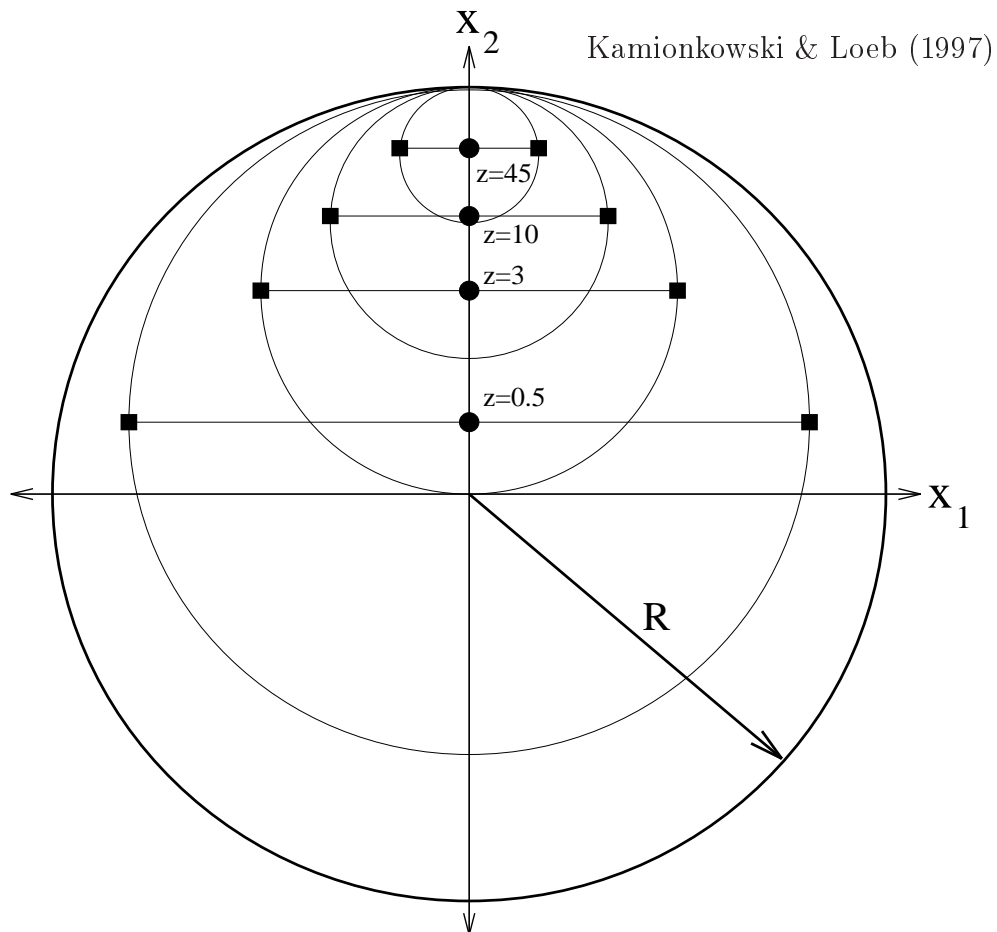
$$C_2^{\text{SW}} = \frac{4\pi}{9} \int_0^\infty \frac{dk}{k} \Delta_{\Phi\Phi}^2(k, r_{\text{rec}}) j_2^2[kr_{\text{rec}}]$$

$$C_2^{\text{ISW}} = 16\pi \int_0^\infty \frac{dk}{k} \Delta_{\Phi\Phi}^2(k, r_{\text{rec}}) \times \left[ \int_0^{r_{\text{rec}}} dr' \frac{1}{g(z_{\text{rec}})} \frac{d}{dr'} g(z') j_2(kr') \right]^2$$

But cosmic variance rules...



# Getting Around Cosmic Variance!

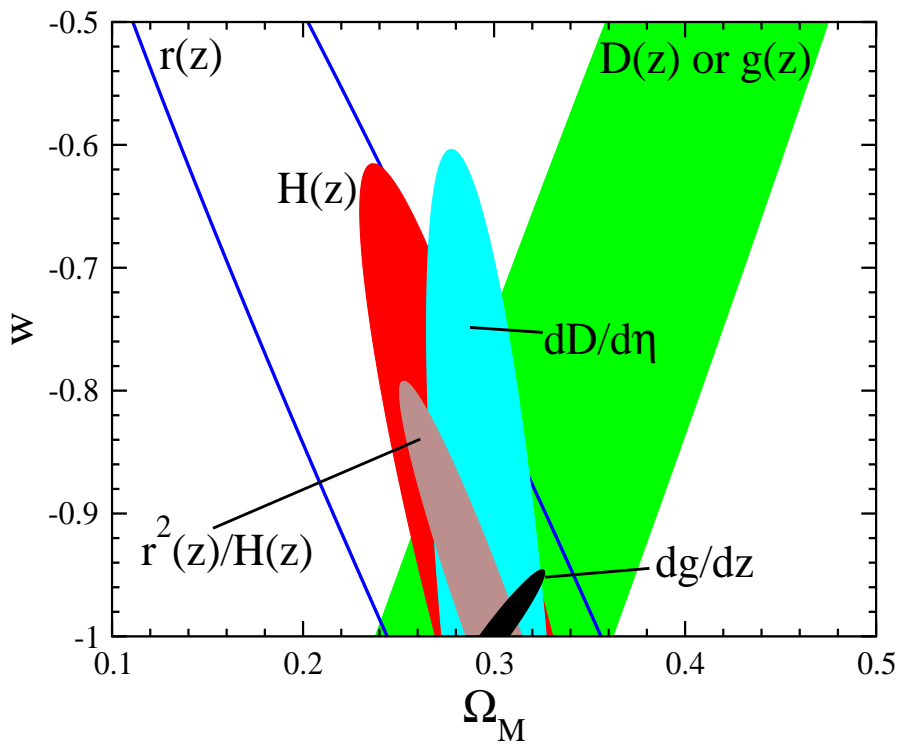


# Growth Rate of fluctuations

$$\delta(z) = \delta(0) D(z)$$

$$\ddot{\delta} + 2(\dot{a}/a)\dot{\delta} - 4\pi G\rho_M\delta = 0$$

$$g(a) \equiv D(a)/a$$



(Cooray, Huterer & Baumann 2003)

# Polarization measurements from clusters

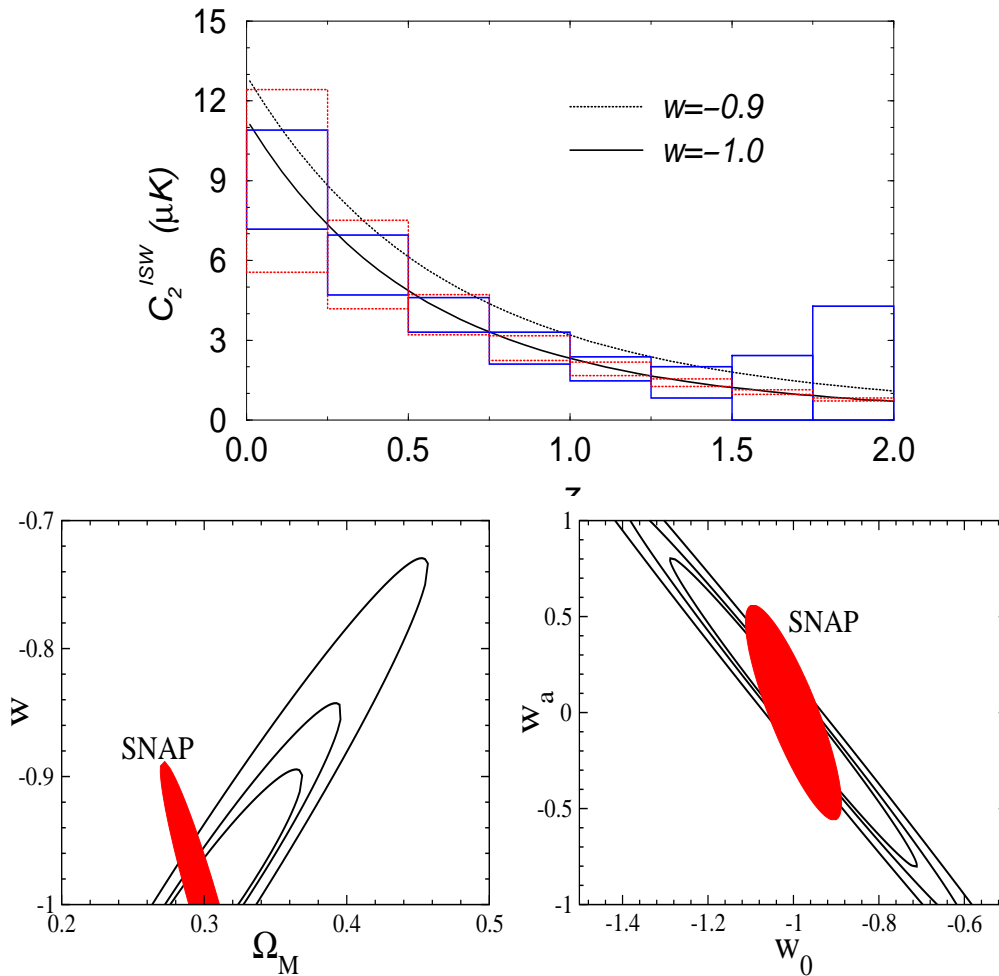
$$P^{\text{prim}} = \frac{\sqrt{6}}{10} \langle \tau \rangle \frac{Q^{\text{rms}}(z)}{T_{\text{CMB}}}$$

$$P^{\text{kin}} = \frac{1}{10} g(x) \langle \tau \rangle \langle \beta^2 \rangle$$

(Sazonov & Sunyaev 1999, Challinor et al. 2000,  
Cooray & Baumann 2003)

- Maximum signal:  $0.1(\tau/0.02) \mu\text{K}$ .
- Other effects:  $\propto \frac{kT_e}{m_e c^2} \tau^2, \beta \tau^2$
- Frequency separation possible;  $P^{\text{prim}}$  dominates





Most optimistic ellipses assume:

- 3 arcmin resolution  $\rightarrow \sim 2$  meter dish
- a few  $\mu K \sqrt{\text{sec}}$  resolution per detector with  $\sim 100$  detectors

so that  $\sim 10,000$  sq. deg. can be covered in about 1/2 year at  $\sim 0.1 \mu K$  sensitivity.

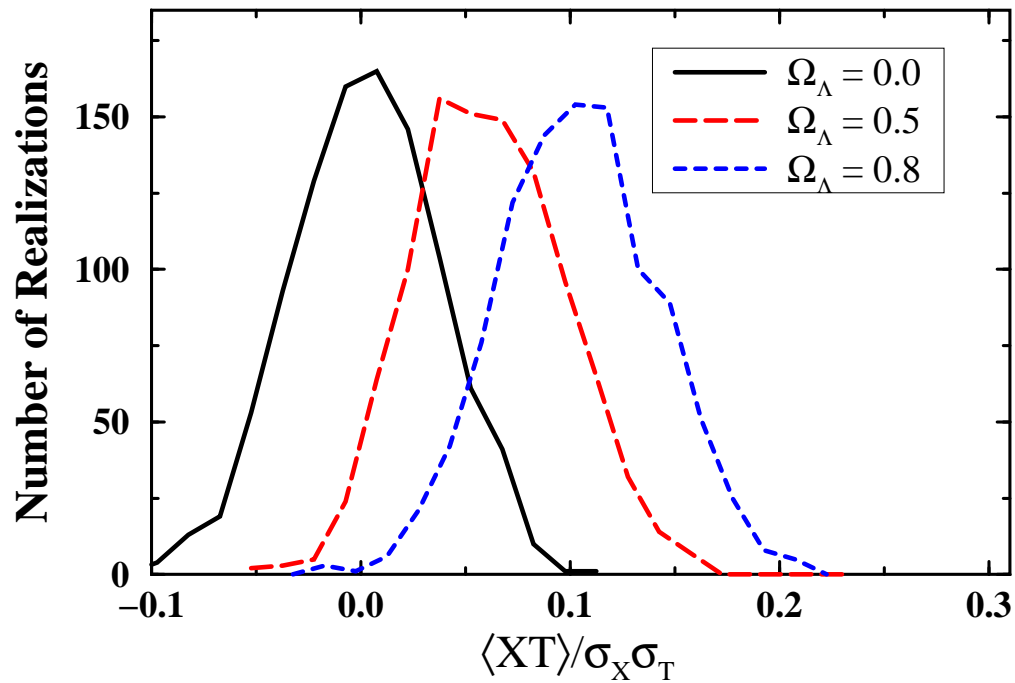
or, in several years of running, a few tens of  $nK$  sensitivity.

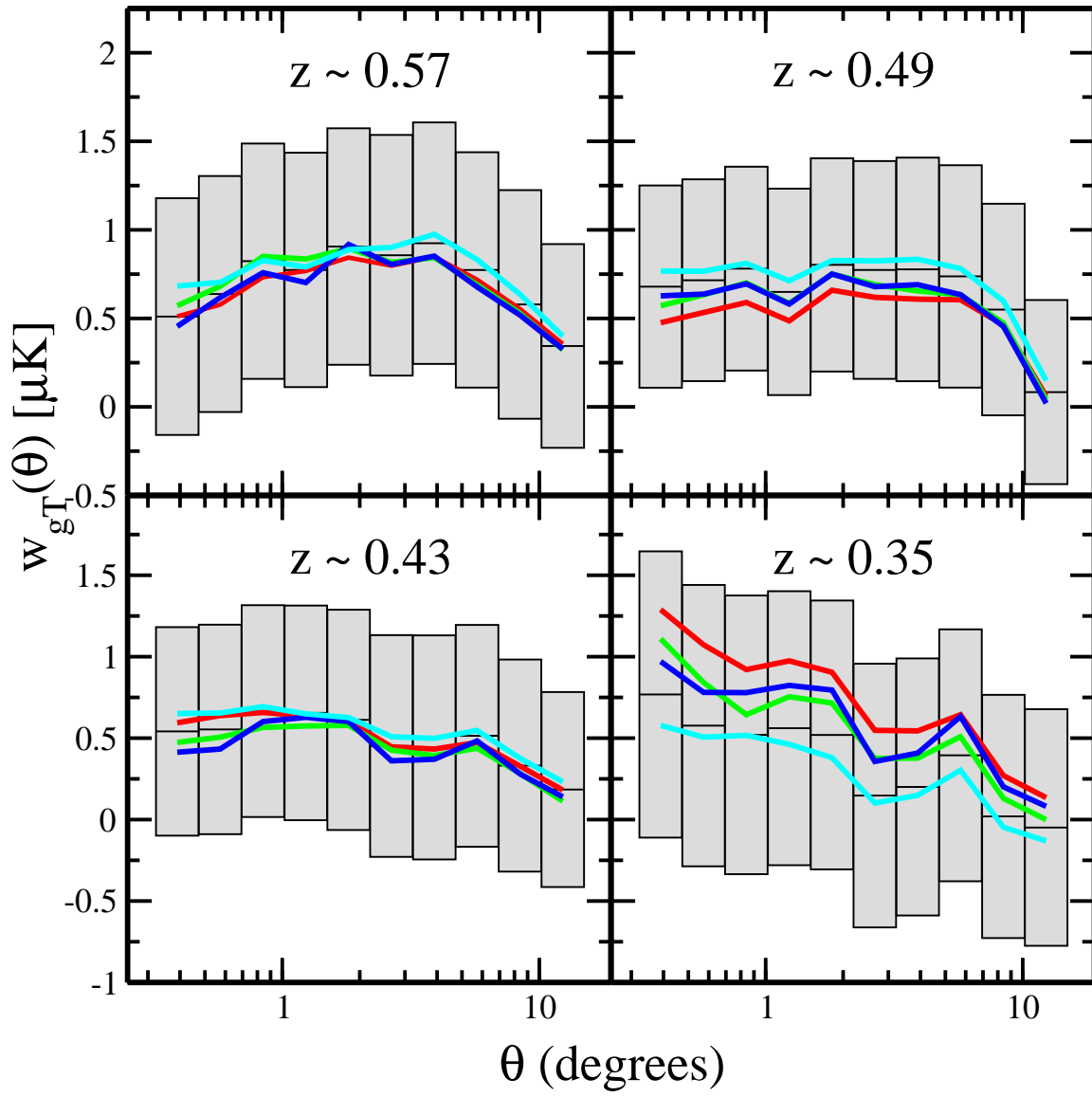
# CMB-LSS Cross-Correlation

$$\langle TX(\theta) \rangle = \frac{\sum_{\theta_{ij}=\theta} X_i T_j w_i w_j}{\sum_{\theta_{ij}=\theta} w_i w_j}$$

Non-zero correlation would be a signature of ISW  
and, therefore, dark energy!

Boughn, Crittenden, Turok et al. (1997)





Will it work even better with  $C_2(z)$ ?

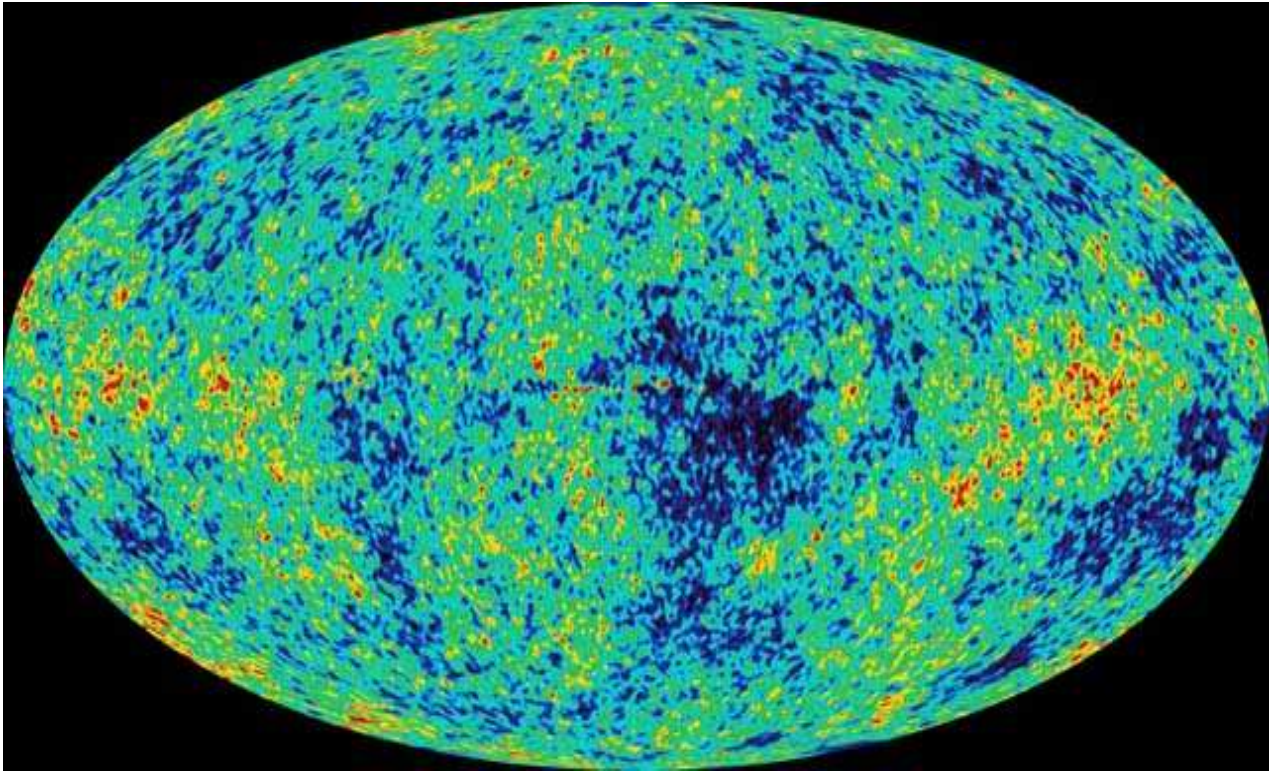
## Polarization from clusters recap:

- Observationally challenging.
- Provides  $C_2(z)$
- Which leads to  $\frac{dg}{dz}(z)$
- Decreases cosmic variance on  $C_2$
- Probes different scales of  $P(k)$

# “Multipole Vectors”!

(Copi, Huterer & Starkman 2003)

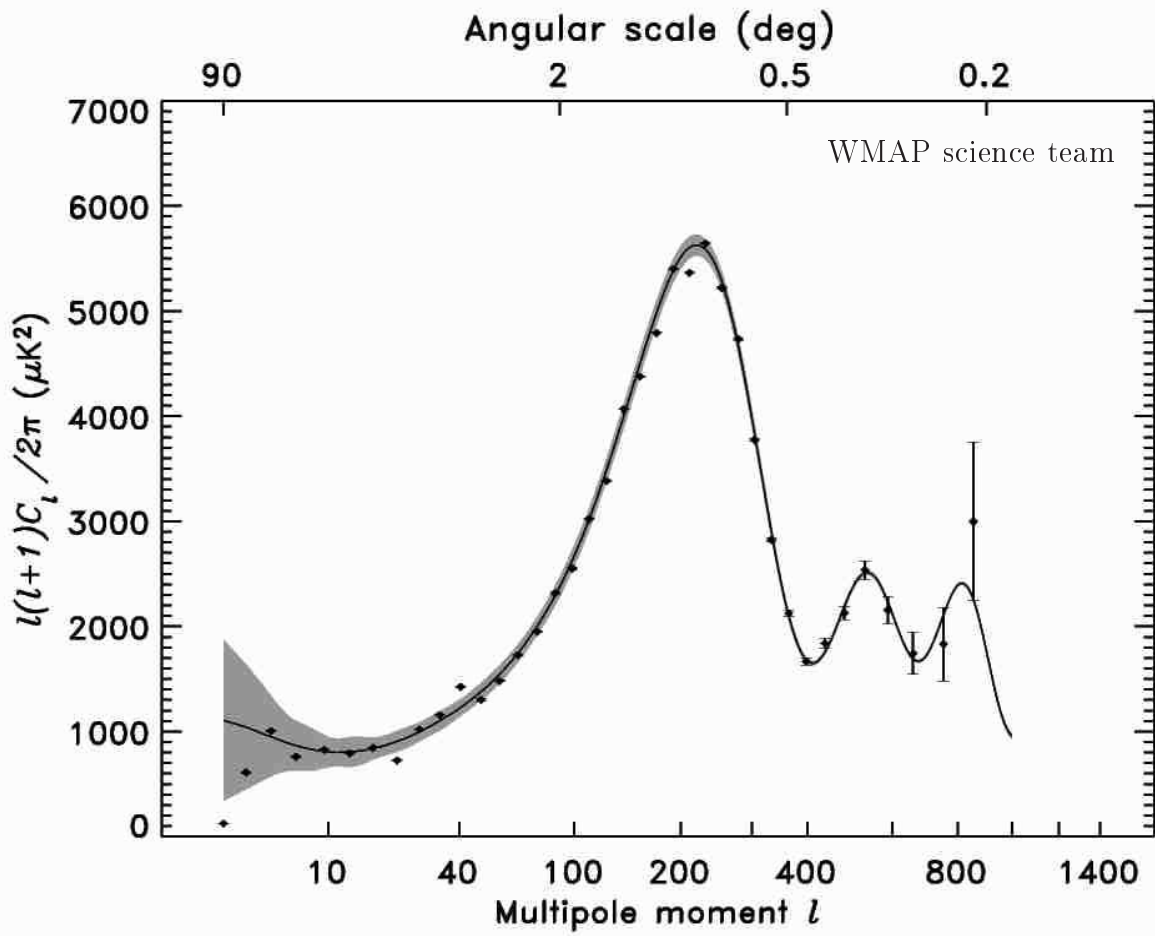
WMAP science team



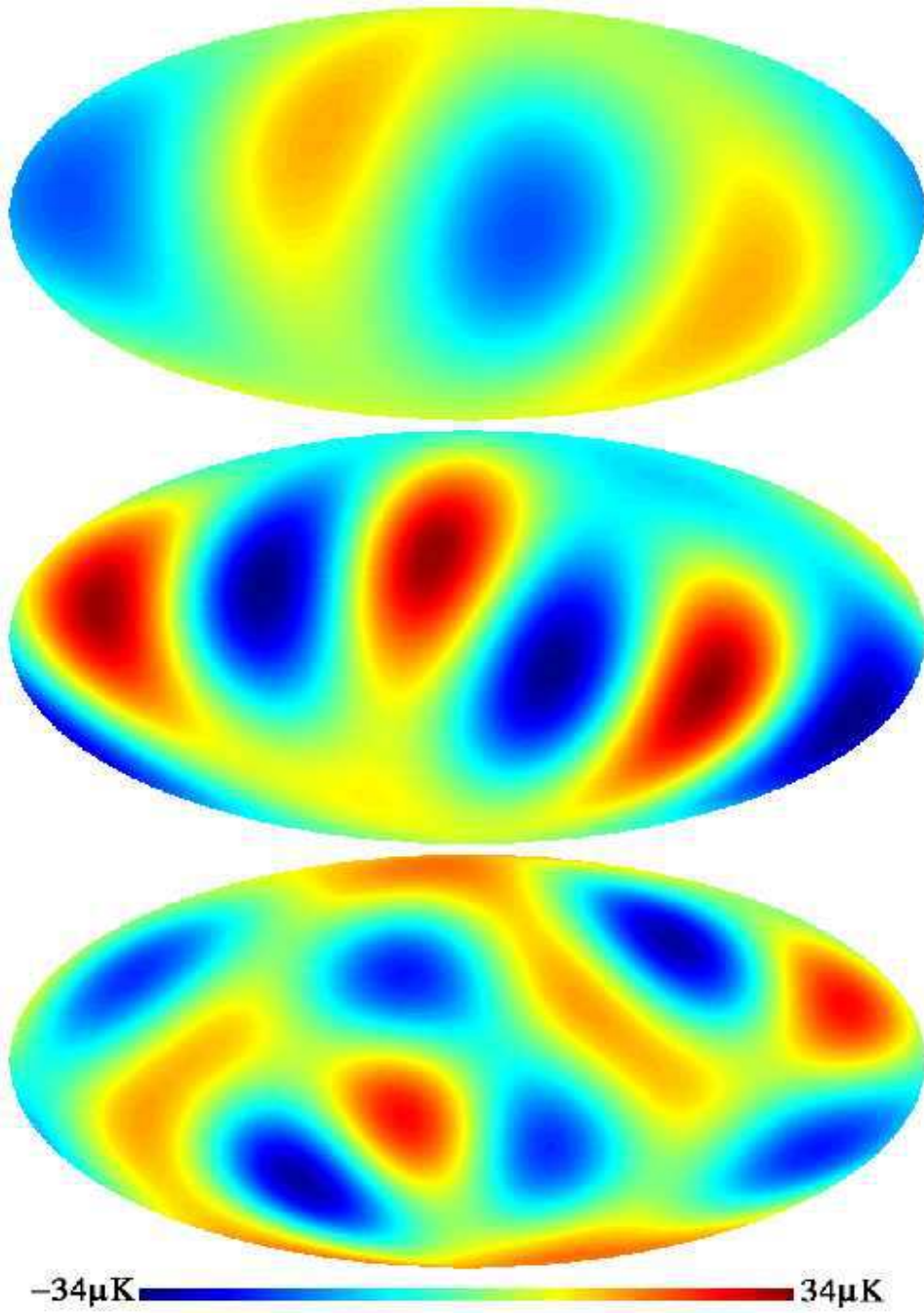
$$\frac{\delta T}{T}(\theta, \phi) = \sum_{l,m} a_{lm} Y_{lm}(\theta, \phi)$$

$$a_{lm} = \int \frac{\delta T}{T}(\Omega) Y_{lm}^*(\Omega) d\Omega$$

$$C_l = \frac{1}{2l+1} \sum_m |a_{lm}|^2$$



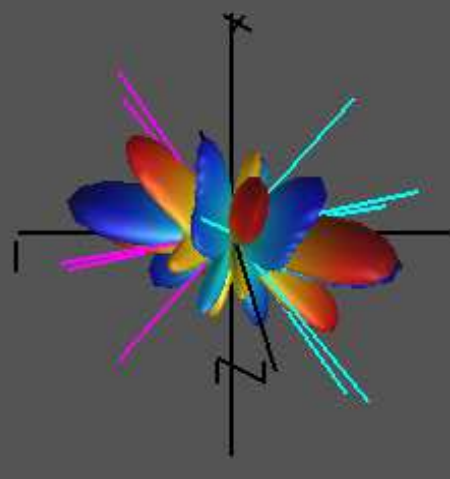
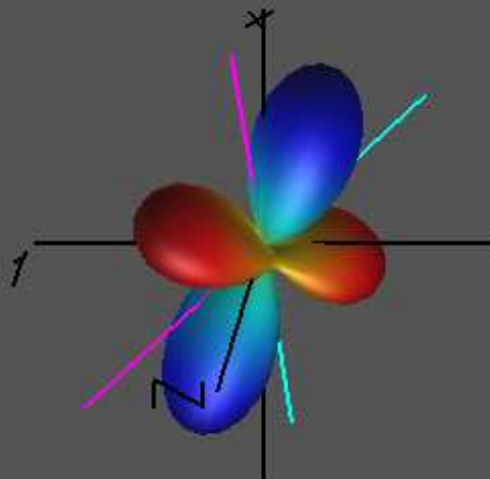
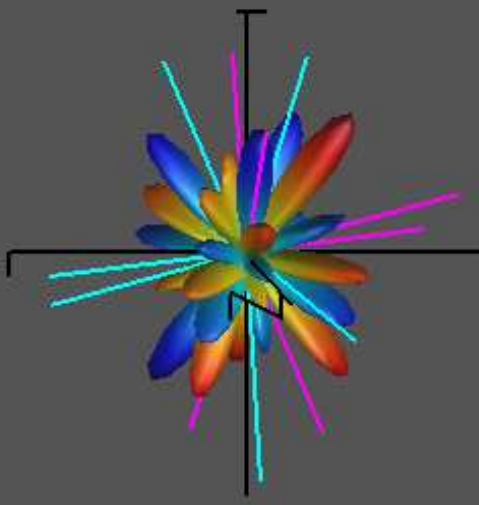
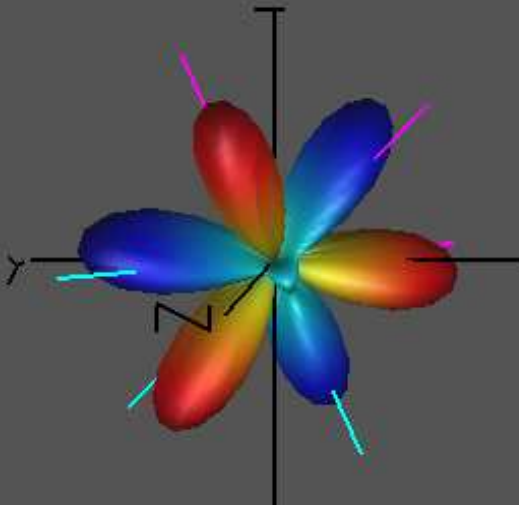
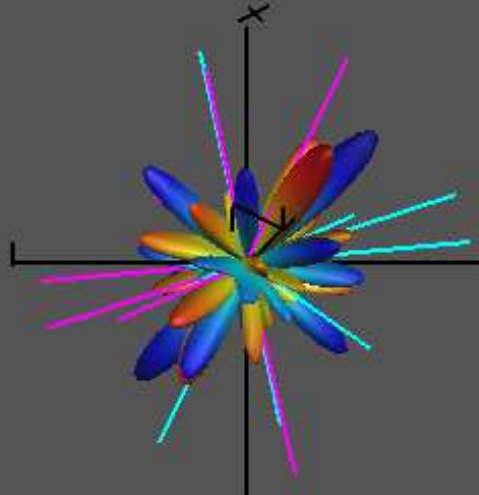
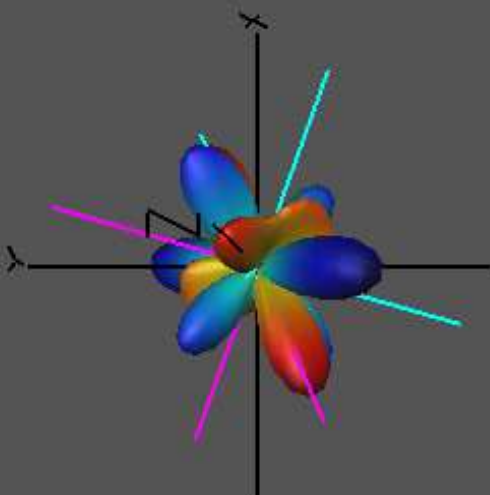
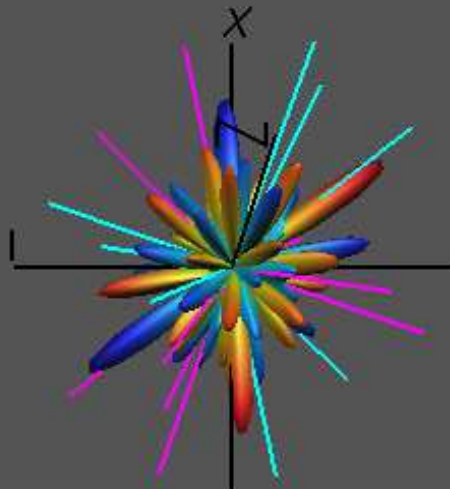
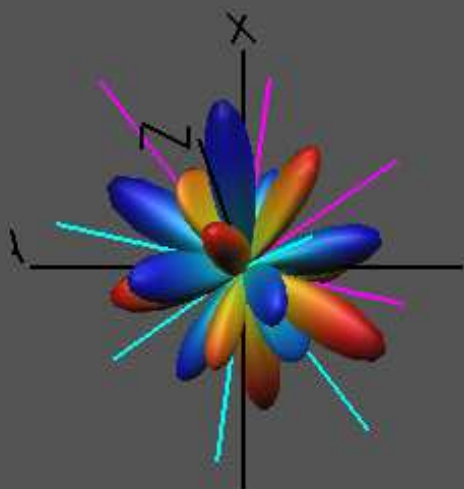
Tegmark et al. 2003



$$\begin{aligned}
\frac{\delta T}{T}(\theta, \phi) &= \sum_{l,m} a_{lm} Y_{lm}(\theta, \phi) \\
&\equiv \sum_l \left\{ a_{l,0} Y_{l,0} + 2 \sum_{m=1}^l [a_{l,m}^{\text{re}} Y_{l,m}^{\text{re}} - a_{l,m}^{\text{im}} Y_{l,m}^{\text{im}}] \right\} \\
&= \sum_l \mathbf{A}^{(l)} \left[ \hat{\mathbf{v}}_1^{(l)} \otimes \hat{\mathbf{v}}_2^{(l)} \otimes \dots \otimes \hat{\mathbf{v}}_l^{(l)} \right]
\end{aligned}$$

To compute the vectors  $\hat{v}_i^{(l)}$ :  
solve  $3l$  coupled  $l^{\text{th}}$  order equations?!  
Thankfully, no...





# Tests of Non-Gaussianity with Multipole Vectors

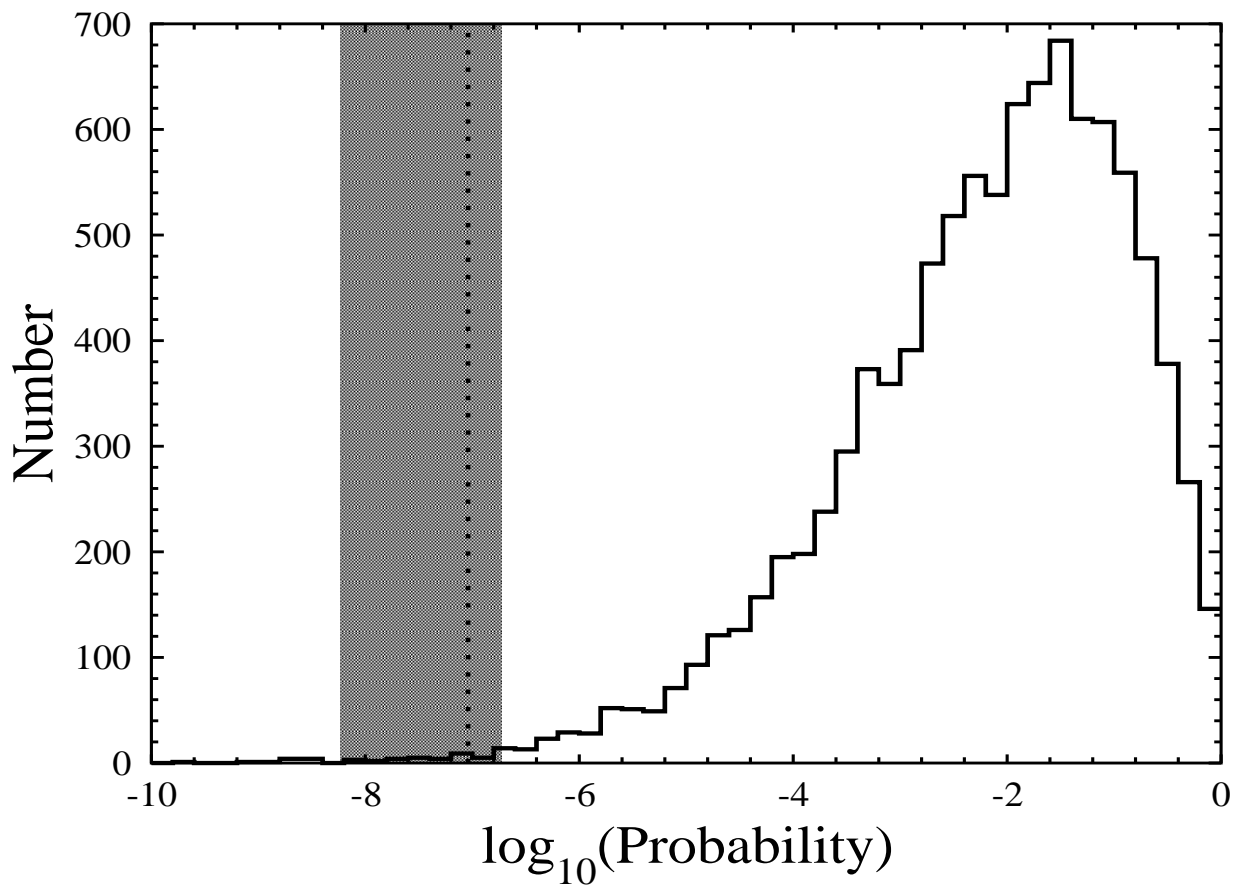
(PRELIMINARY!)

1. Bipolarity (for any multipole  $l_i$ )

$$T = \begin{pmatrix} \sum x_i^2 & \sum x_i y_i & \sum x_i z_i \\ \sum x_i y_i & \sum y_i^2 & \sum y_i z_i \\ \sum x_i z_i & \sum y_i z_i & \sum z_i^2 \end{pmatrix}$$

2. Dot products (for any  $l_i, l_j$ )

3. Dot products of cross products (for any  $l_i, l_j$ )



Probability of dot product of cross products for  $[l_i, l_j] \in \{2, 3, \dots, 8\}$ :

4 parts in a thousand!

But: goes away when quadrupole, octupole are omitted.

