

# **Strong Lensing and the Density Profiles of Halos and Galaxies**

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## Is Cold Dark Matter in trouble?

“There does not seem to be a coherent pattern to the present list of challenges to the CDM model”

(Peebles & Ratra, astro-ph/0207347).

Density profiles of galaxies are shallower than predicted;

few dwarf satellite galaxies observed than predicted

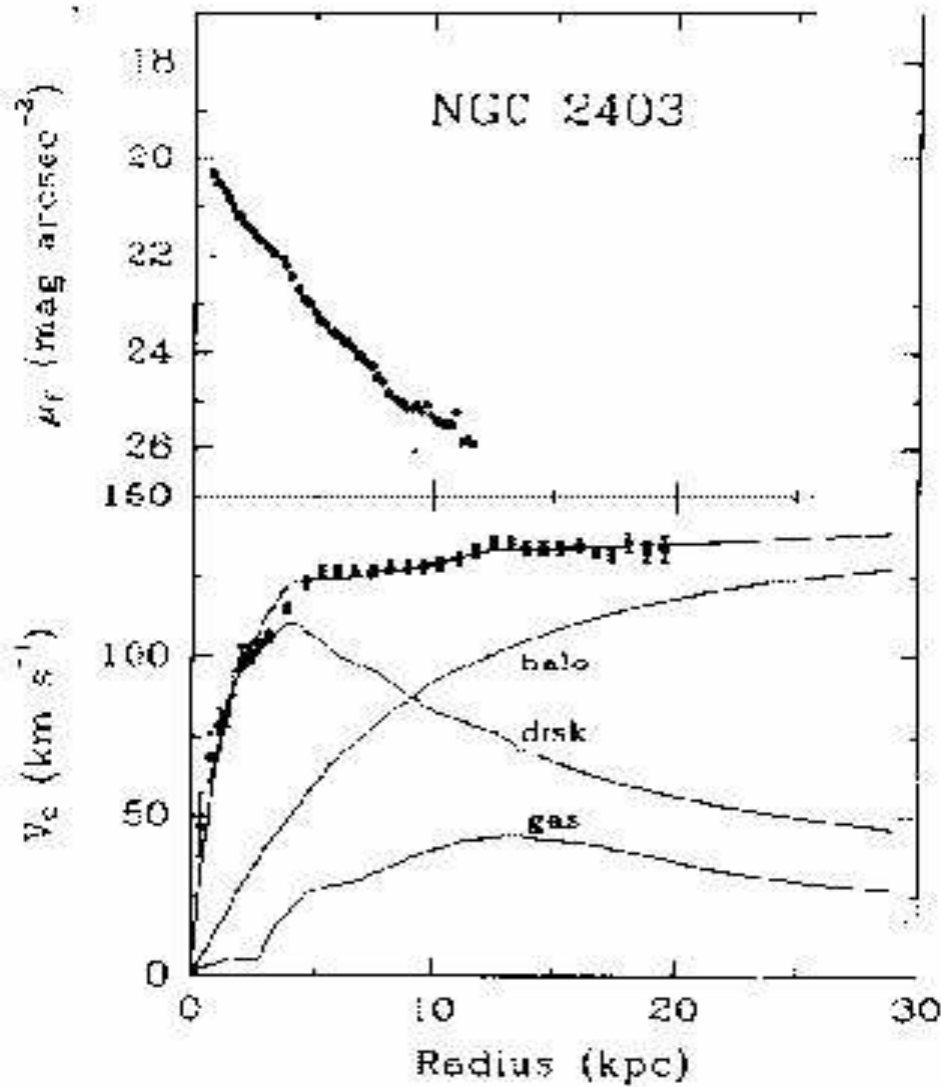
⇒ **decrease** the small-scale power

But

Quasars (powered by black holes?) observed out to  $z \approx 6$

⇒ **increase** the small-scale power

# Galaxy Rotation curves, circa 1975



$$\frac{mv^2}{r} = \frac{GMm}{r^2}$$

Luminous  $M$  only, expect

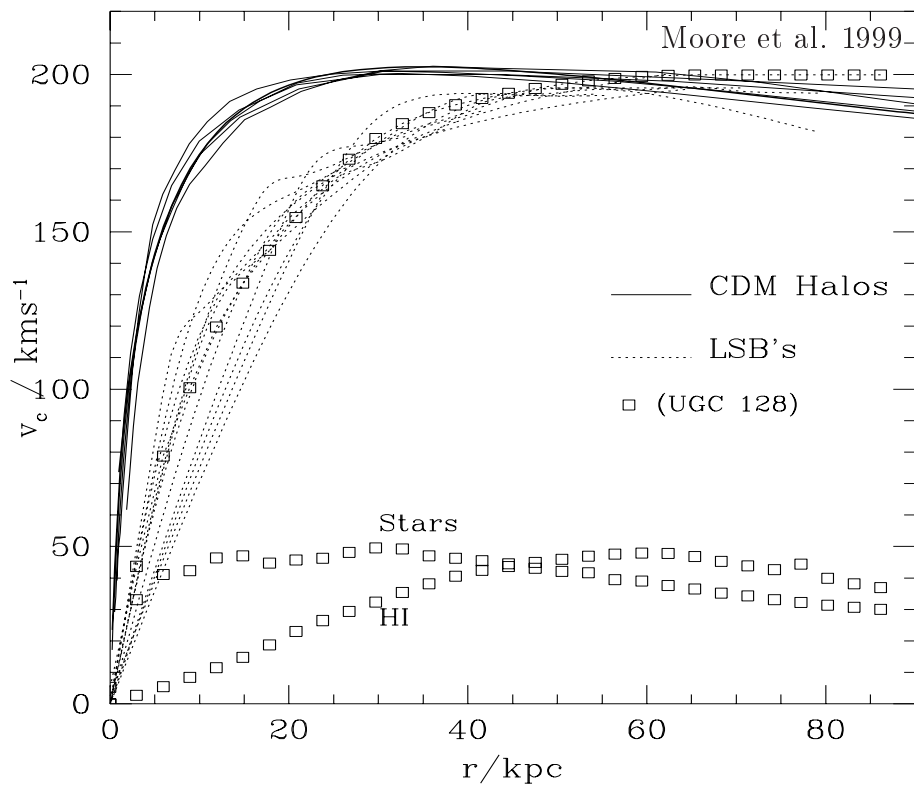
$$v \propto r^{-0.5}$$

Including the dark halo, expect

$$v^2 \propto r^2 \rho(r)$$

# Galaxy Rotation curves, circa 2000

UGC 128, de Blok & McGaugh



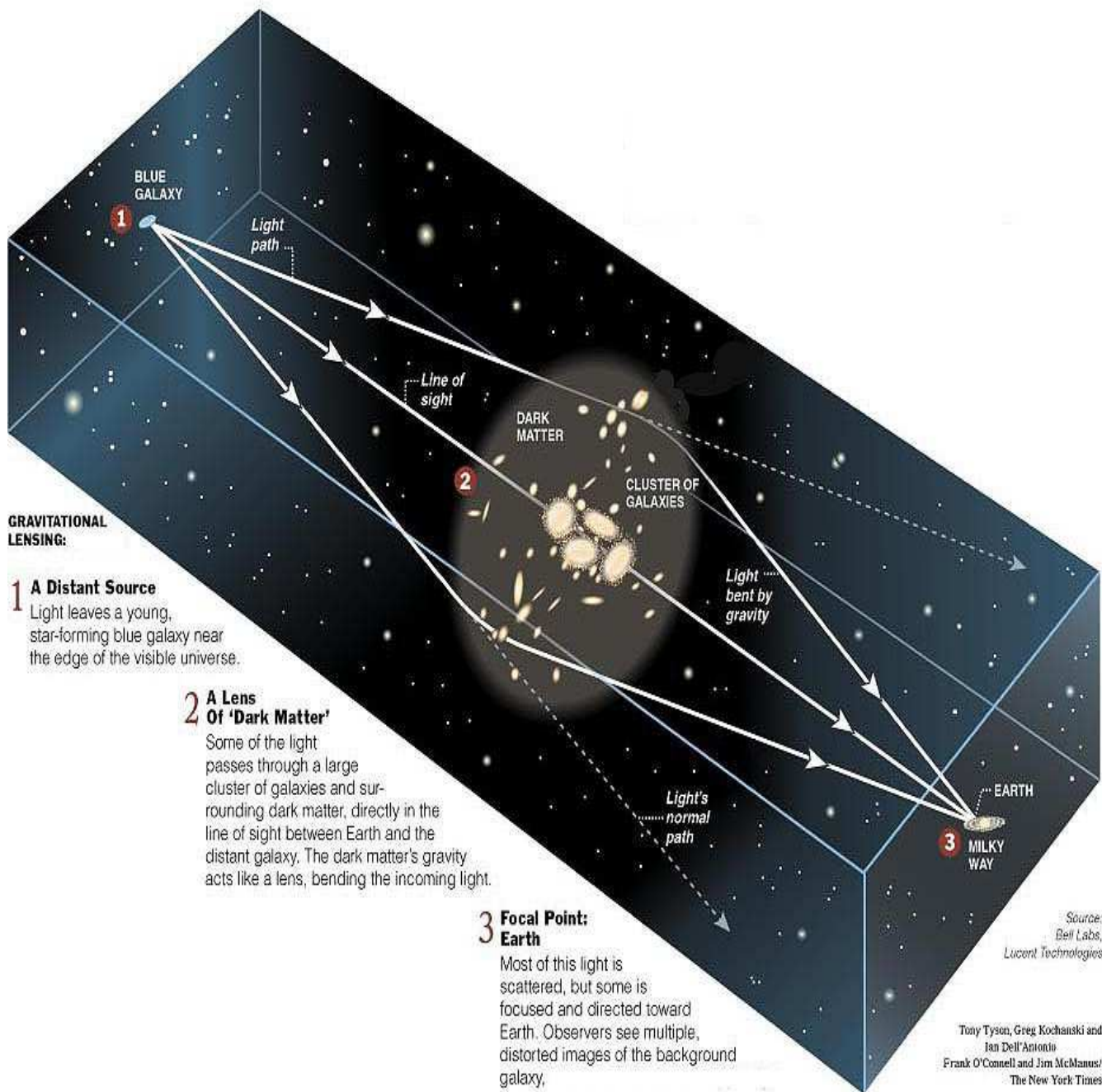
## Suggested modifications to CDM

- Feedback by supernovae, star formation, or winds
- Self-Interacting Dark Matter  
but: unnatural particle physics; makes halos that are too round; may lead to core collapse
- Primordial spectrum with less power at small scales  
but: small scale power will be regenerated in nonlinear collapse
- Warm Dark Matter  
but: may leave the Ly- $\alpha$  forest with too little power

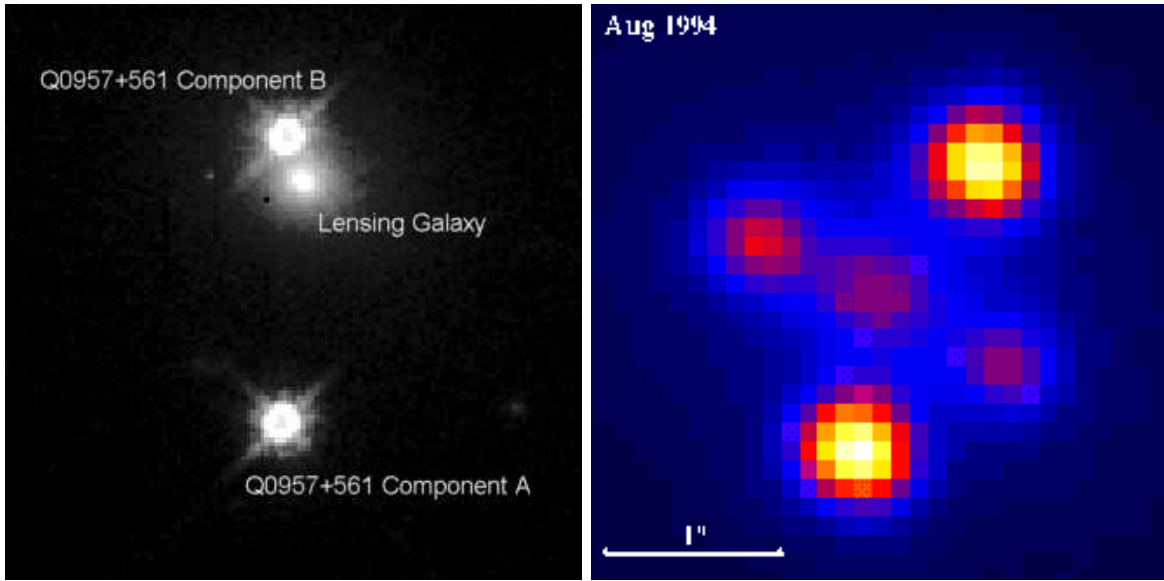
Which modification of CDM works?

Is a modification of CDM needed?

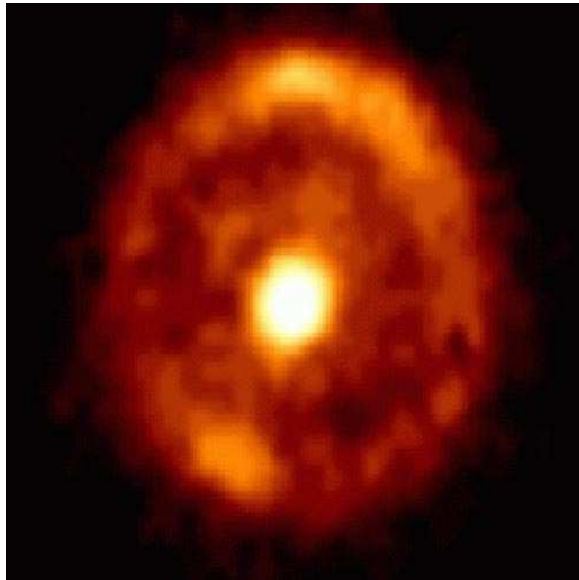
# Strong Gravitational Lensing !



Q 2237+0305 (G. Lewis and M. Irwin, WHT)



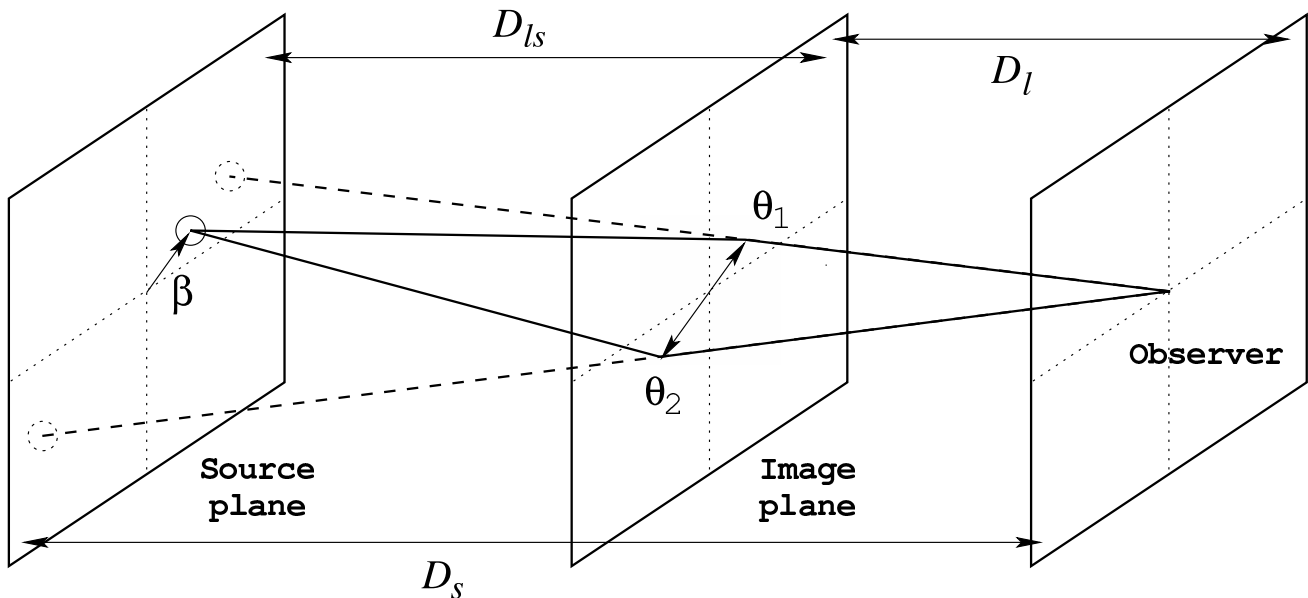
B1938+666 (L.J. King, NICMOS, HST)



# Gravitational lensing theory

From general relativity,

$$n(\mathbf{r}) = 1 - \frac{2\phi(\mathbf{r})}{c^2}$$



*The lens potential:*

$$\psi(\boldsymbol{\theta}) \equiv \frac{2}{c^2} \frac{D_{ls}}{D_l D_s} \int_{-\infty}^{+\infty} \phi(\boldsymbol{\theta}, l) dl$$

*Apply Fermat's principle:*

$$t = \frac{1+z_l}{c} \int_S^O n(\mathbf{r}) dr = \frac{1+z_l}{c} \int_S^O dr - \frac{1+z_l}{c^3} \int_S^O 2\phi(\mathbf{r}) dr$$

$$t(\boldsymbol{\theta}) = t_0 + \left(\frac{1+z_l}{c}\right) \left(\frac{D_l D_s}{D_{ls}}\right) \left[\frac{1}{2}|\boldsymbol{\theta} - \boldsymbol{\beta}|^2 - \psi(\boldsymbol{\theta})\right]$$

*The lens equation:*

$$\boldsymbol{\theta} = \boldsymbol{\beta} - \nabla\psi(\boldsymbol{\theta})$$



## Optical depth for lensing

$d\tau = (\text{density of galaxies}) \times (\text{lensing cross-section}) \times c dt$

$$\tau = \int_0^{z_s} dz_l \frac{dD_l}{dz_l} (1 + z_l)^3 \times \int_0^\infty dL \frac{d\phi}{dL}(L) \sigma_{\text{SIS}}(L, z_l, z_s) B(L, z_l, z_s)$$

Required input:

- Cosmology ( $\Omega_M, \Omega_\Lambda$ )
- Luminosity (or mass) function:

$$\frac{d\phi}{dL}(L) dL = \phi_* \left( \frac{L}{L_*} \right)^\alpha \exp(-L/L_*) \frac{dL}{L_*}, \quad \text{or}$$

$$\frac{d\phi}{d \ln M} = 0.315 \frac{\rho_M}{M} \frac{d \ln \sigma^{-1}}{d \ln M} \exp(-|\ln \sigma^{-1} + 0.61|^{3.8})$$

- Cross-section for lensing, e.g. SIS:  $\rho(r) = \sigma^2/(2\pi G r^2)$

$$\sigma_{\text{SIS}} = \pi(\theta_E D_l)^2 = 16\pi^3 \left( \frac{\sigma}{c} \right)^4 \left( \frac{D_l D_{l_s}}{D_s} \right)^2$$

- relate  $L$  to  $\sigma$ : Faber-Jackson relation

$$\frac{L}{L_*} = \left( \frac{\sigma}{\sigma_*} \right)^\gamma \quad \gamma \sim 4$$

# Strong Lensing as a Probe of Cosmology

- Many papers since the late 1980s; the idea is

$$N_{\text{lens}} \propto \text{Volume} = \frac{r^2(z)}{H(z)}$$

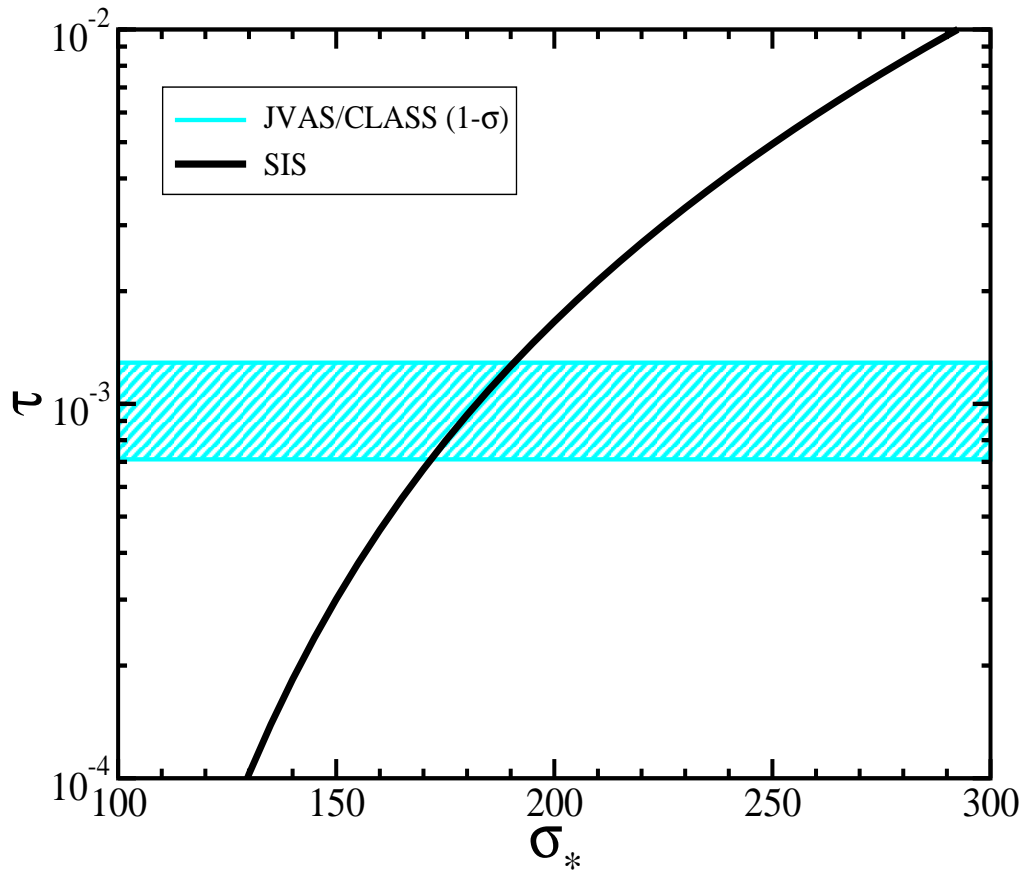
- Kochanek (1996)  $\Omega_{\Lambda} < 0.66$  at 95% CL.
- But other parameter choices may favor high  $\Omega_{\Lambda}$   
(Cheng & Krauss 1999, Chiba & Yoshii 1999)
- also Kochanek 1994, 1995, Im et al. (1997), Waga & Miceli 1999, Cooray, Quashnock & Miller 1999, Sarbu, Rusin & Ma 2001, Chae et al. 2002

but much better...

## Strong Lensing as a Probe of Lens Properties !

(Krauss, 1990s)

$$d \ln \tau = 2.07 d \ln z_s + 1.00 d \ln \phi_* + 0.69 d \ln \alpha + \\ 4.16 d \ln \sigma_* + 0.69 d \ln \gamma - 0.61 d \ln \Omega_M + 0.61 d \ln w.$$

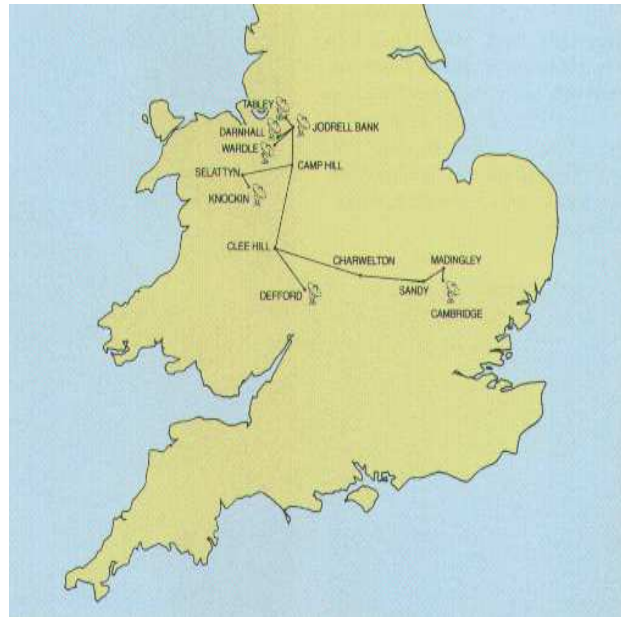


# Jodrell-Very Large Array Astrometric Survey / Cosmic Lens All-Sky Survey (JVAS/CLASS)

Very Large Array (VLA)

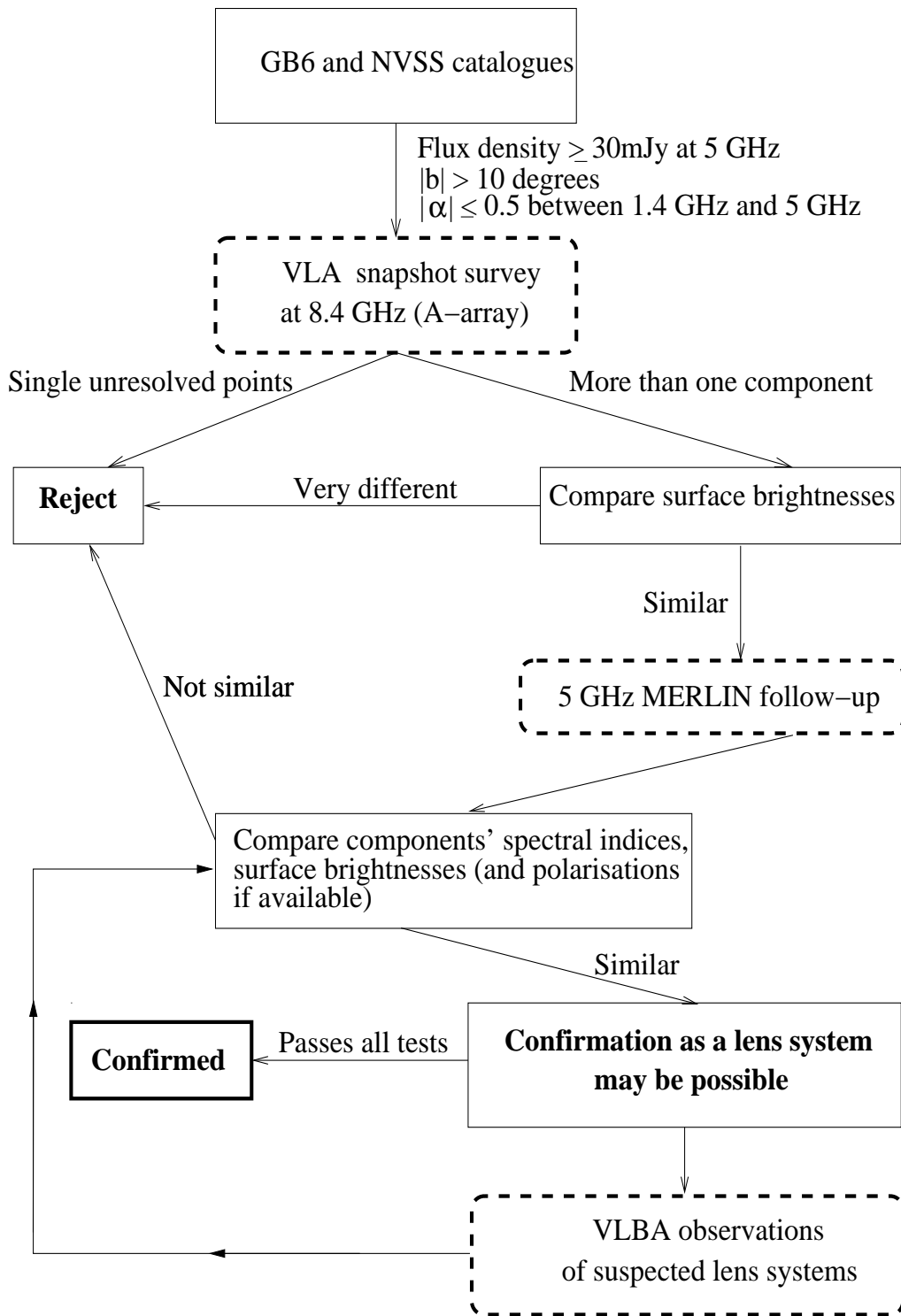


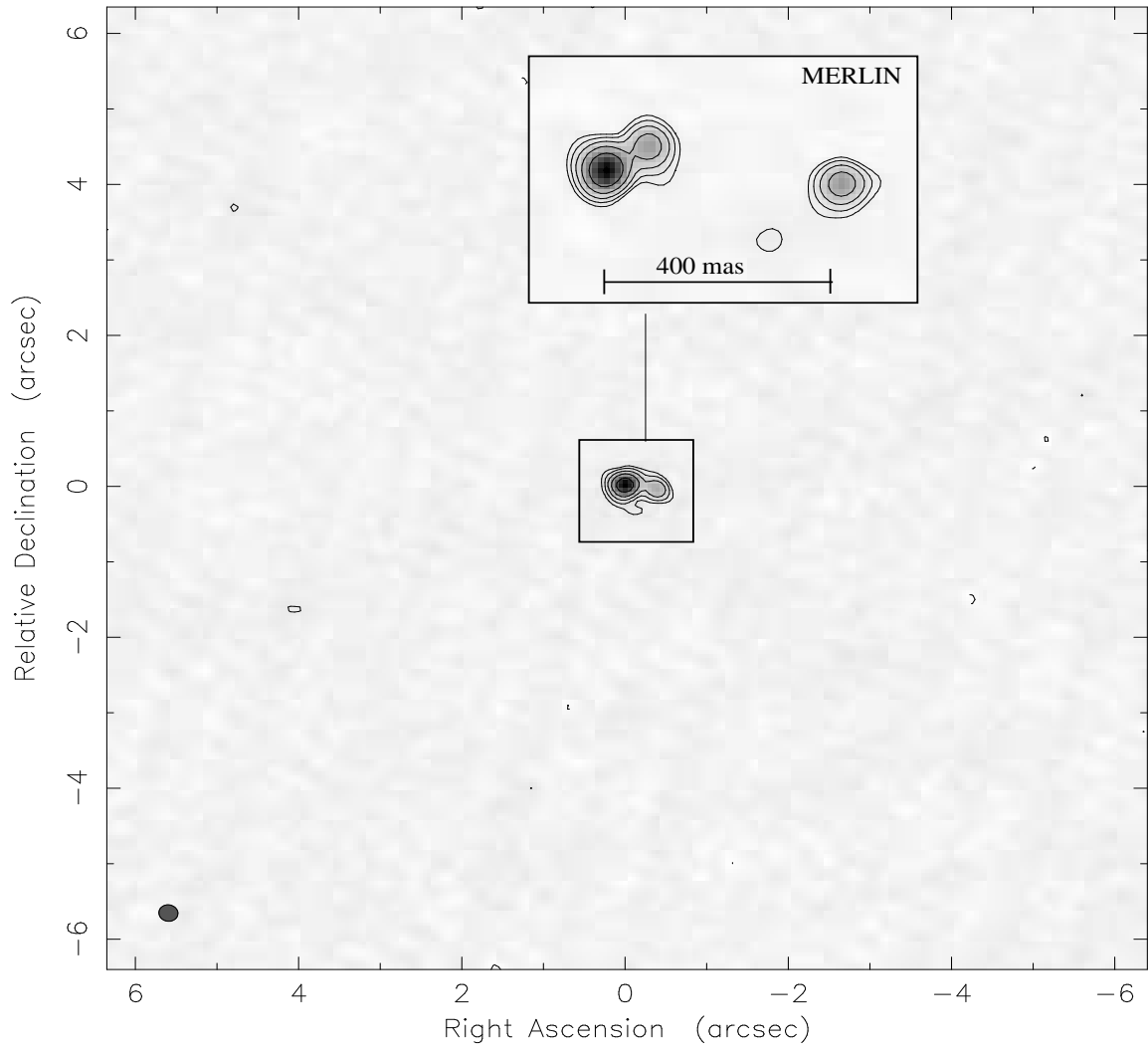
Multi-Element Radio-Linked  
Interferometer Network (MERLIN)



Very Long Baseline Array (VLBA)







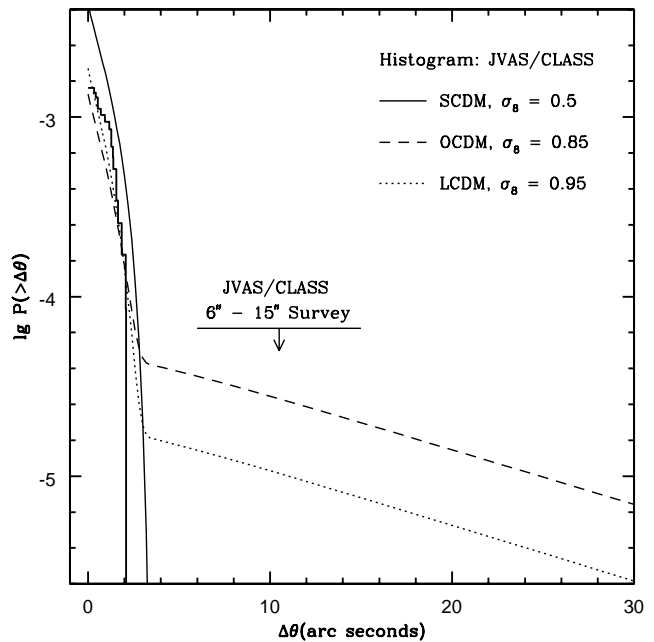
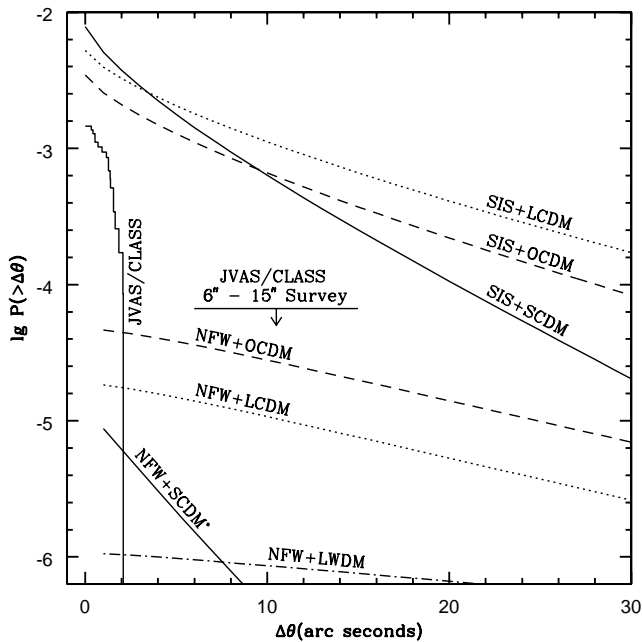
# JVAS/CLASS: some peculiar results

- All images are doubles or quads (odd image number theorem?)
- Surprisingly many spiral galaxy lenses
- No dark lenses
- No image separations larger than 3 arcsec  $\Rightarrow$

Li & Ostriker 2001: there are two populations of objects

$M \lesssim 10^{13} M_{\odot}$  (“galaxies”) with  $\rho(r) \propto r^{-2}$

$M \gtrsim 10^{13} M_{\odot}$  (“clusters”) with  $\rho(r) \propto r^{-1}$

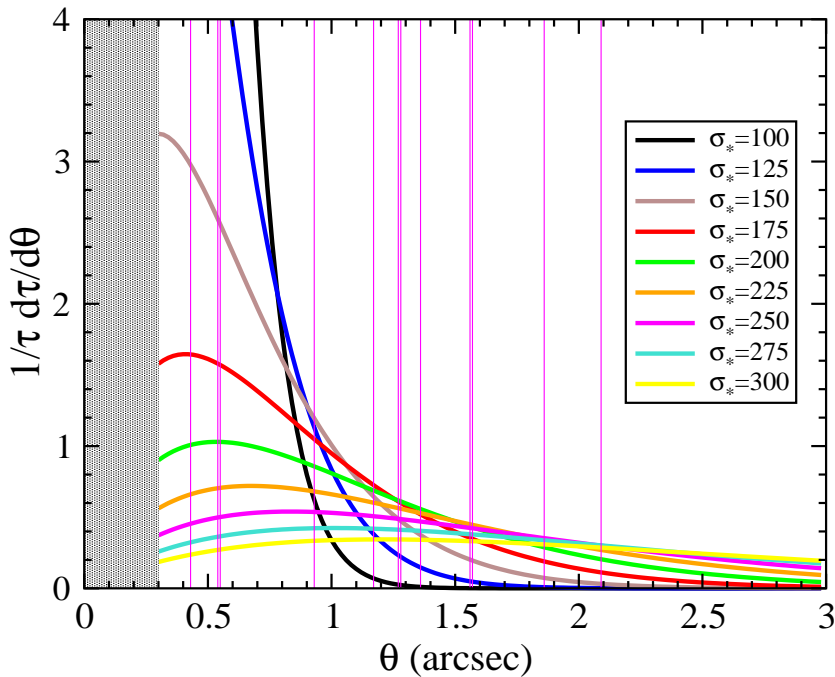


# Elliptical galaxies in JVAS/CLASS

Survey	Lens	$z_l$	$z_s$	$\theta$	Reference
JVAS/MG	J0414+054	0.96	2.64	2.09	Hewitt et al. 1992
CLASS	B0128+437	—	—	0.54	Phillips et al. 2000
CLASS	B0712+472	0.41	1.34	1.27	Jackson et al. 1998
CLASS	B0739+366	—	—	0.54	Marlow et al. 2001
JVAS	B1030+074	0.60	1.54	1.56	Xanthopoulos et al. 1998
CLASS	B1152+199	0.44	1.02	1.56	Myers et al. 1999
JVAS	B1422+231	0.34	3.62	1.28	Patnaik et al. 1992b
CLASS	B1555+375	—	—	0.43	Marlow et al. 1999
CLASS	B1933+503	0.76	—	1.17	Sykes et al. 1998
JVAS	B1938+666	0.88	—	0.93	King et al. 1997
CLASS	B2045+265	0.87	1.28	1.86	Fassnacht et al. 1999
CLASS	B2319+051	0.62	—	1.36	Rusin et al. 2001a

$$\tau = \int_0^{z_s} dz_l \frac{dD_l}{dz_l} (1 + z_l)^3 \times \int_{L_{\min}}^{\infty} dL \frac{d\phi}{dL}(L) \sigma_{\text{SIS}}(L, z_l, z_s) B(L, z_l, z_s)$$

$$\rho(r) = \frac{\sigma^2}{2\pi G r^2}$$

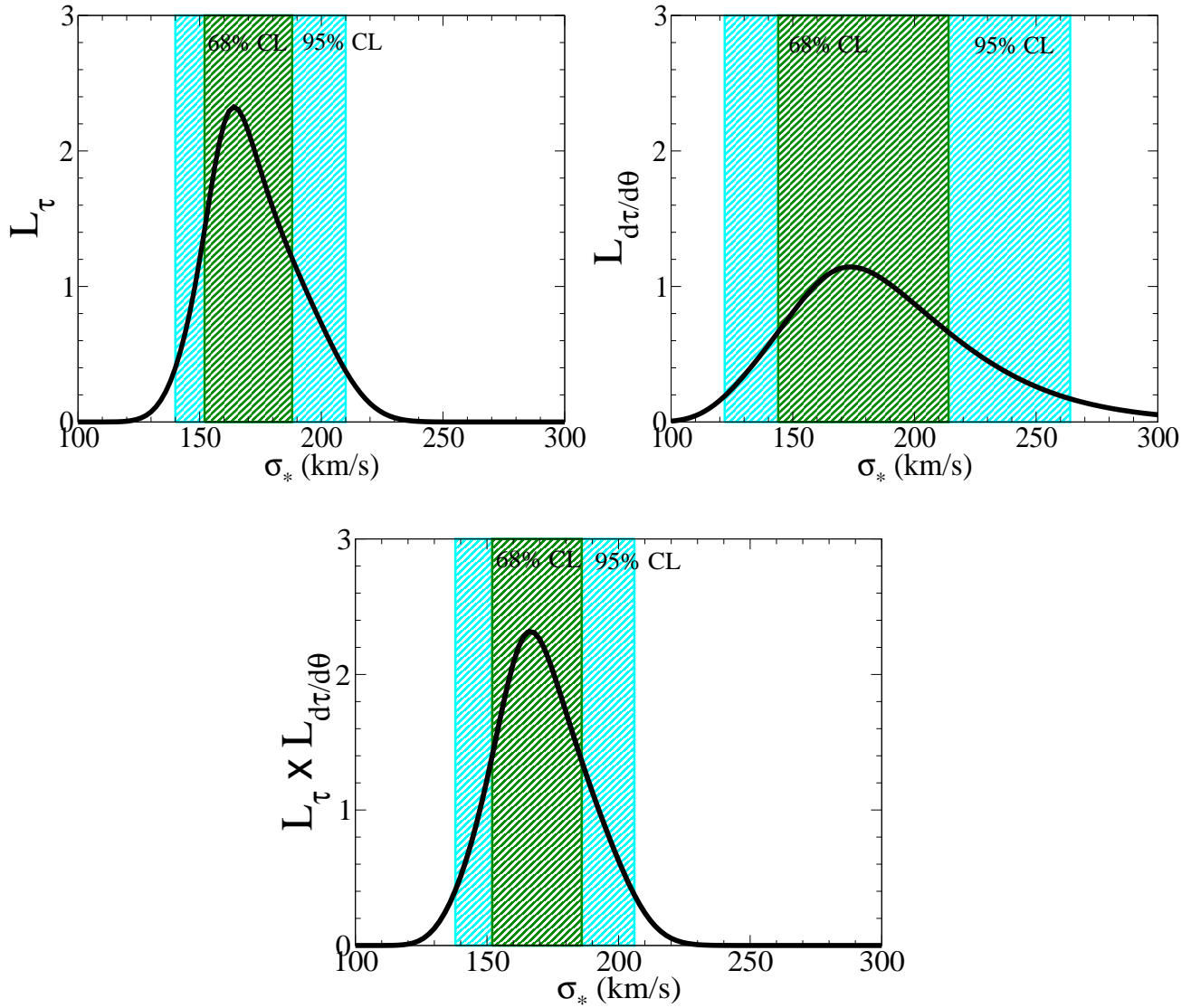


Parameters:

$\sigma_*, \gamma, \phi_*, \alpha$



# Results – SIS Profile



$138 \text{ km/s} \leq \sigma_* \leq 206 \text{ km/s}$  (95% CL)

less than the widely used 225 km/s!

# Constraining the Profile of Elliptical Galaxies

Generalized Navarro-Frenk-White profile:

$$\rho(r) = \frac{\rho_s}{\left(\frac{r}{r_s}\right)^\beta \left[1 + \left(\frac{r}{r_s}\right)\right]^{3-\beta}}$$

N-body simulations:  $\beta = 1$  (NFW 1997),  $\beta = 1.5$  (Moore et al. 1999)

Halo mass:

$$M \equiv M_{200} = 200 \left( \frac{4\pi}{3} r_{200}^3(z) \rho_c(z) \right)$$

Halo concentration:

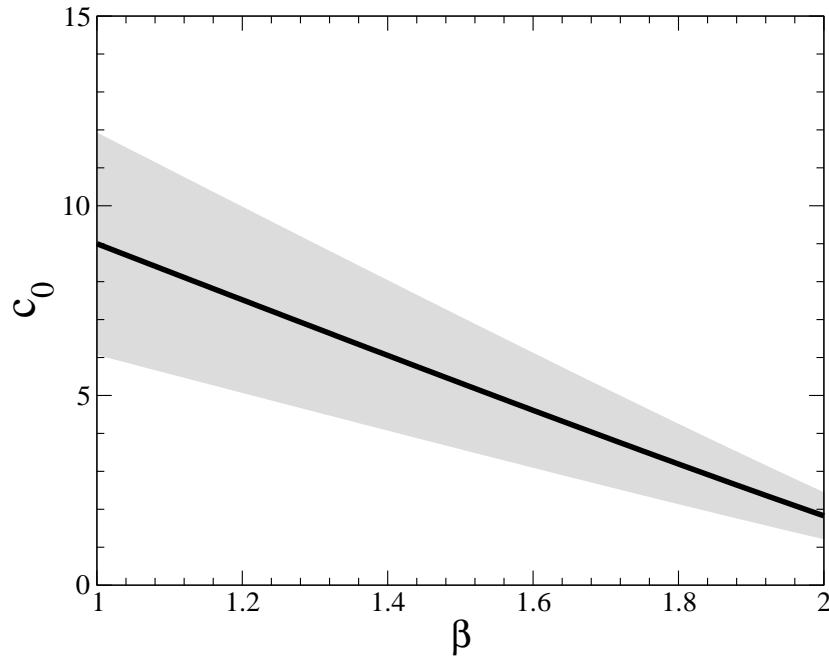
$$c(z) \equiv \frac{r_{200}(z)}{r_s(z)}$$

Given  $\beta$ ,  $c(z)$  and for a given mass  $M$ :

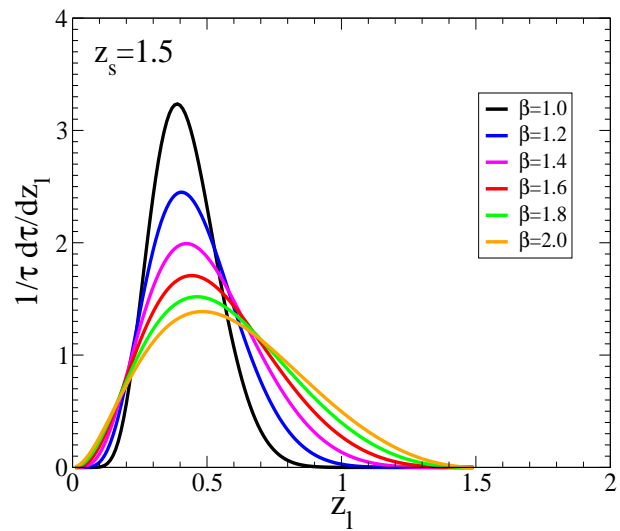
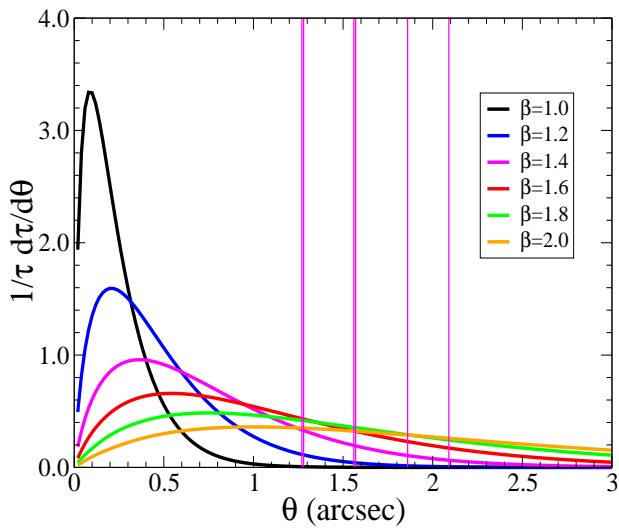
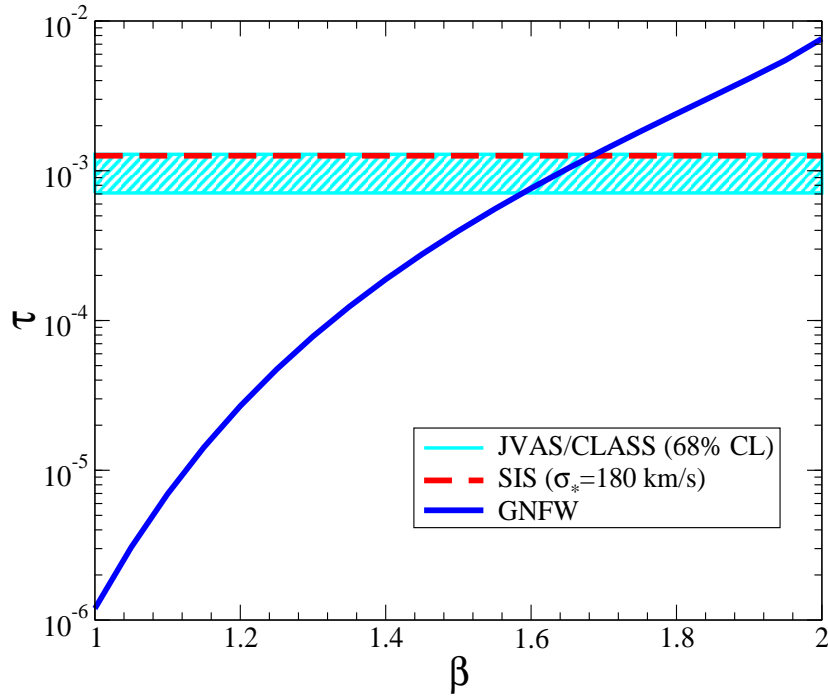
1. Compute  $r_s(z)$ ,  $\rho_s(z)$
2. Numerically solve the lens equation to get splittings, lensing cross-section
3. Numerically solve to get magnification bias

## Concentrating on the concentration...

$$c(z) = \frac{c_0}{(1+z)} \left( \frac{M}{M_*} \right)^{-0.13} \quad (\text{Bullock et al. 2001})$$



# Dependences on $\beta$



# Parameter set

Parameter we want:  $\beta$

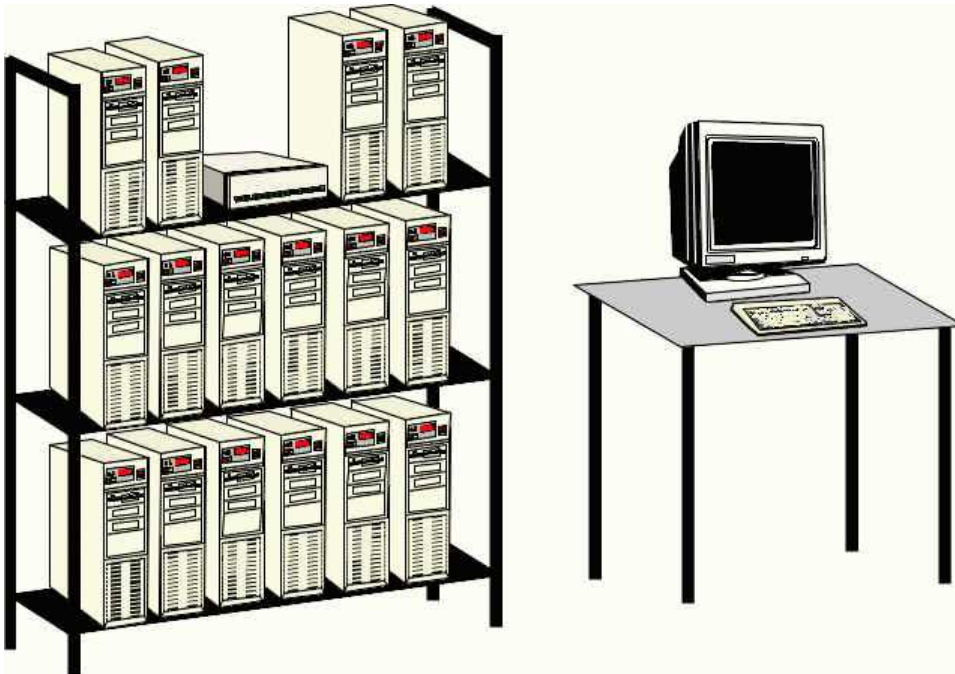
Parameters to marginalize over:

$$\sigma_* \in [120, 270] \text{ km/s}$$

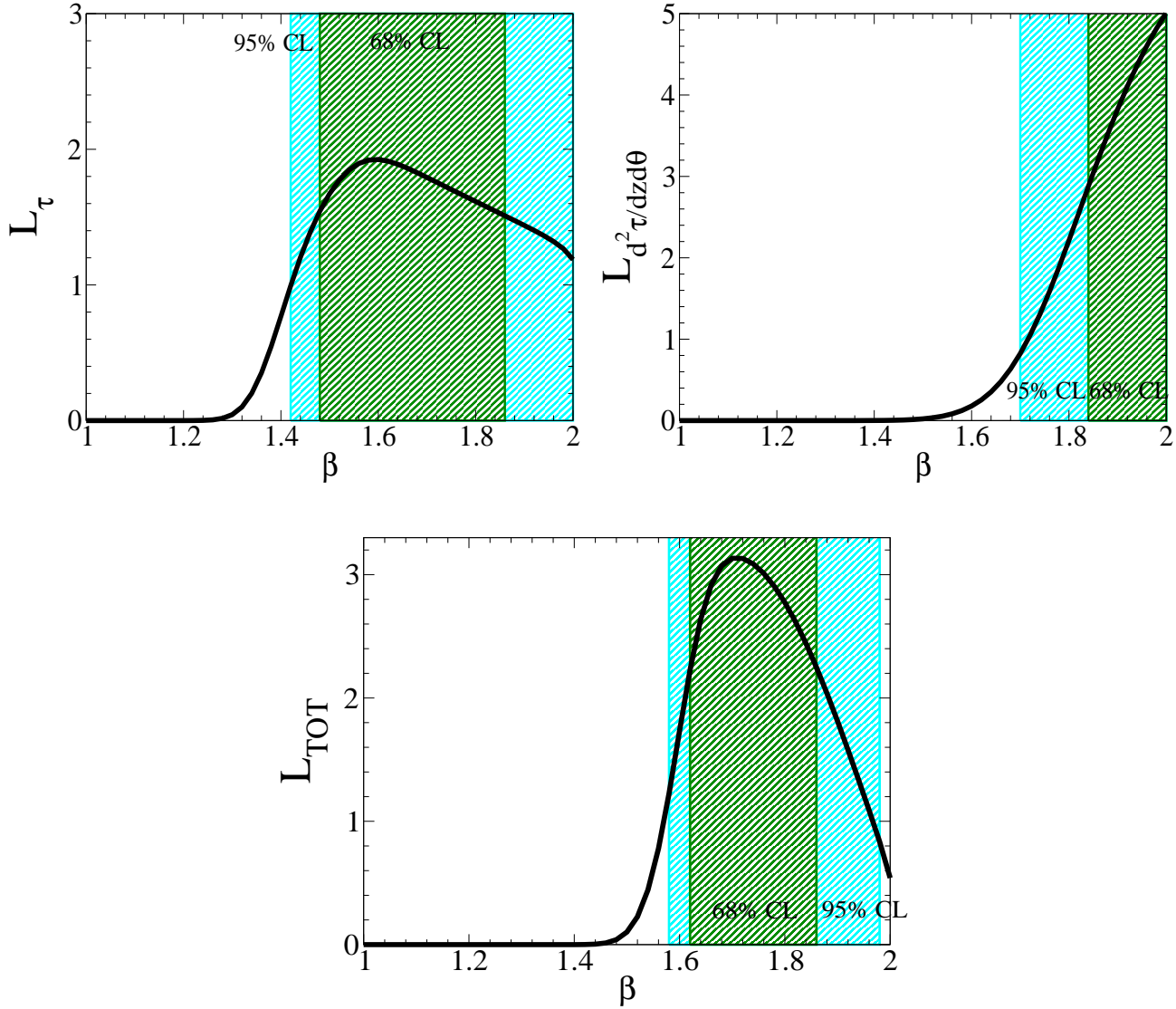
$$\phi_* \in [0.5, 1.5] \times 0.6 \times 10^{-2} h^3 \text{ Mpc}^{-3}$$

$$\gamma \in [3.5, 4.0]$$

$$\alpha \in [-1.3, -0.7]$$



# Results – GNFW Profile

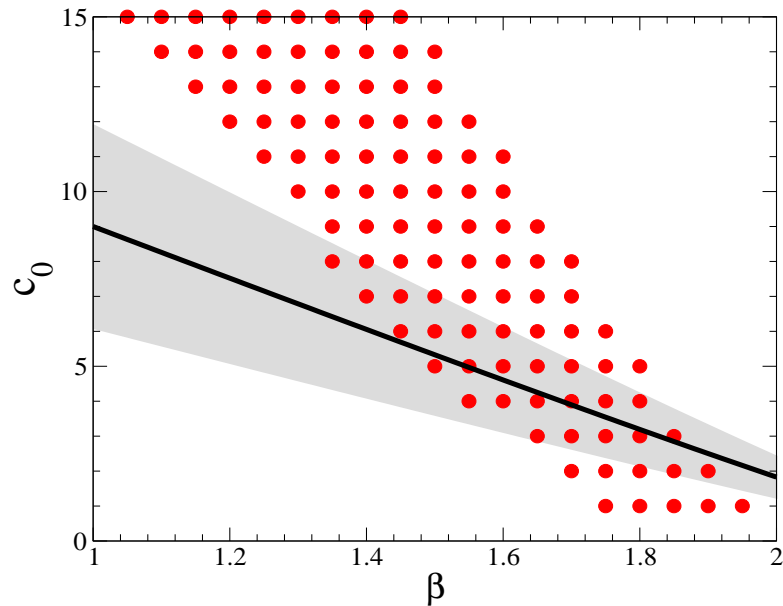


$$1.58 \leq \beta \leq 1.98 \text{ (95\% CL)}$$

→ NFW profile ( $\beta = 1$ ) strongly ruled out

# GNFW Results

Let  $c_0$  be a completely free parameter, then



Shallow NFW profiles are expected to become steeper galaxy profiles because:

1. Baryonic infall will make the profile steeper, roughly SIS
2. Massive black holes in centers of galaxies would steepen the profile
3. No detected odd images  $\Rightarrow \beta \simeq 2$  (Rusin & Ma 2001)

or maybe not:

1. Direct lens modelling gives  $\beta$  anywhere between 1 and 2
2. Merging BHs would have the opposite effect (Meritt & Milosavljević)

# Redshift evolution

The main caveat: redshift evolution of the

1. Number density of galaxies, and/or
2. Lens profile

So, we have measured the “redshift-averaged  $\sigma_*$ ” (and  $\beta$ )



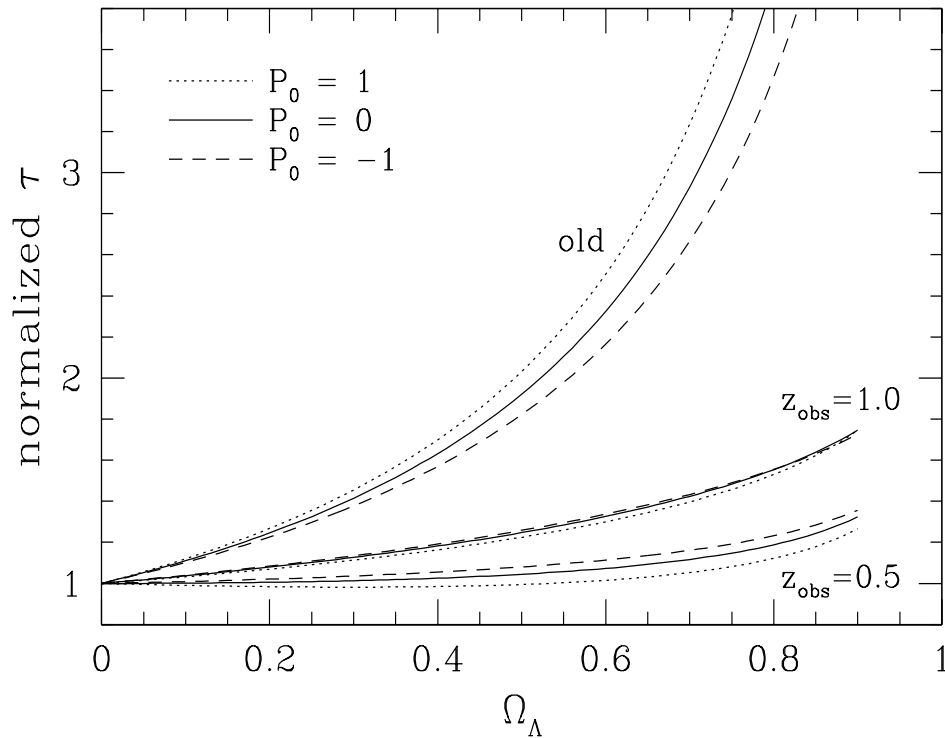
# Redshift evolution

Assume e.g.

$$\phi_*(z) = \phi_*(0)10^{0.4Pz}$$

Keeton (2001): why not **normalize** lensing statistics to the number counts of galaxies at some  $z$ ?

- Old way: fixed, “known”  $\phi_*(z)$
- New way: variable  $\phi_*(z)$  so as to reproduce  $N_{\text{gal}}(z) = \phi_*(z) \times \text{Volume}(\Omega_M, \Omega_\Lambda)$



But:  $\sigma_*$  will also be  $z$ -dependent, may cancel the effect of  $\phi_*(z)$ .

## Conclusions

- CDM may have some problems on small scales / central galaxy profiles
- Strong (and weak) lensing are ideal tools to test/constrain CDM
- Strong lensing best used to probe lens (galaxy, halo) parameters, and not cosmology
- In fact, cosmological constraints due to strong lensing cannot be trusted
- No lenses at  $\theta > 3'' \implies$  two populations of galaxies:  
SIS “galaxies” and NFW “clusters”
- Our analysis:  $\sigma_* \approx 180$  km/s,  $\beta \approx 1.6-2.0$  for ellipticals
- More lenses needed!