

# Turbulence and Conduction in Galaxy Clusters

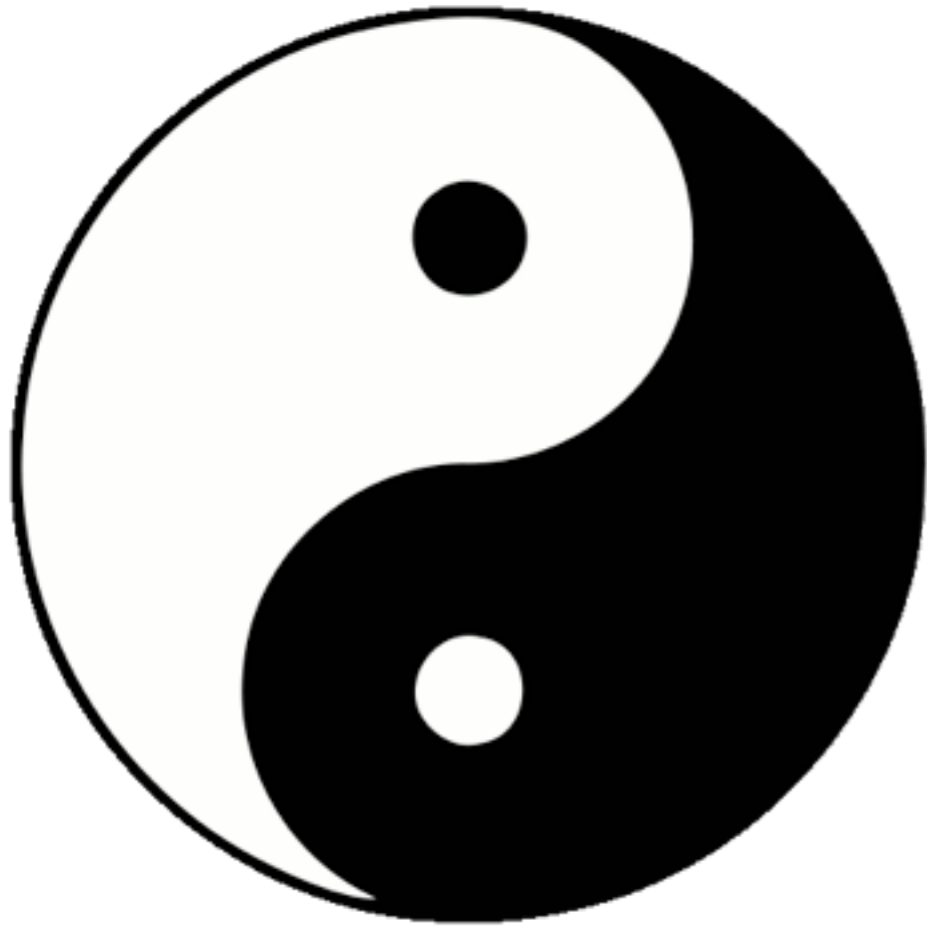
Peng Oh (UCSB)

# Co-Conspirator



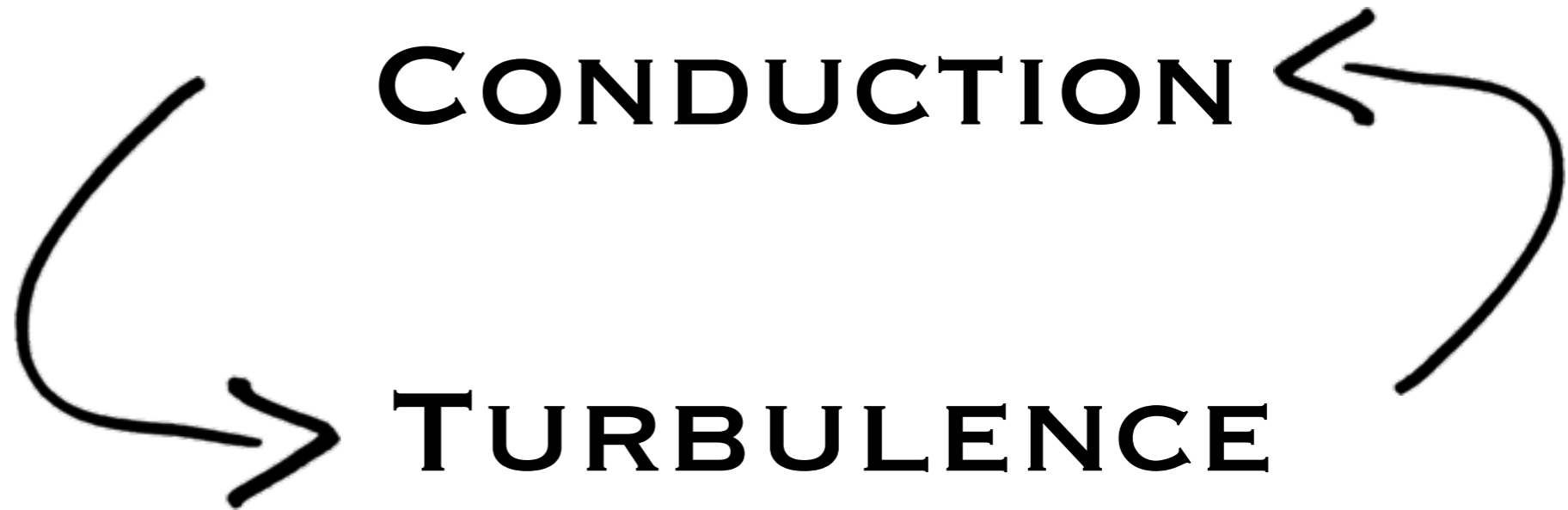
Comrade Ruszkowski

# THE TAO OF HEAT TRANSPORT



**CONDUCTION**

**TURBULENCE**





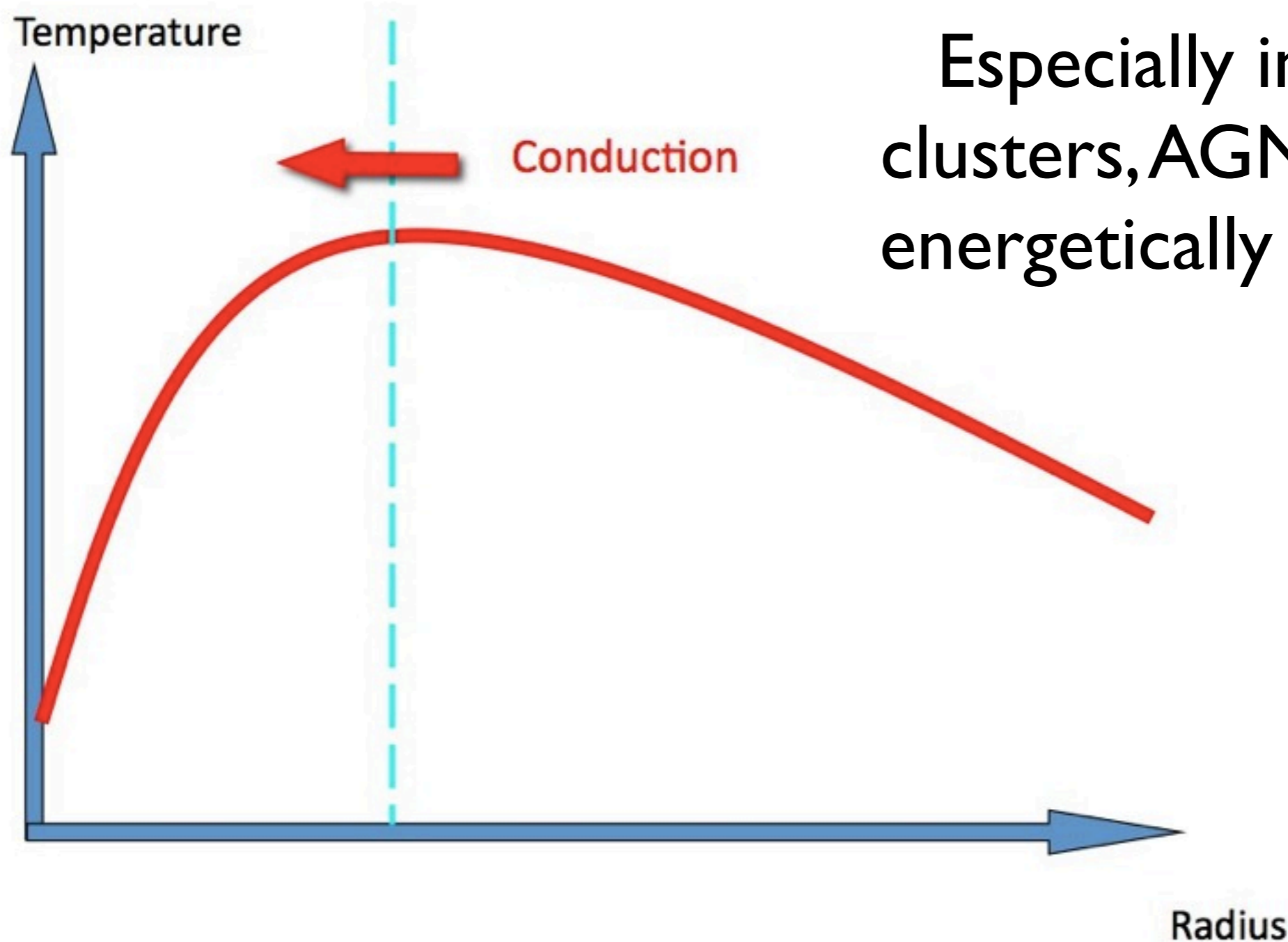
# Shaken and Stirred: Turbulence and Conduction in Clusters

Ruszkowski & Oh, 2010, ApJ, 713, 1332

Also Parrish et al (2010)

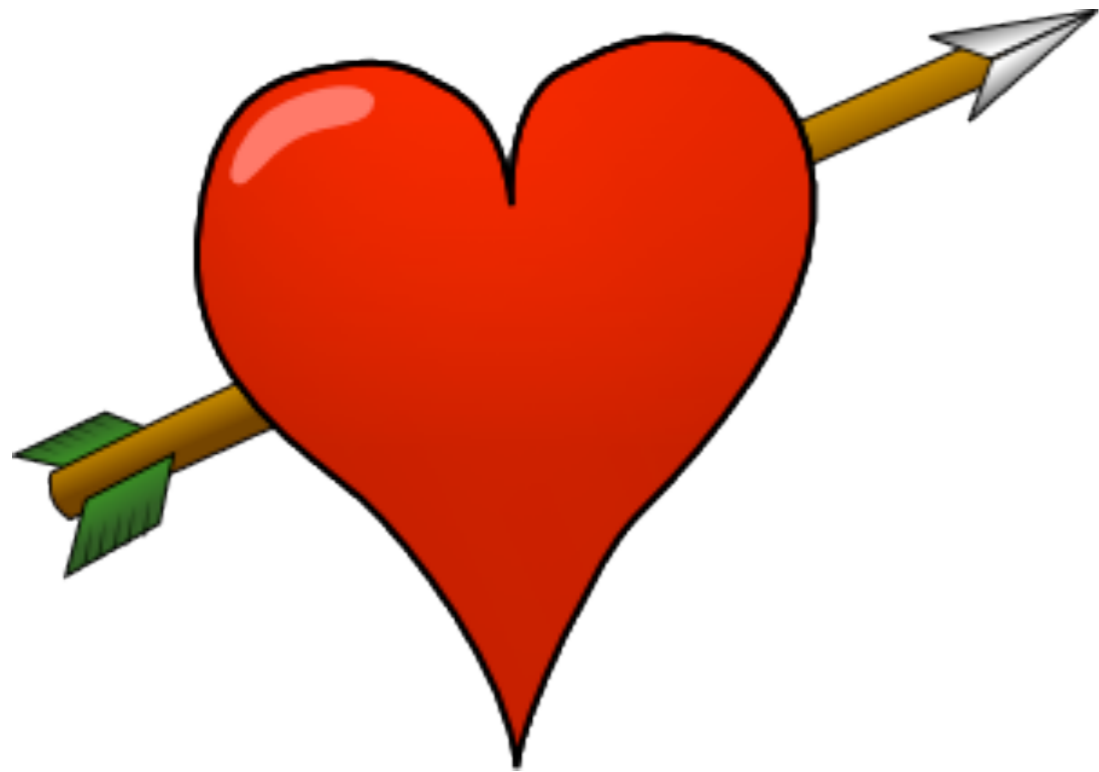


# Thermal Conduction can supply heat to cool core



Especially in massive clusters, AGN heating is energetically insufficient

# Conduction is a fickle lover



Cluster becomes  
isothermal



Cooling catastrophe

There's a thin line between love and hate...

# Use global stability analysis to understand

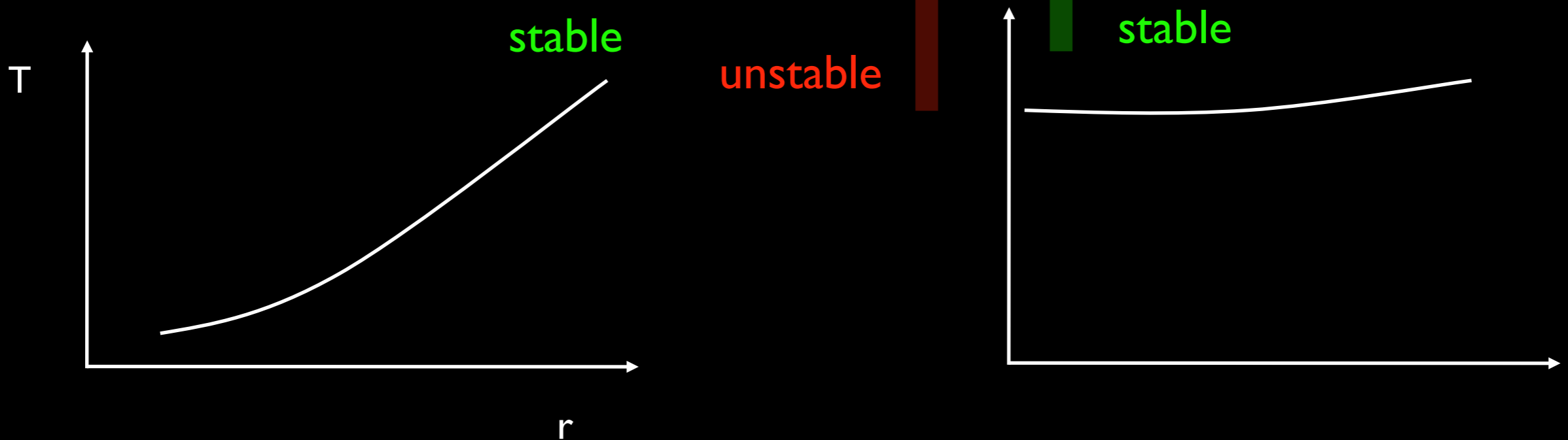
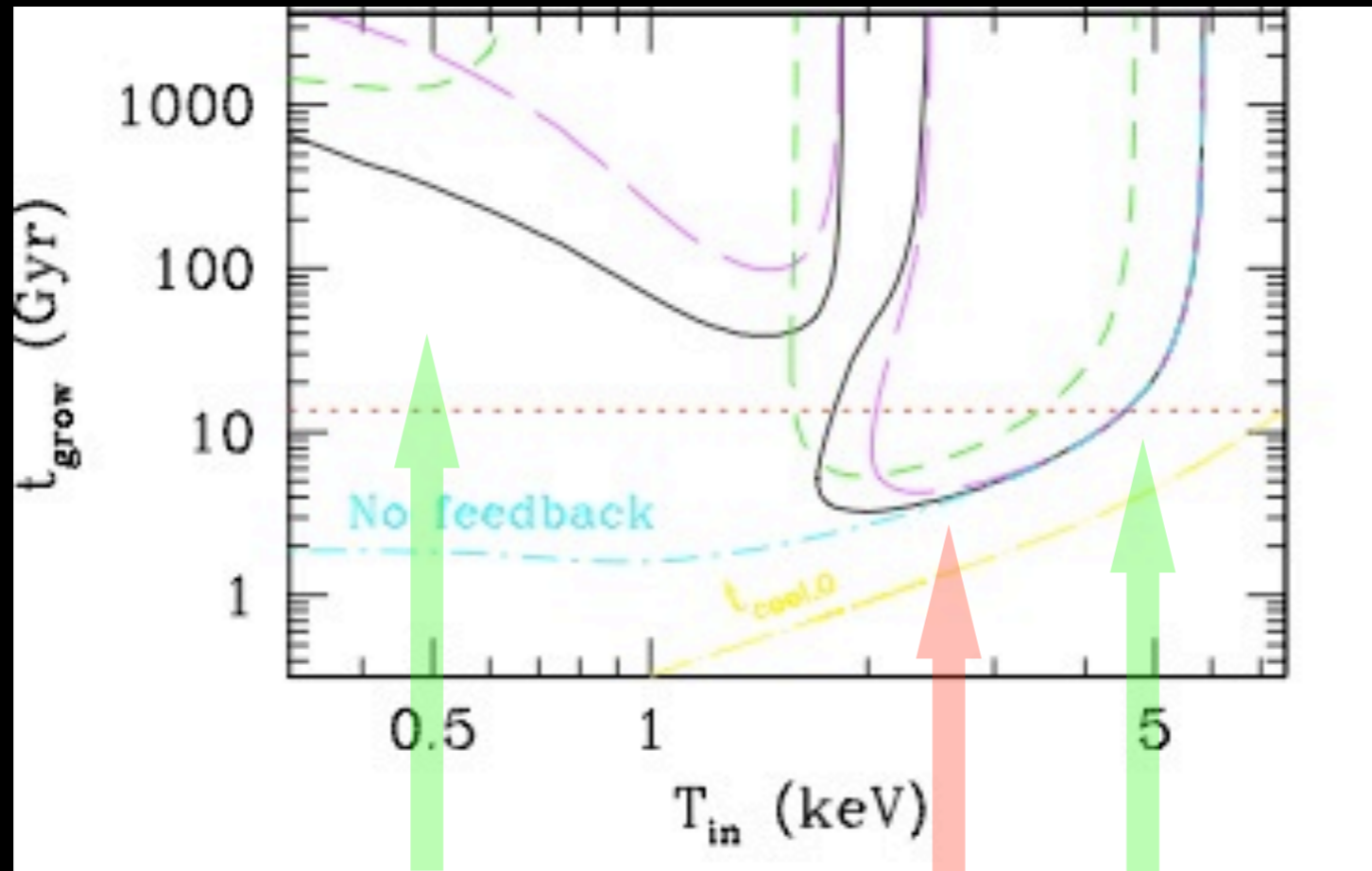


‘Pick any two’

Semi-analytic study: Explore parameter space **quickly**

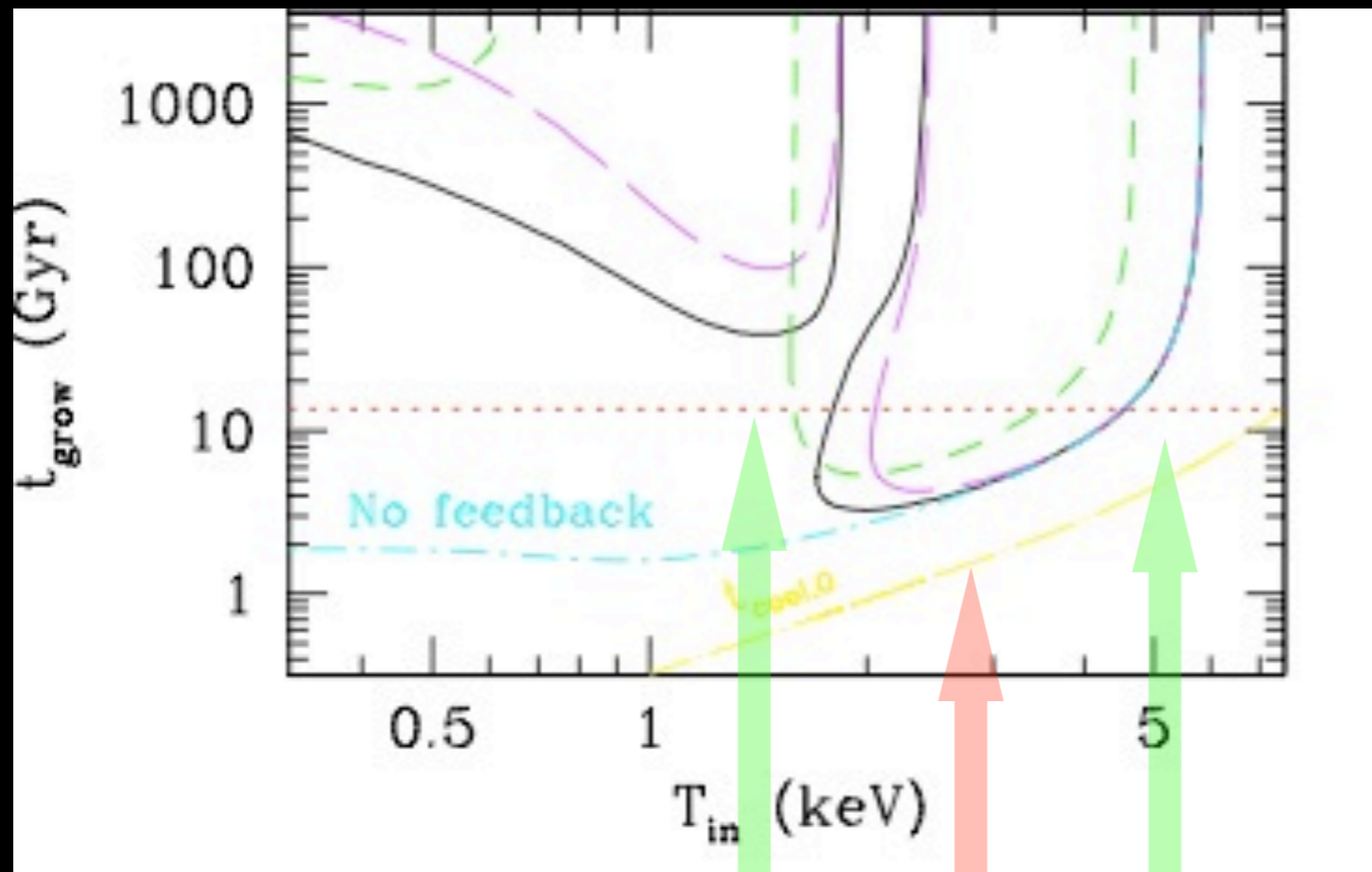
Set up equilibrium model, then study globally unstable modes

**Guo, Oh & Ruszkowski (2008)**



**Stability is bimodal!**





stable

unstable

stable

## Bimodality !

cool core



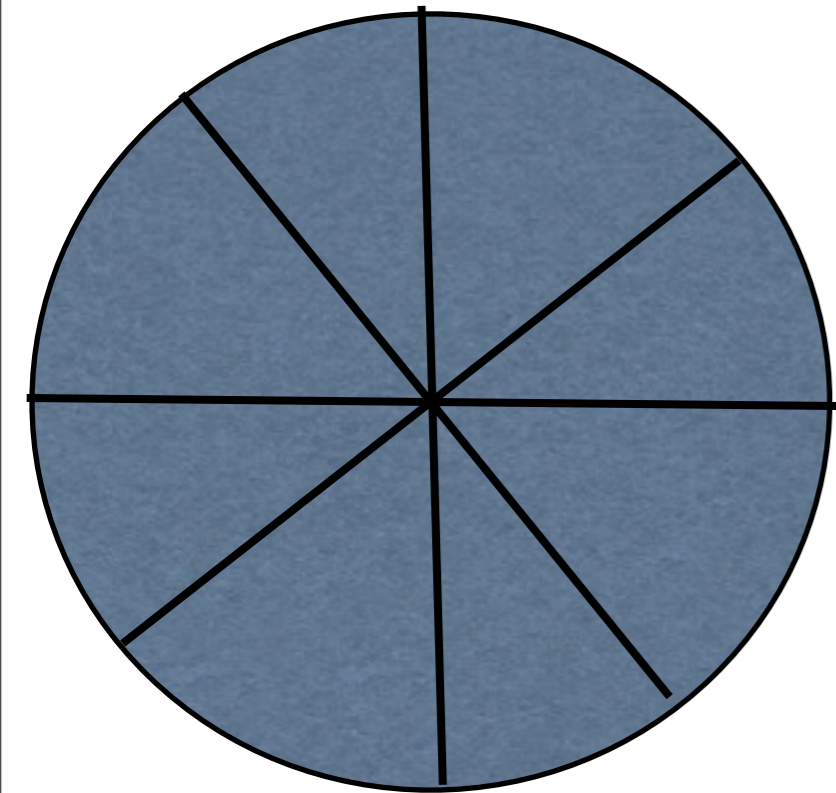
stabilized by AGN + conduction

non cool core

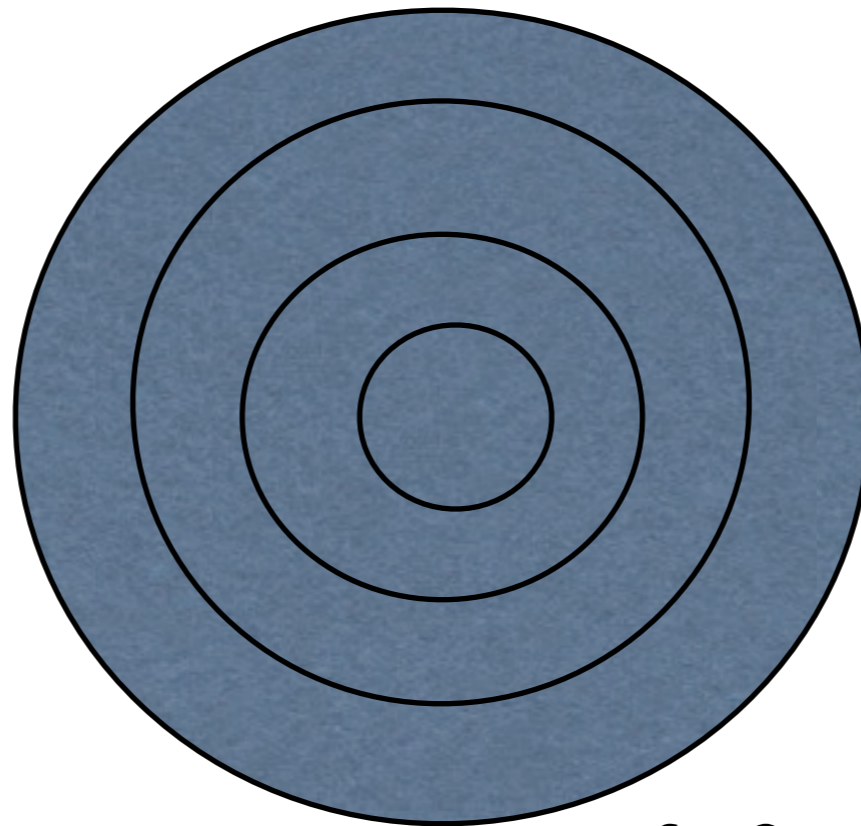


stabilized by conduction

# Efficiency depends on unknown B-field topology



$f=1$



$f=0$



$$f = \langle \cos^2 \theta \rangle = 0.33$$

Note: this presumes  $l_T \gg l_B \gg \lambda_e$

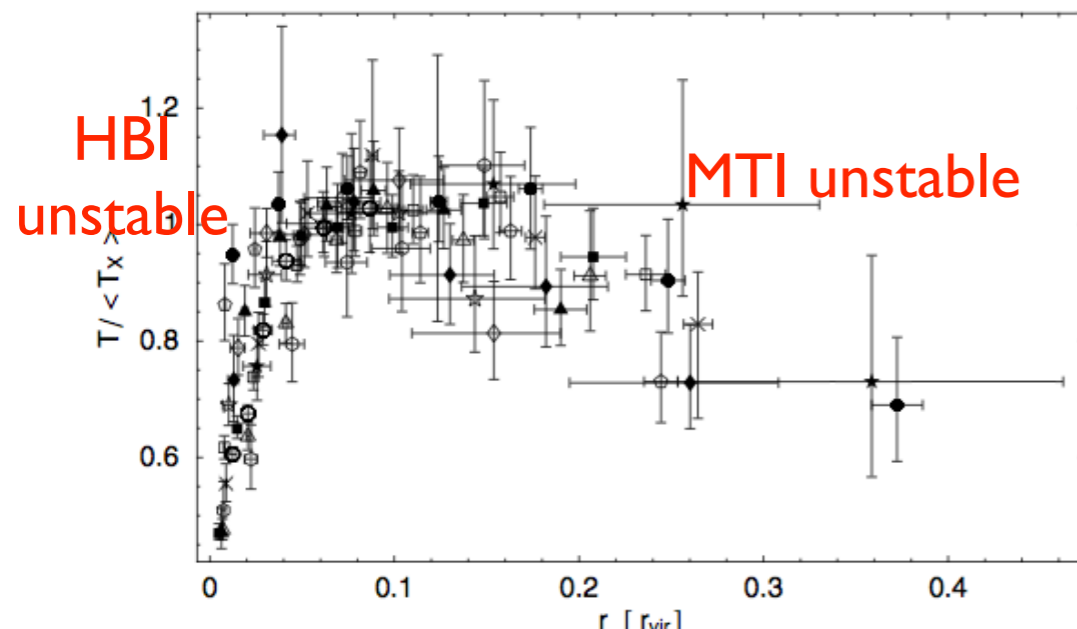
# Buoyancy instabilities realign magnetic field

$\nabla T < 0$   Magnetothermal Instability (MTI)  
(Balbus 2000)

Radial magnetic fields

$\nabla T > 0$   Heat buoyancy instability (HBI)  
(Quataert 2008)

Tangential magnetic fields



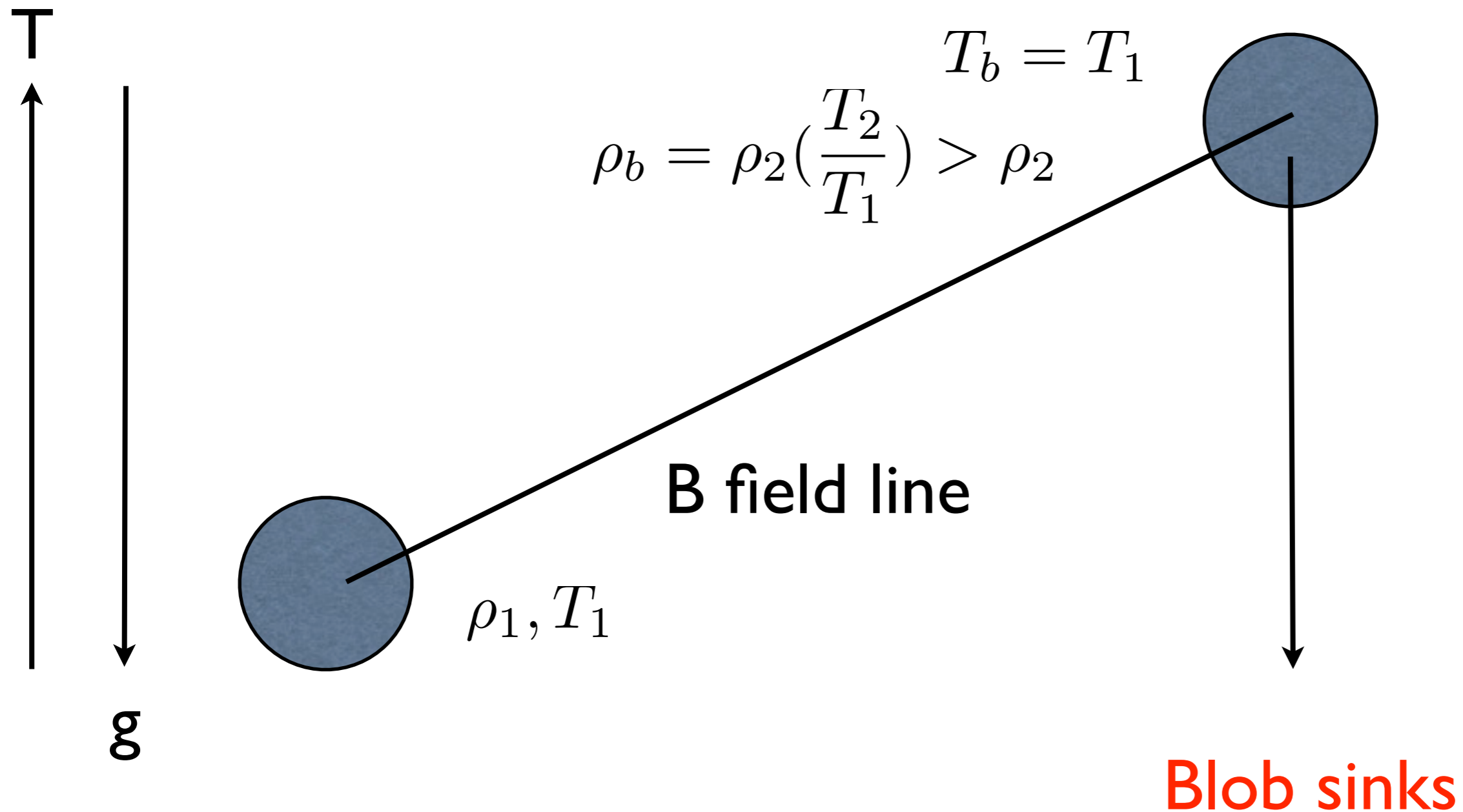
# How does it work?

Consider HBI case

$$T_b = T_1$$

$$\rho_b = \rho_2 \left( \frac{T_2}{T_1} \right) > \rho_2$$

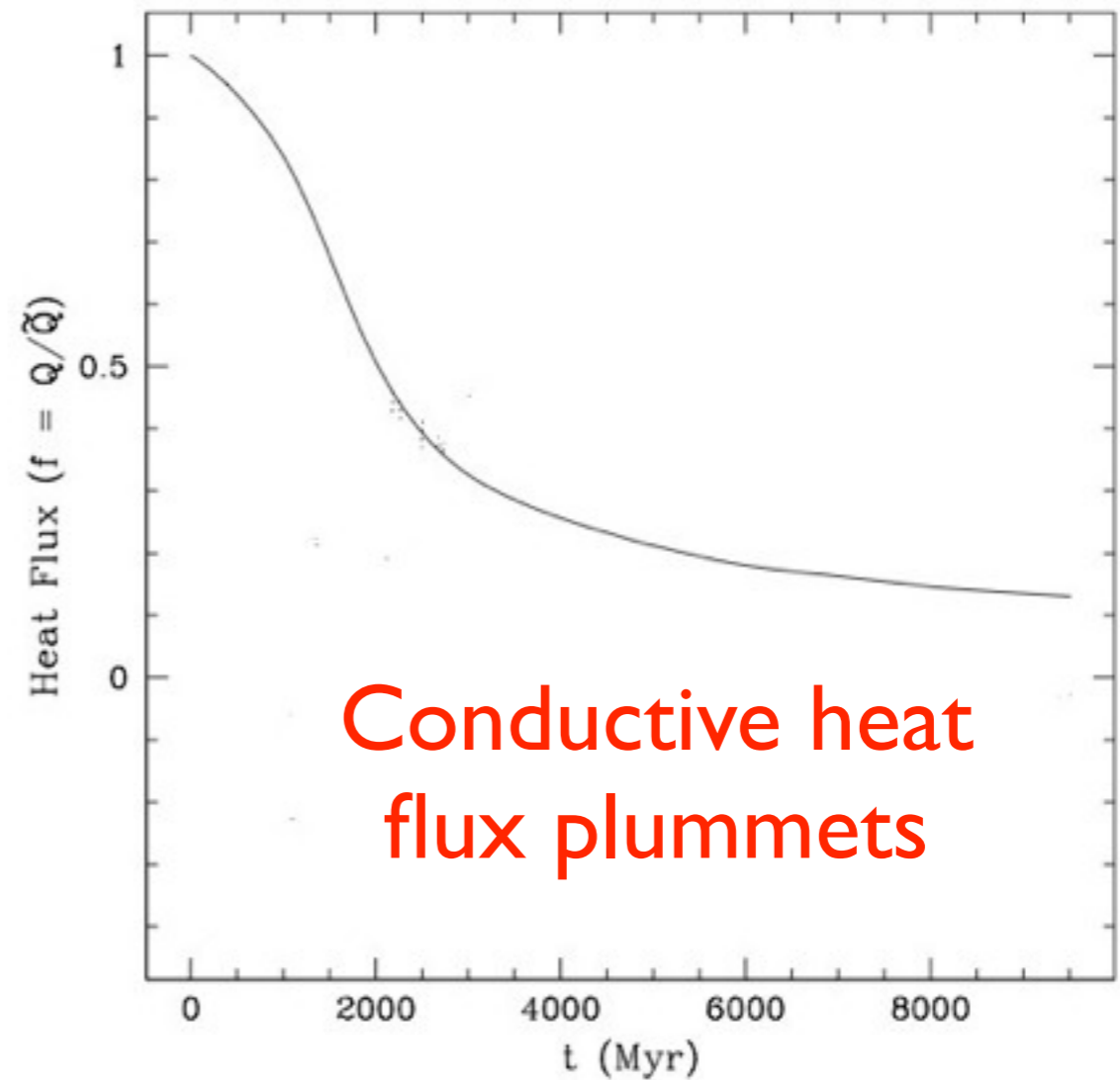
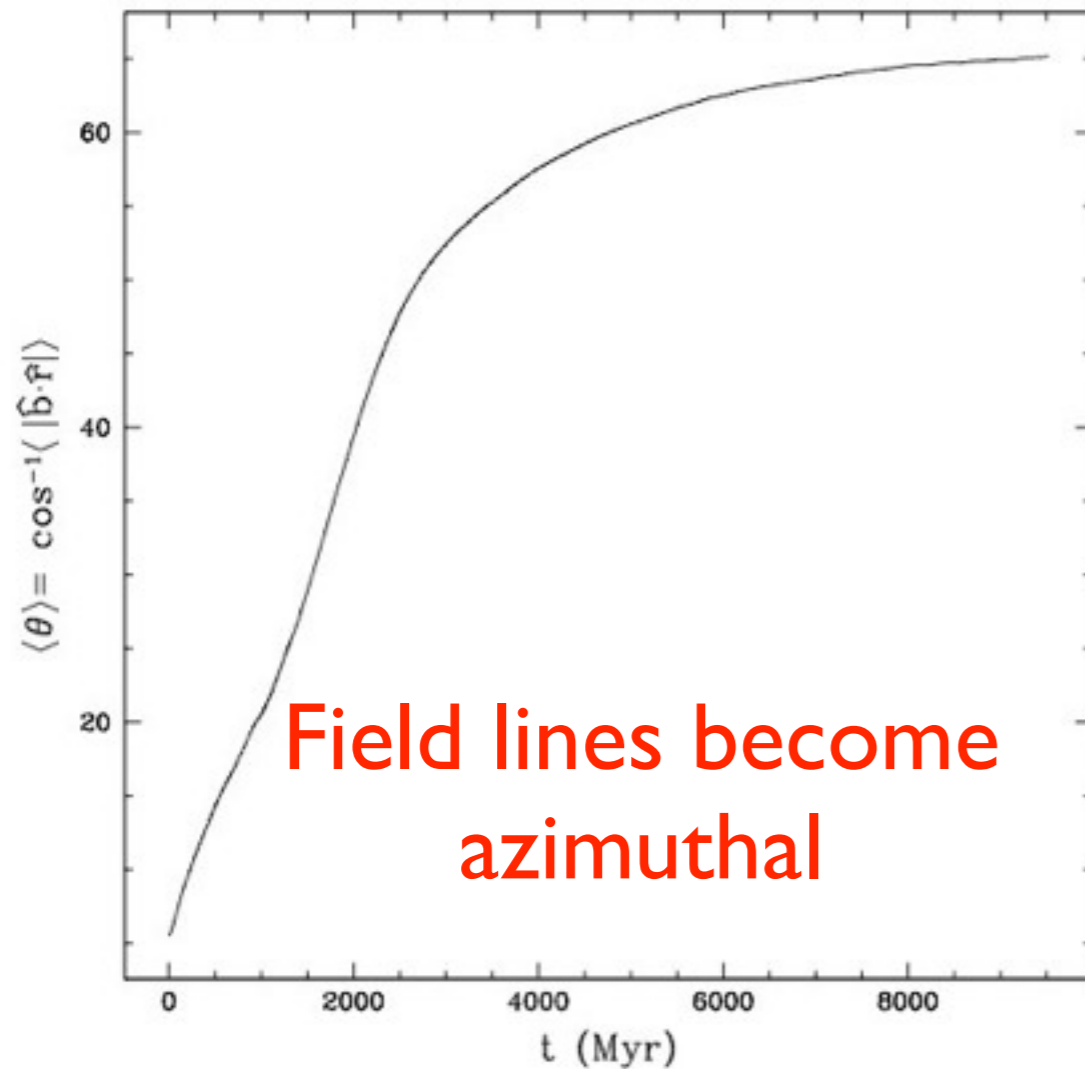
$\rho_2, T_2$



**Blob sinks**



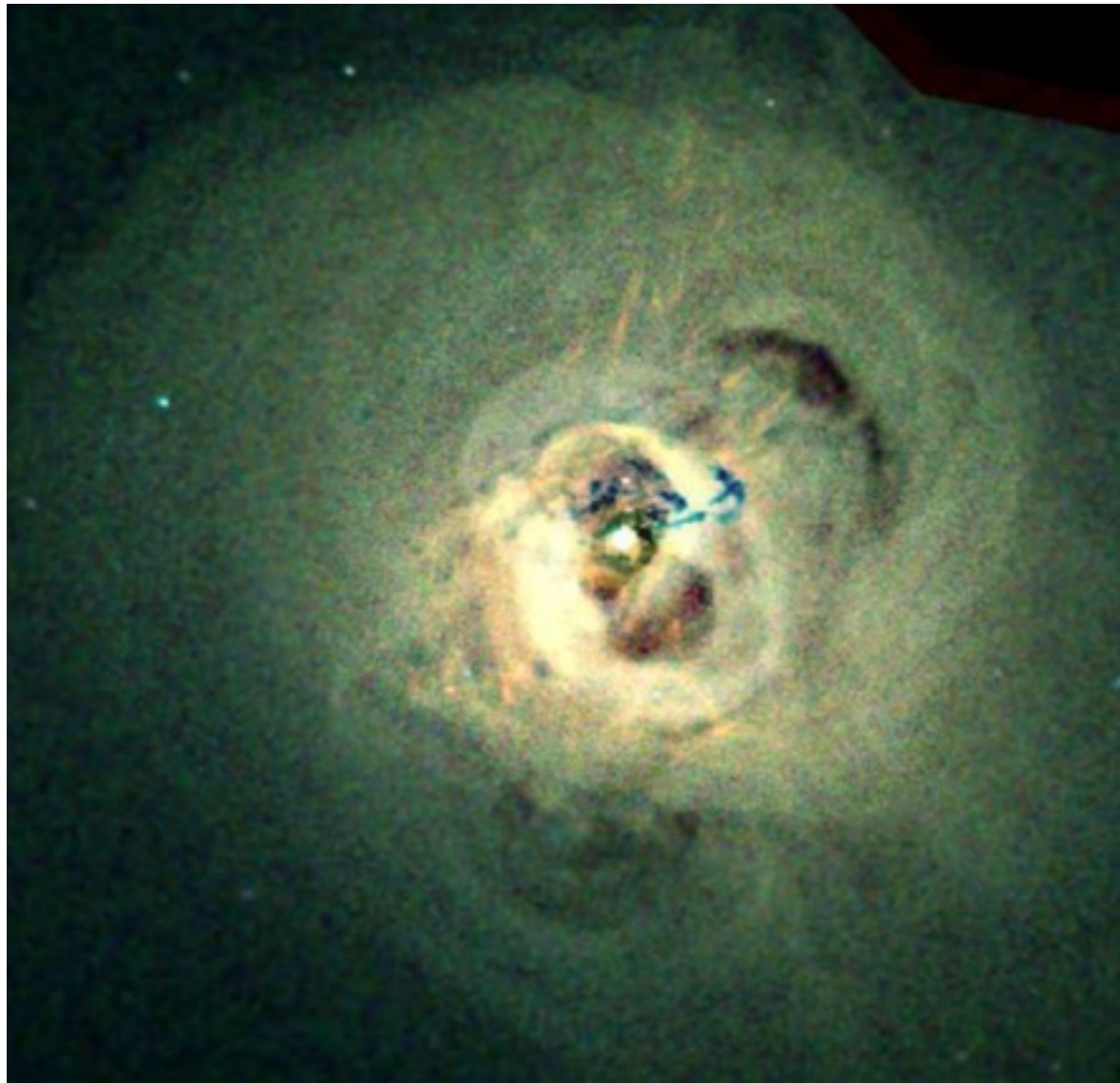
...this shuts off thermal conduction



and brings on a cooling catastrophe

Parrish et al (2009) (see also Bogdanovic et al 2009)

# We expect the ICM to be turbulent



Evidence from:

Lack of resonance scattering lines (Churazov et al 2004)

Analysis of pressure maps (Schuecker et al 2004)

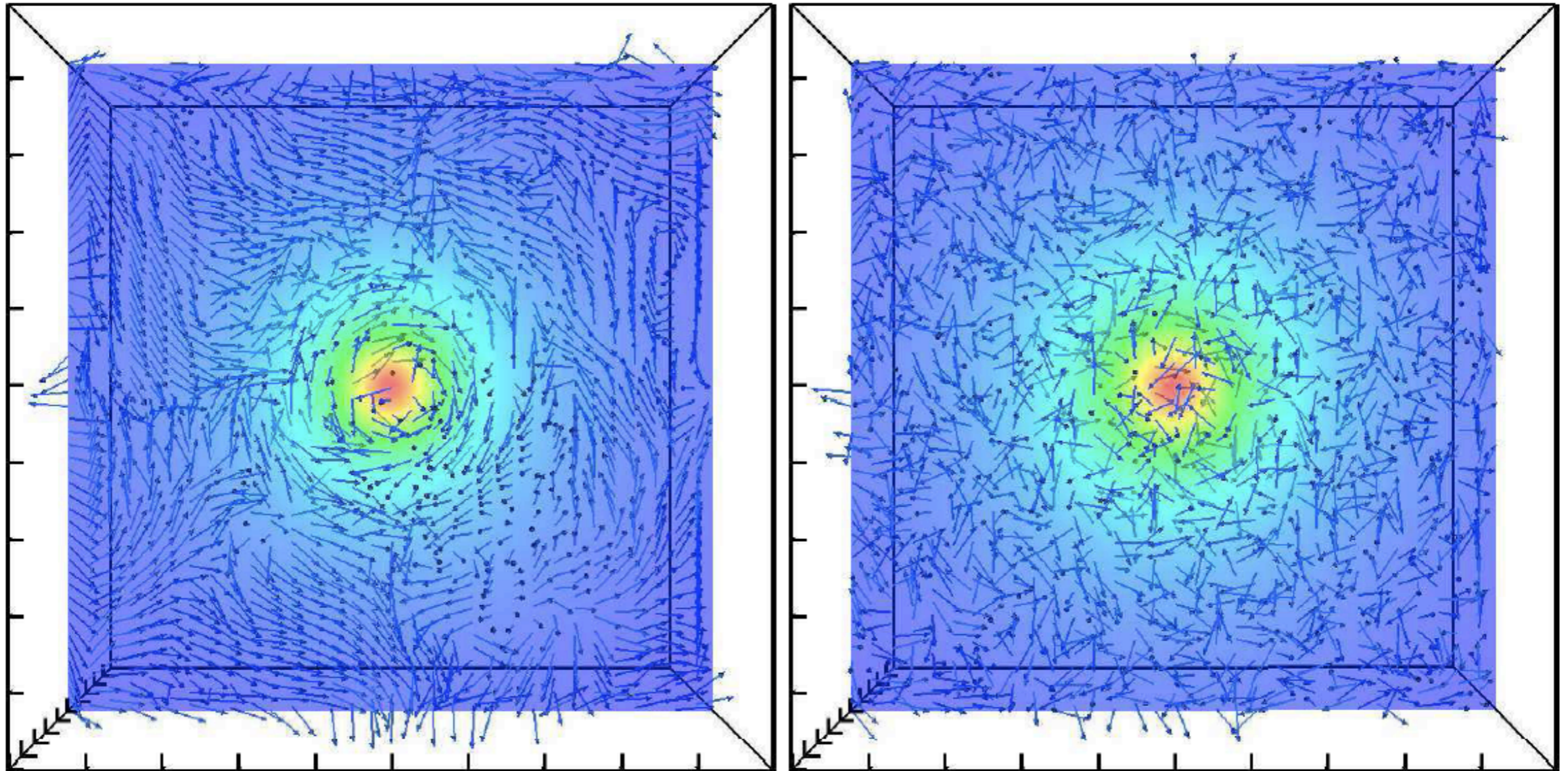
Faraday rotation maps (Ensslin & Vogt 2006)

X-ray spectroscopy upper bounds (Sanders et al 2010)

Could turbulence randomize field lines and restore conduction?



# Bottom Line: YES

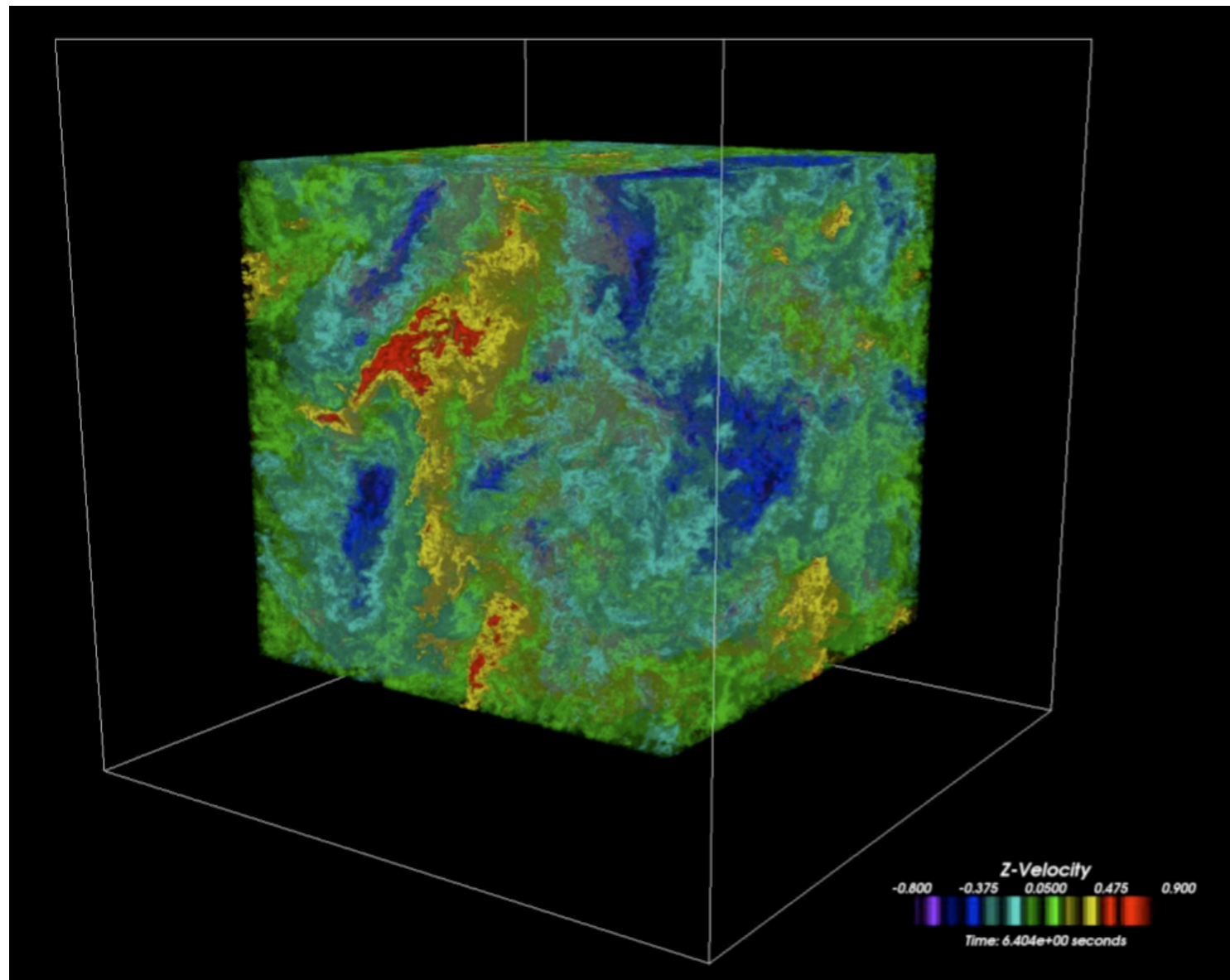


Ruszkowski & Oh (2010a), using FLASH

Similar results by Parrish et al (2010), using ATHENA



# Simulation details



FLASH turbulence module  
(Fisher et al 2008)

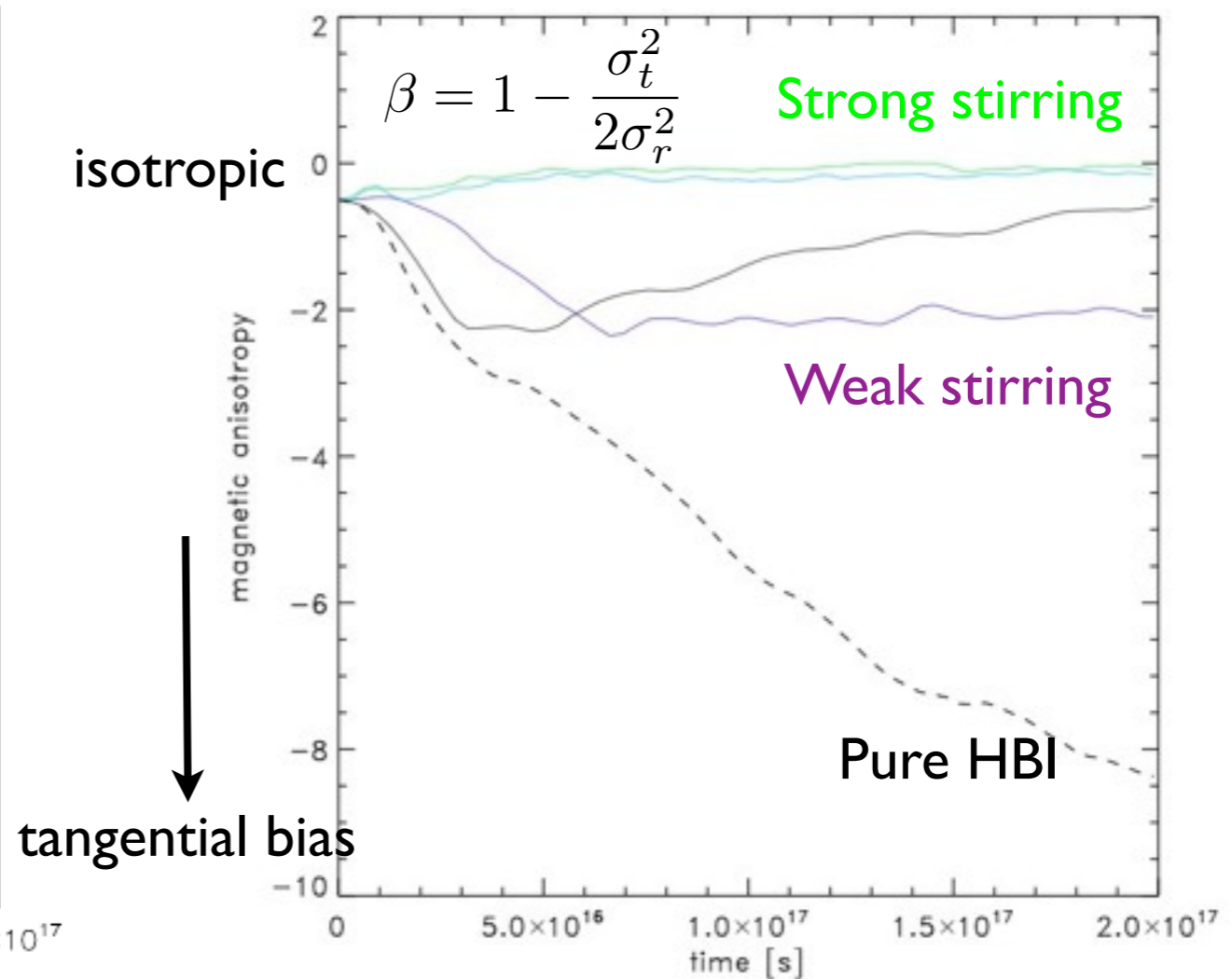
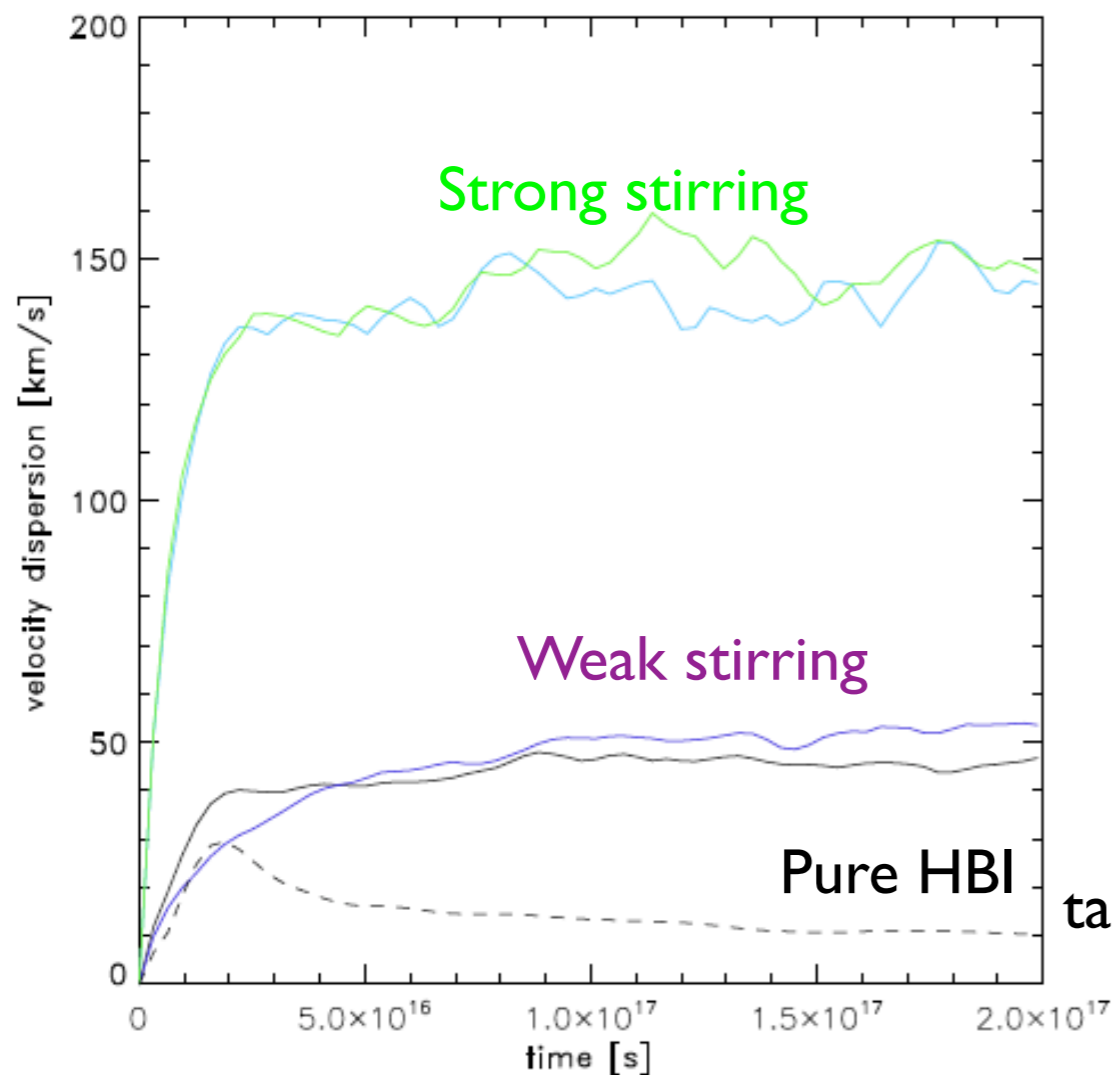
3D MHD FLASH  
simulations

1 Mpc<sup>3</sup>, max resolution  
2.7 kpc h<sup>-1</sup>

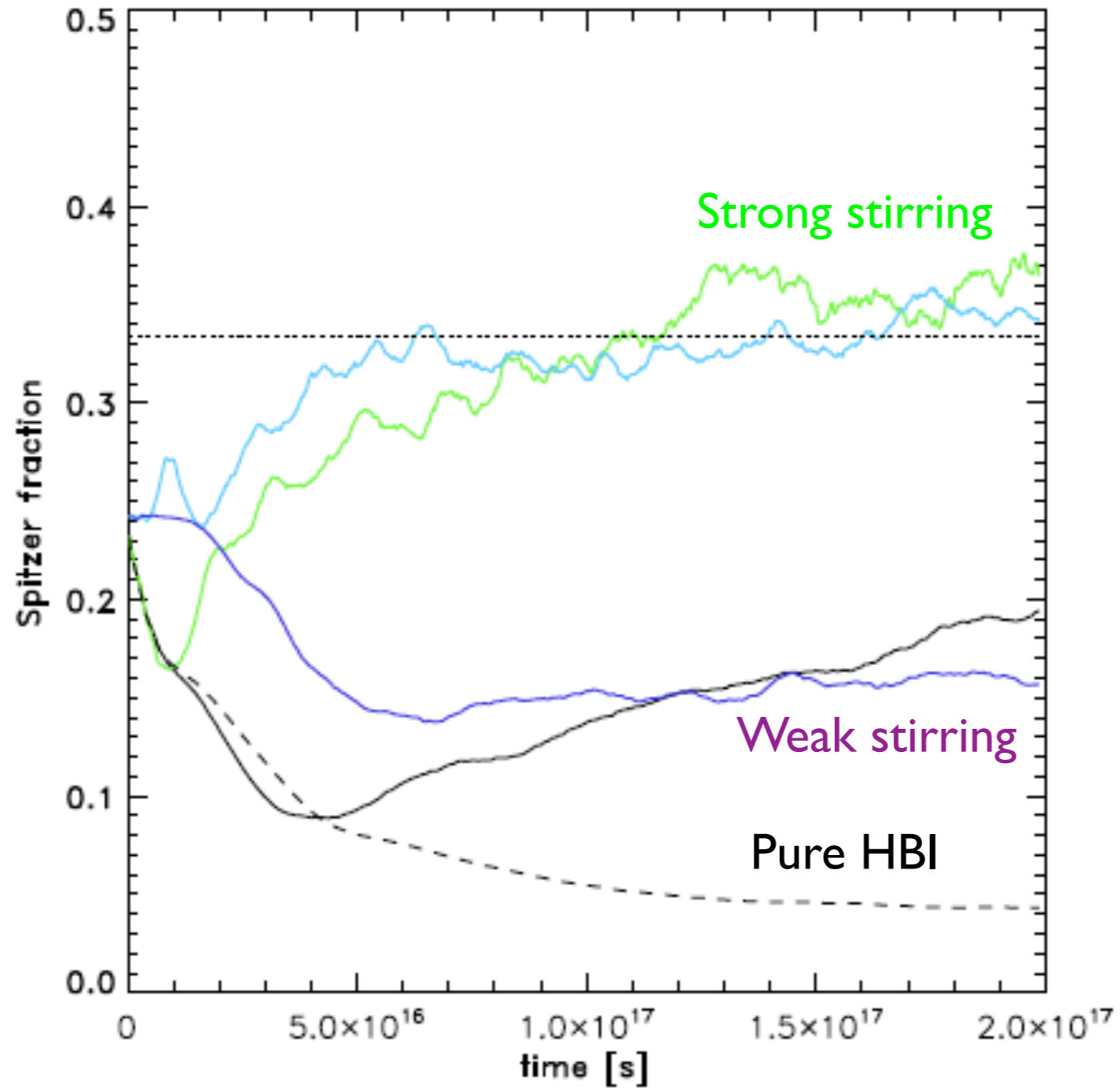
Spectral  
implementation of  
turbulent forcing  
(note: volume-filling by  
construction)



# Required Amount of Turbulence is **small**



Highly subsonic ( $\sim 10\%$  of sound speed) motions overcome HBI, randomize fields

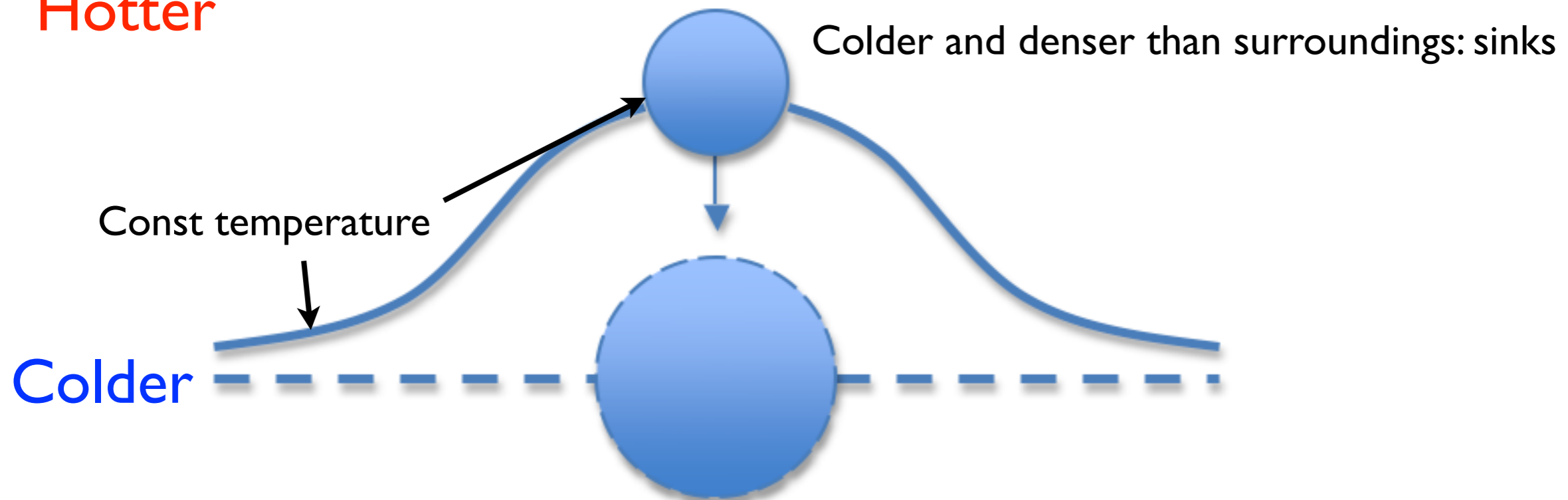


...and conduction is restored!

# Why are such weak motions effective?

Stratified by **temperature**, rather than entropy.

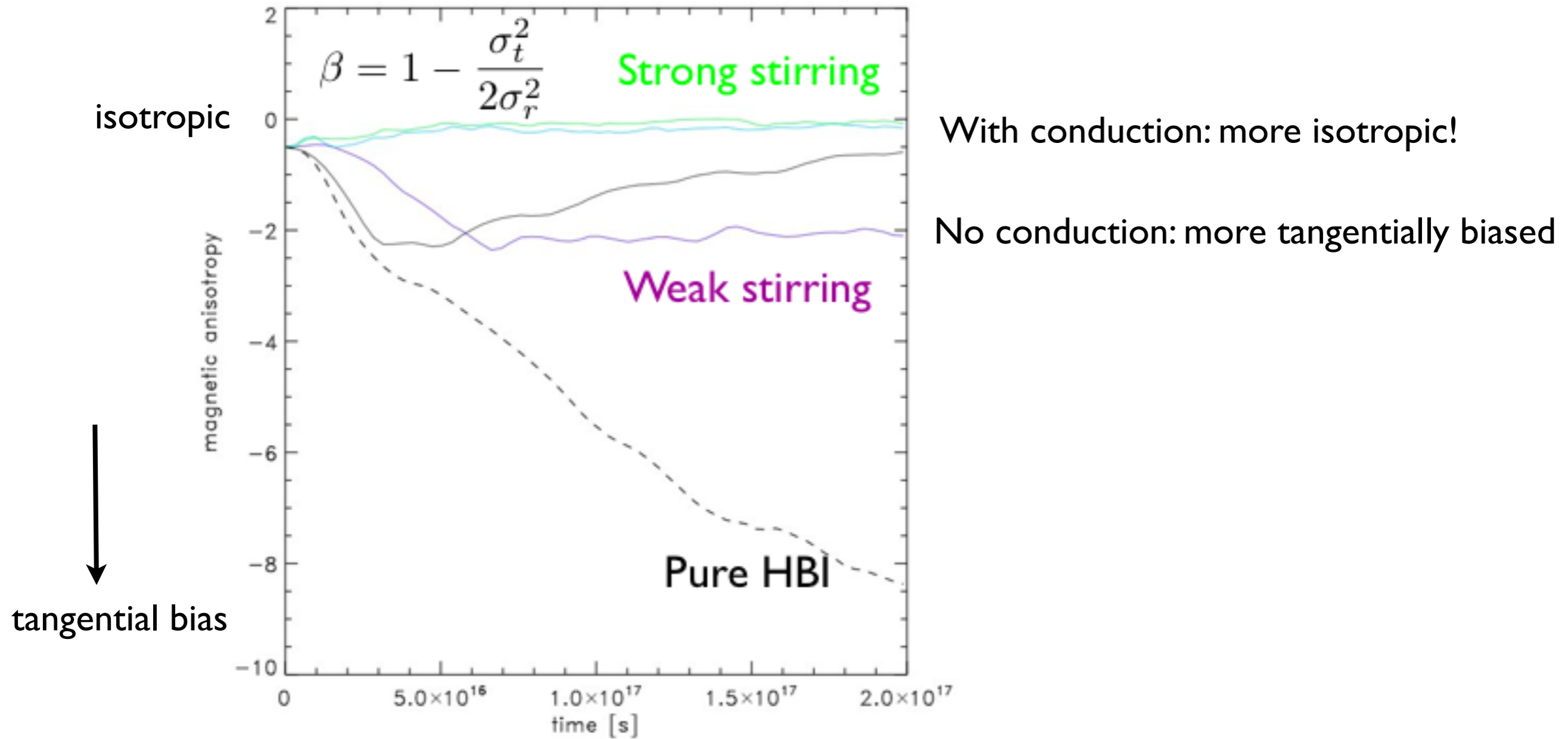
Hotter



Oscillates with frequency  $(\omega_{\text{BV}}^{\text{MHD}})^2 = \frac{g}{r} \frac{d \ln T}{d \ln r}$ ,

**Buoyancy forces are weaker in conducting plasma**

# We can see this already: different anisotropy for same amount of turbulence





# Compare buoyancy force to inertial term (Richardson number, Froude number)

$$\frac{v_z}{v} \sim \left( \frac{\omega}{\omega_{BV}} \right)^2$$

$\omega \ll \omega_{BV} \Rightarrow$  Largely tangential motion

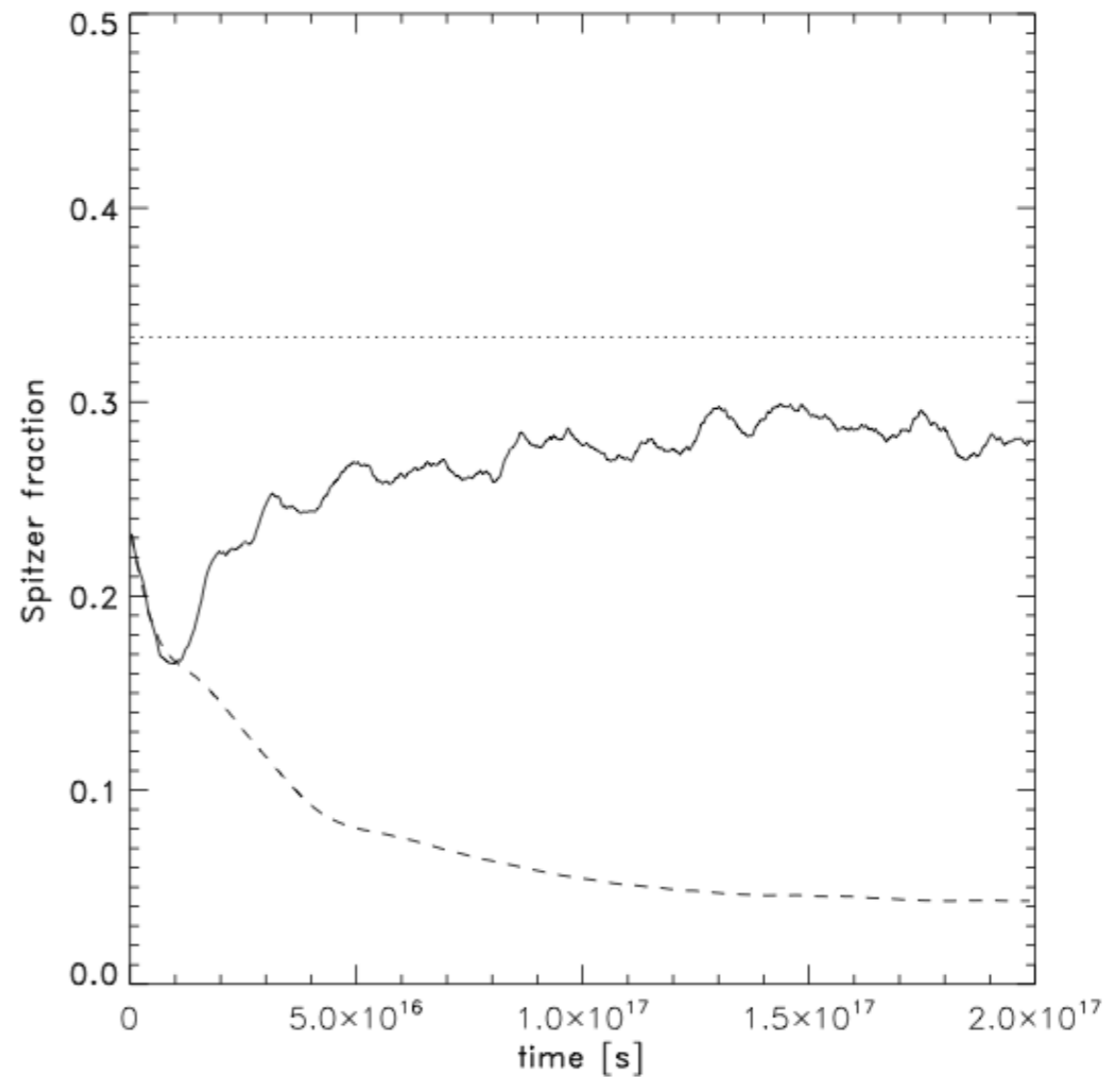
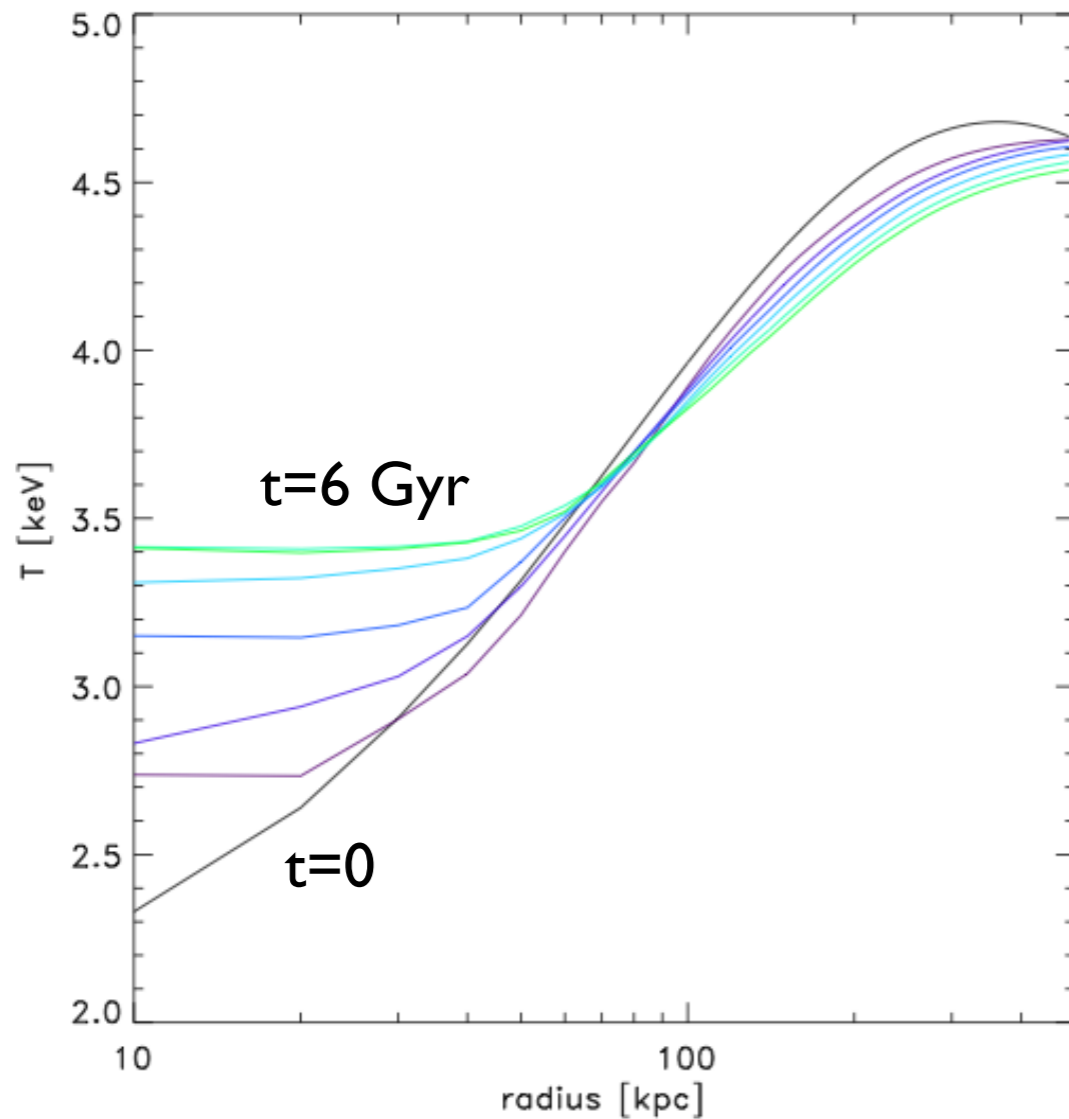
$\omega \gg \omega_{BV} \Rightarrow$  Isotropic motion

Required velocities are small

$$\sigma \approx 135 \text{ km s}^{-1} g_{-8}^{1/2} r_{10}^{1/2} \left( \frac{d \ln T / d \ln r}{0.15} \right)^{1/2} \left( \frac{Ri_c}{0.25} \right)^{-1/2}$$

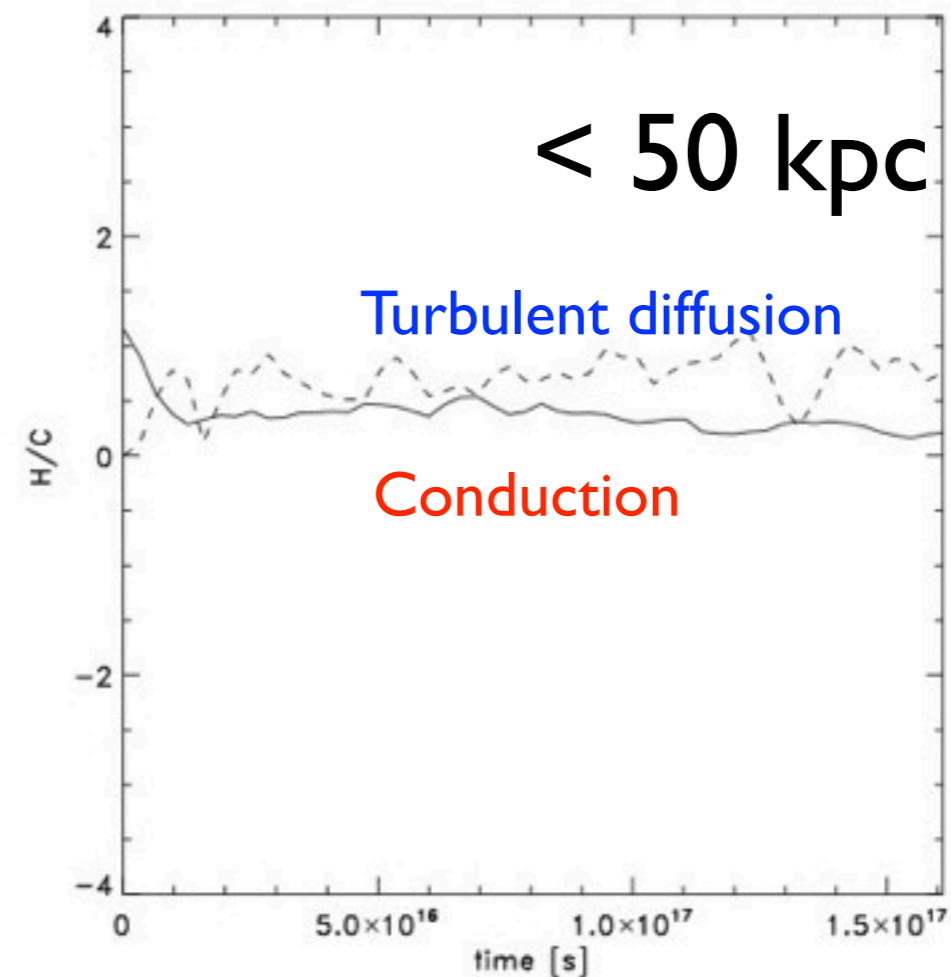
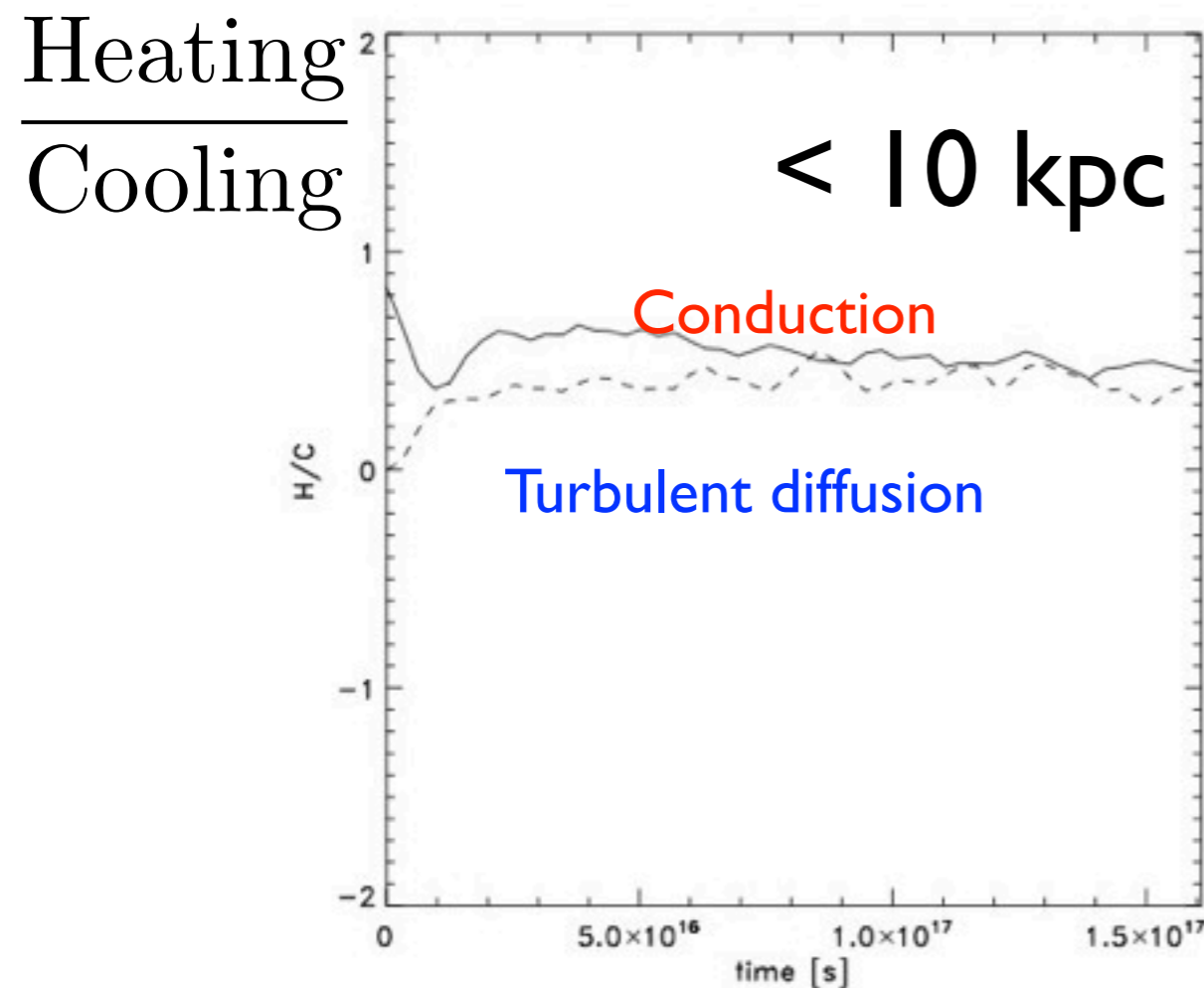
See also Sharma et al (2009)

# No cooling catastrophe!



Cluster transitions from cool core to non cool core...

# Turbulent heat diffusion is important...



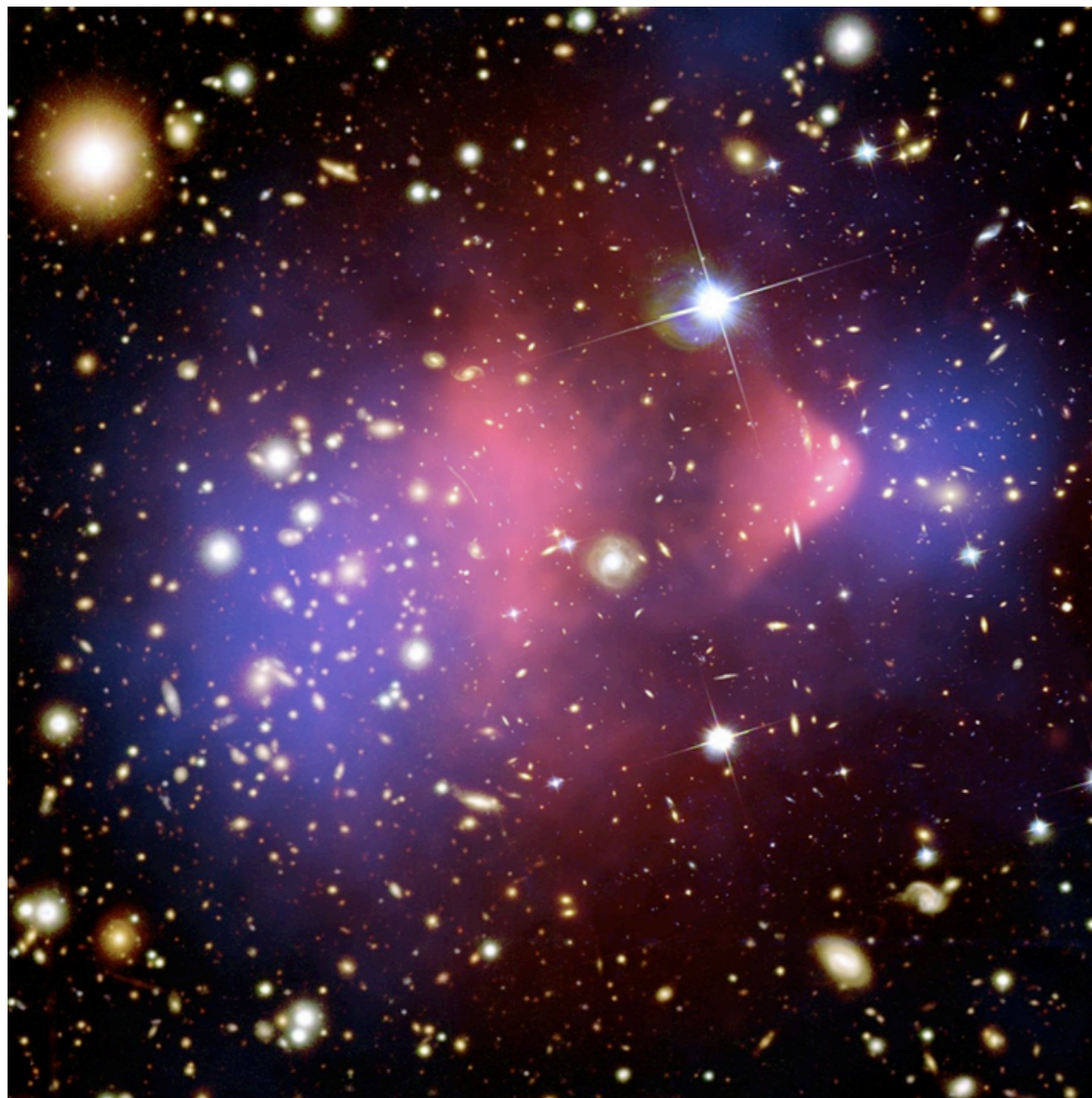
...but even when subdominant, conduction is crucial  
(more later)

Can we have sustained,  
volume-filling turbulence?

Ruszkowski & Oh (2010b, on astro-ph vveerrry soon...)



# Most sources of turbulence are intermittent

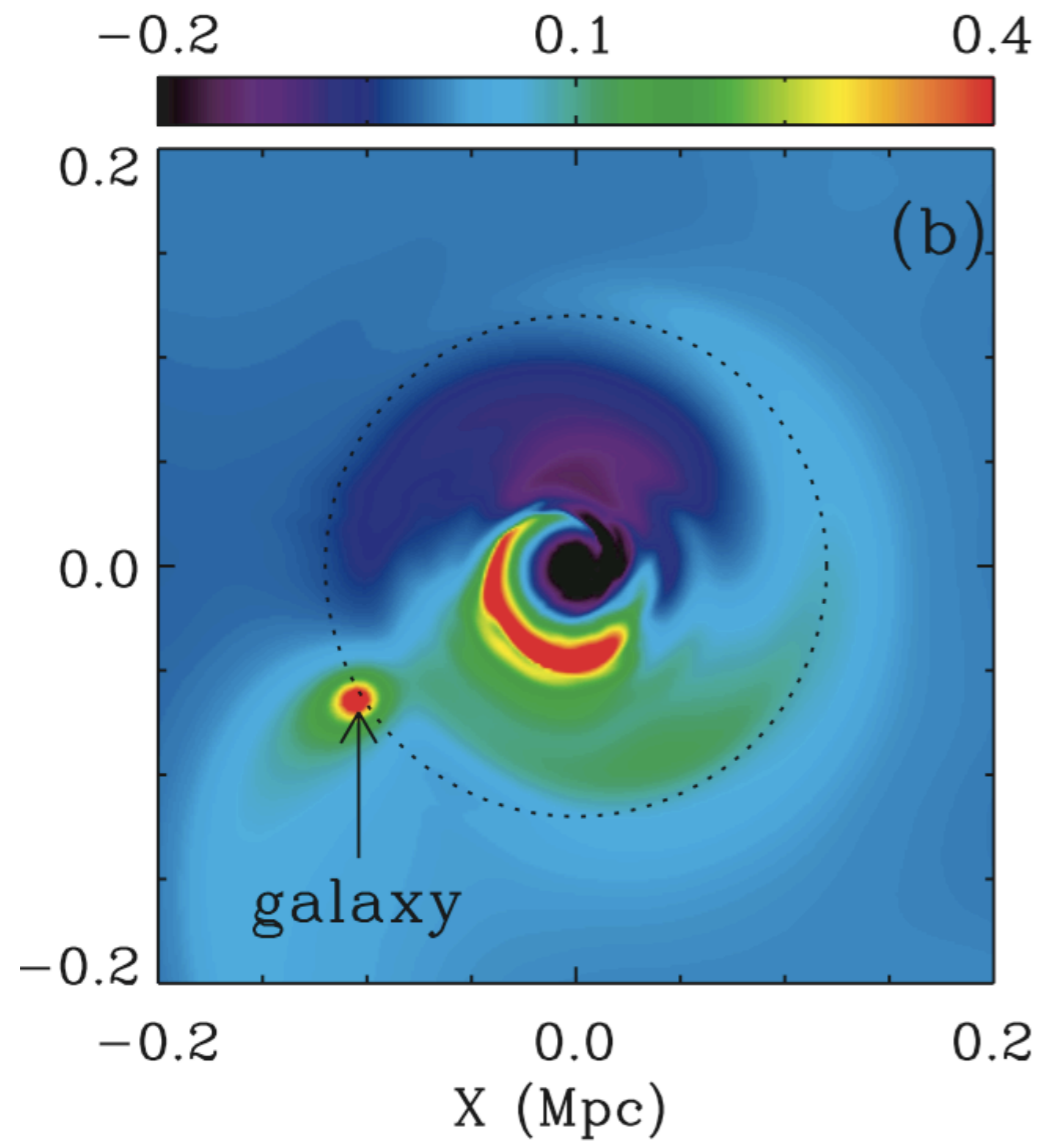


Mergers

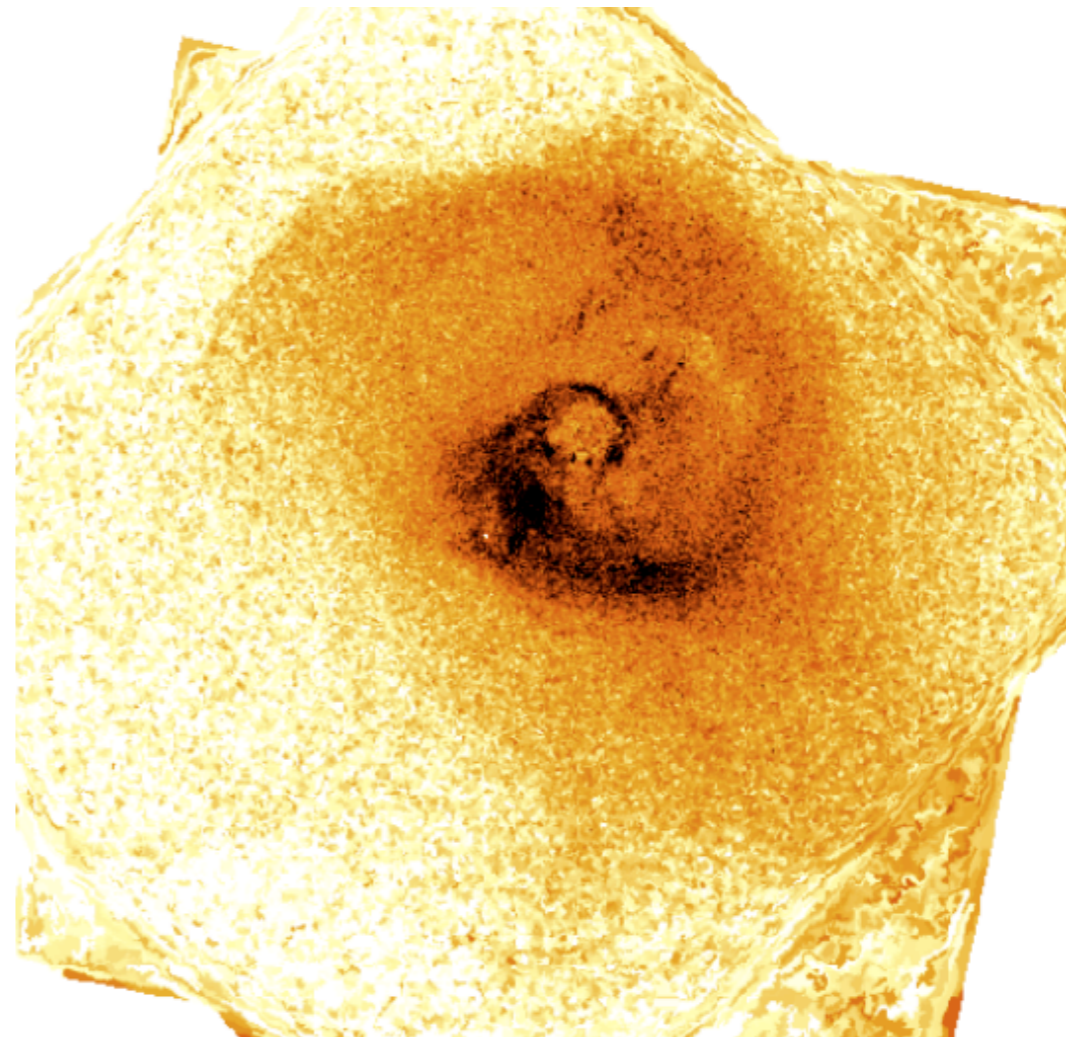


AGN Outbursts

# ...except galaxy motions



Kim (2007)



Fabian (2005)

# But galactic wakes have low volume-filling factor

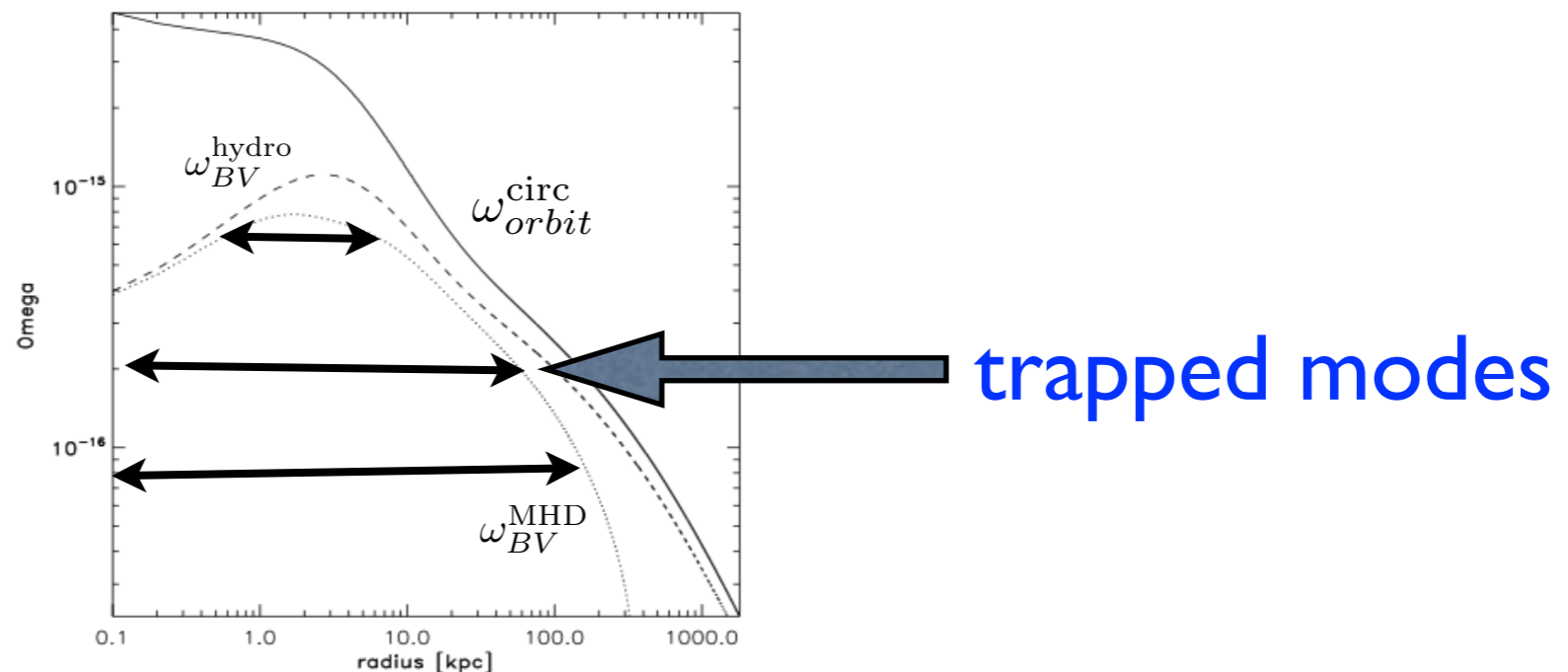
(Subramanian et al 2006, Iapichino & Niemeyer 2008)

## Volume-filling only if can excite g-modes

(Lufkin et al 1995, Kim 2007)

This requires  $\omega_{\text{turb}} \sim v_{\text{turb}}/\lambda < \omega_{\text{BV}}$

cD galaxy contribution to potential is crucial





# Contradictory requirements?

Trapping of g-modes requires

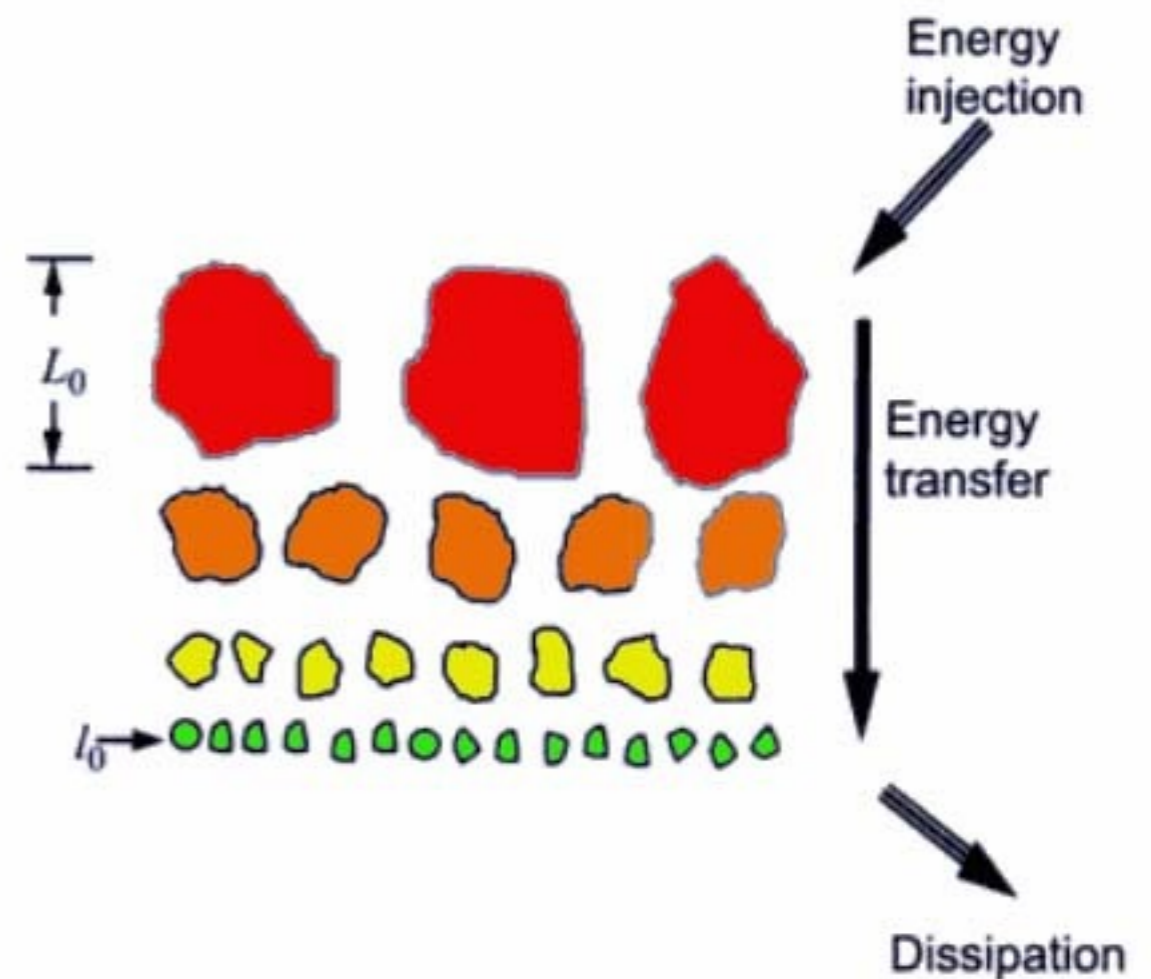
$$\omega_{\text{turb}} < \omega_{\text{BV}}$$

Isotropic velocities requires

$$\omega_{\text{turb}} > \omega_{\text{BV}}$$

But recall  $\varepsilon \sim \frac{v^3}{l} \sim \text{const}$

$$\omega \sim \frac{v}{l} \sim l^{-2/3}$$



Large eddies are trapped, smaller eddies isotropize



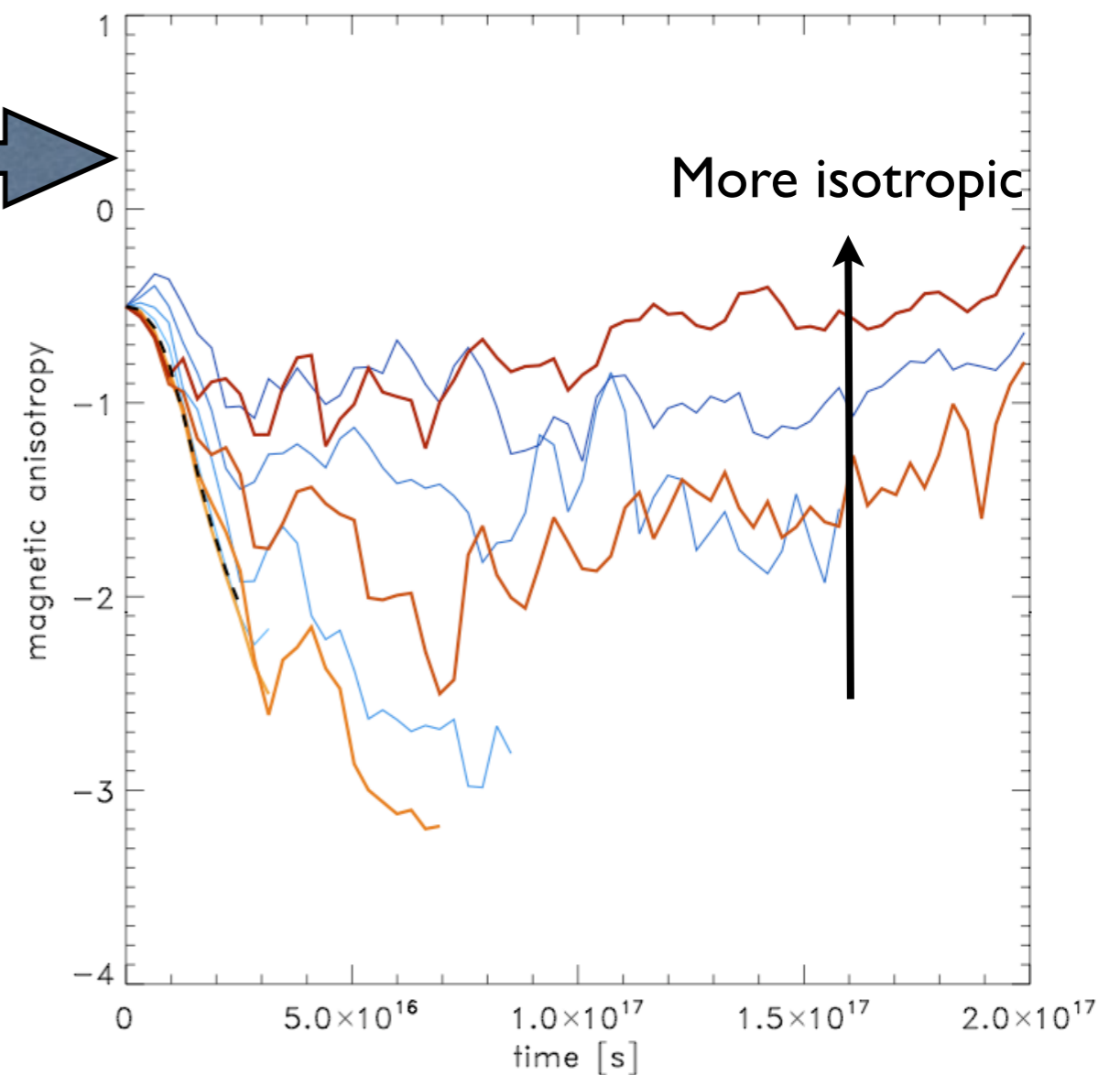
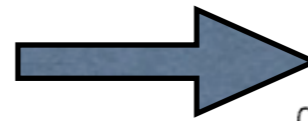
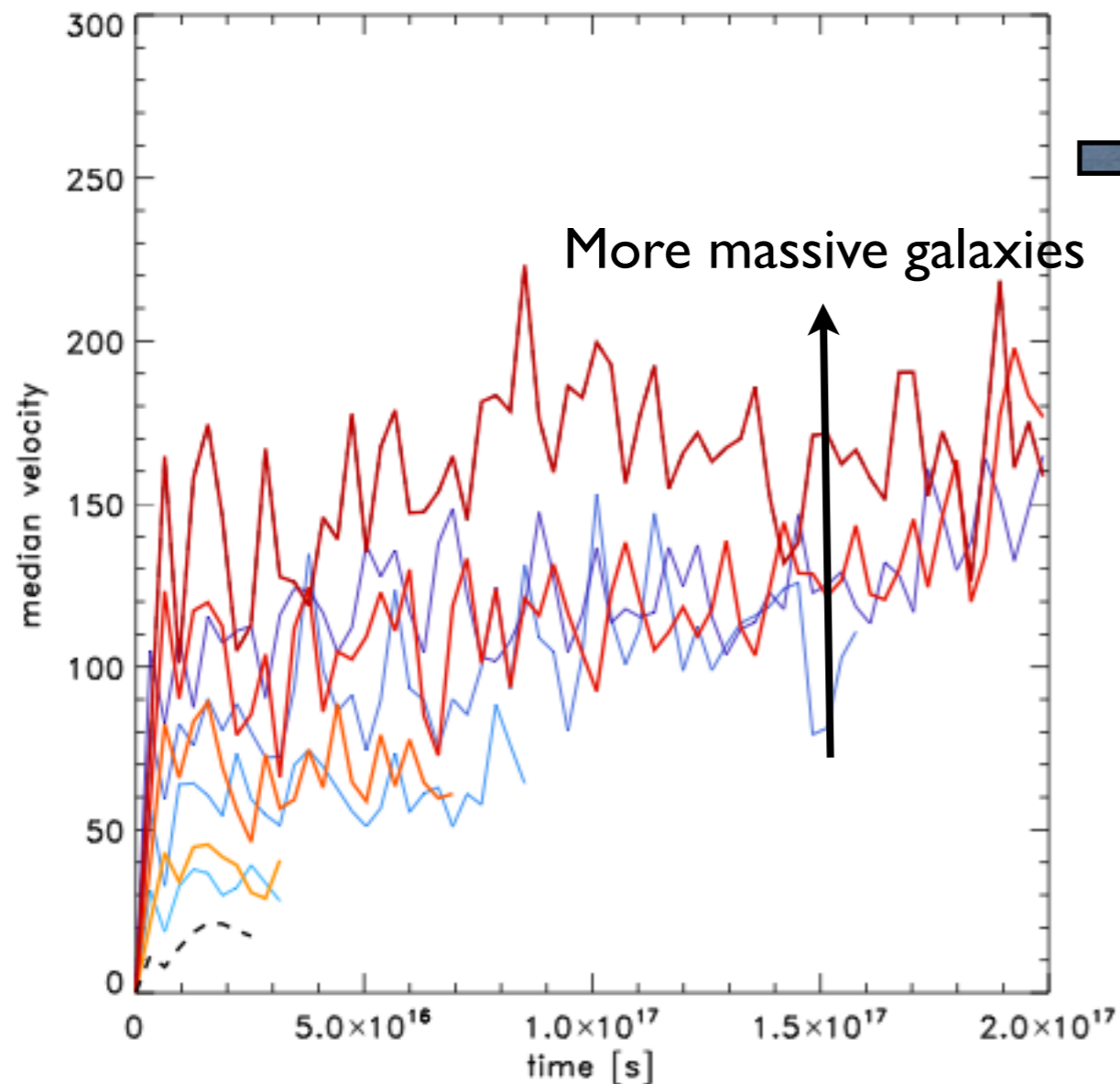
# Simulate this!



Assume galaxies trace NFW profile; solve Jeans equation to get velocities

Mass in substructure from lensing studies (Natarajan et al 2008)

# Magnetic fields are isotropized...

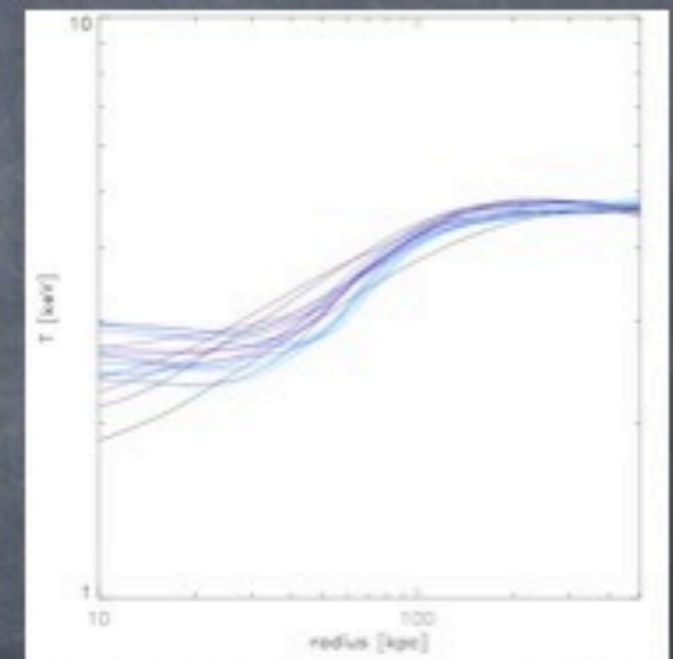
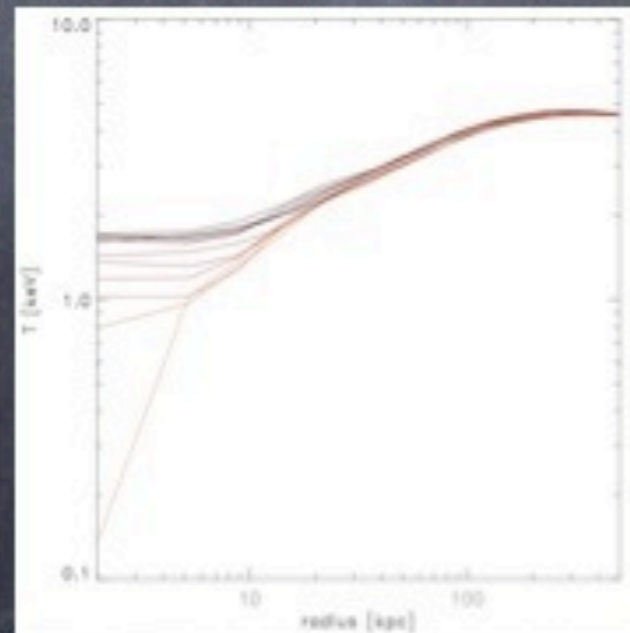
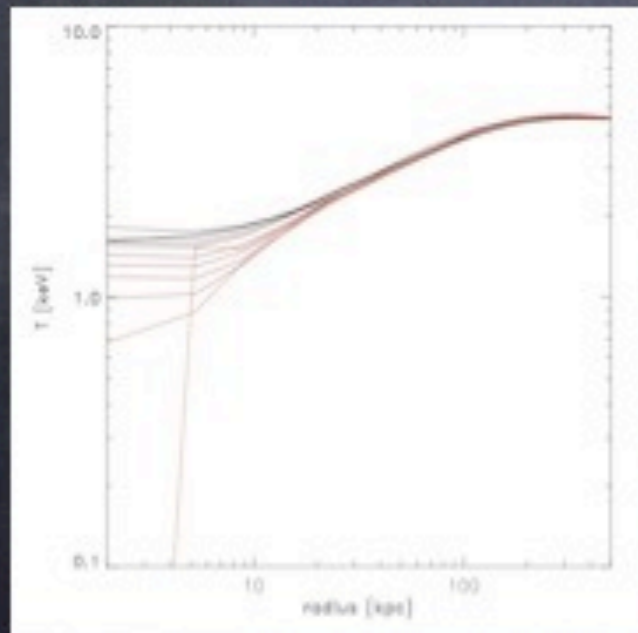
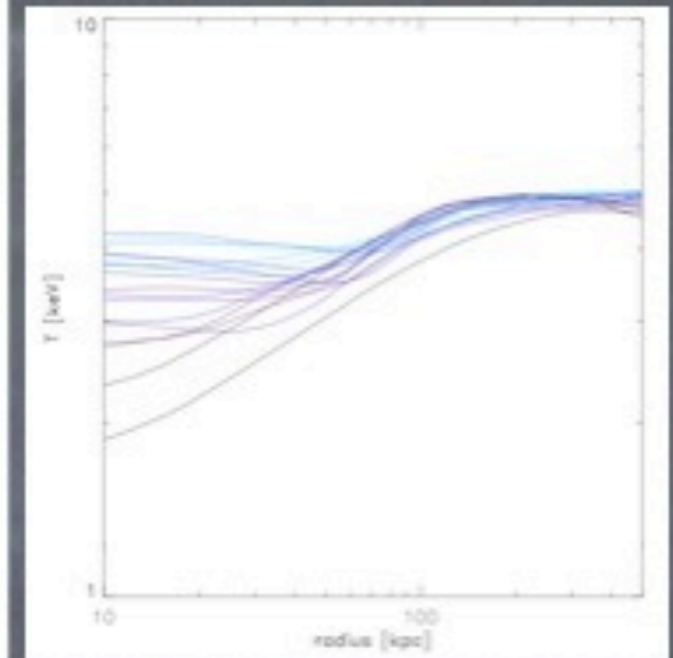
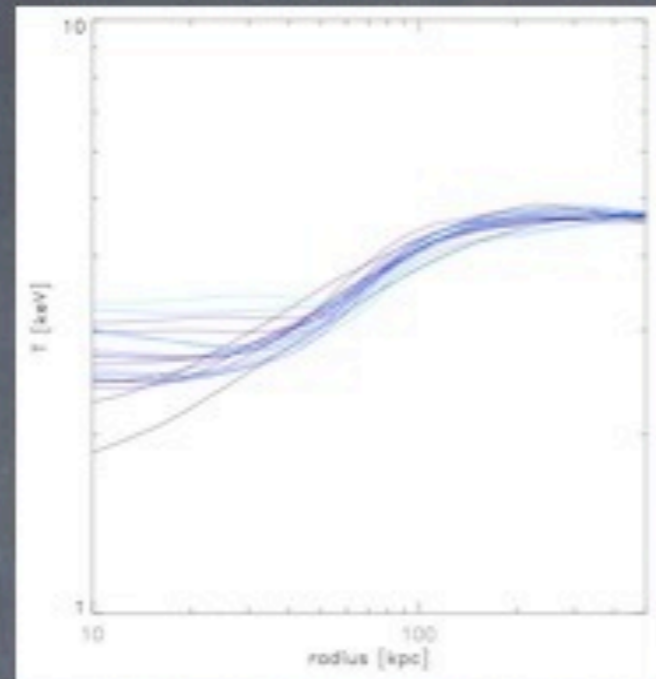
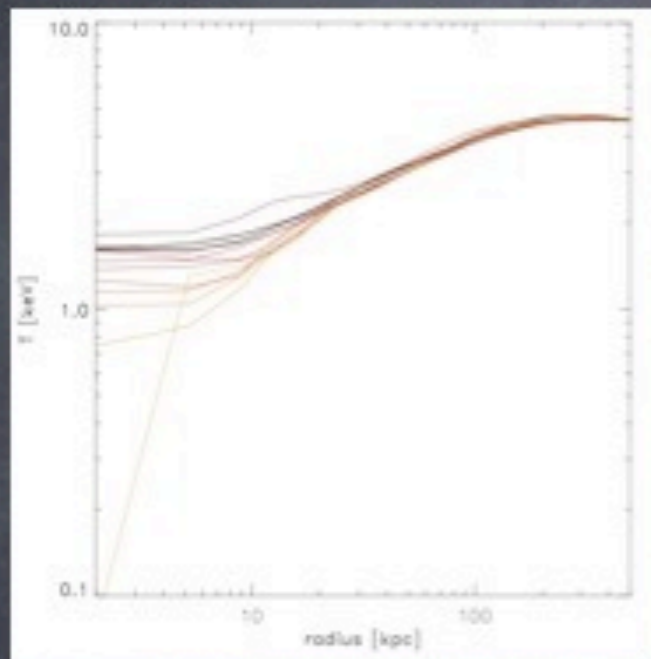




Note: no subgrid physics!

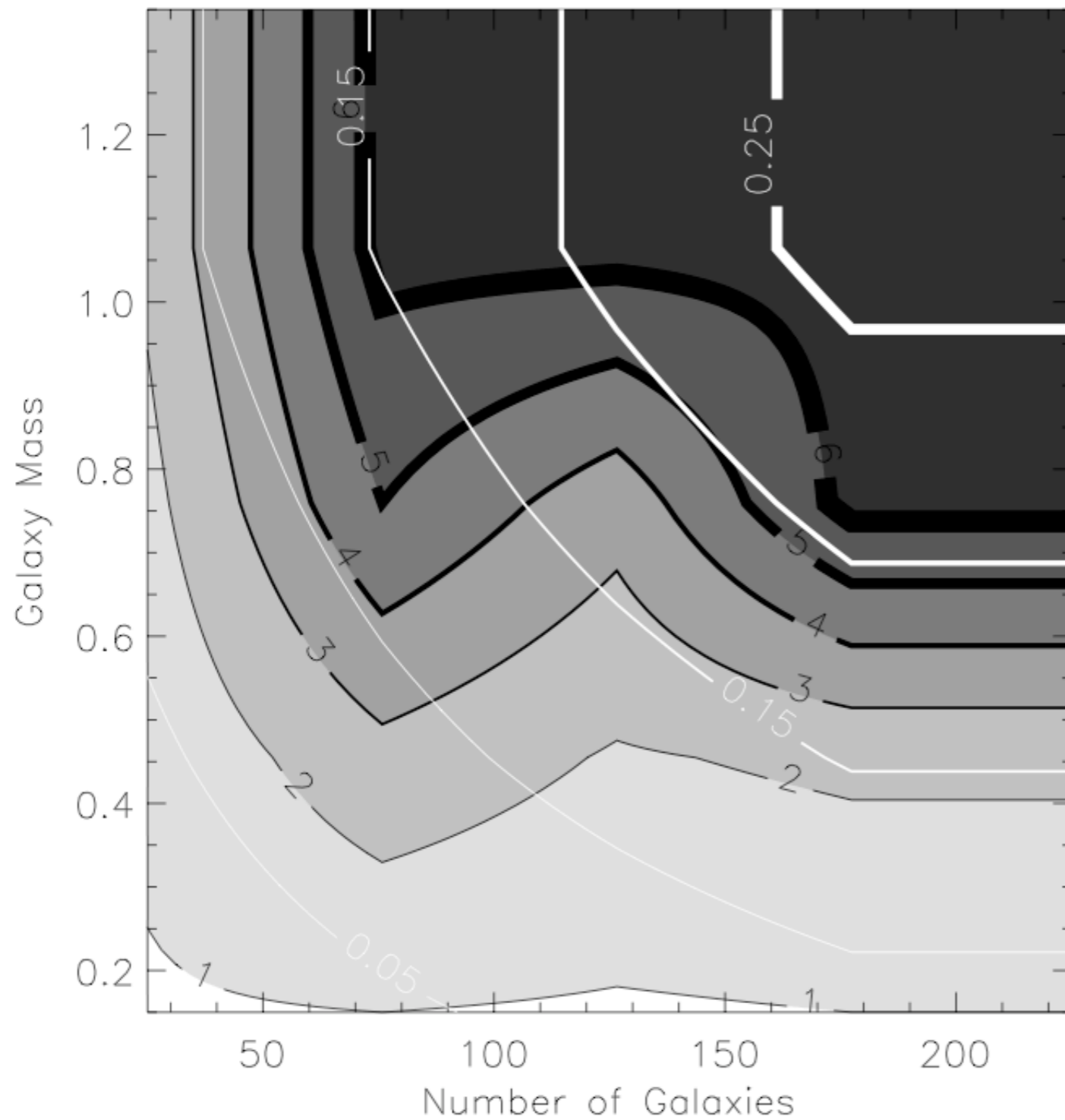


Galaxy mass



Number of galaxies



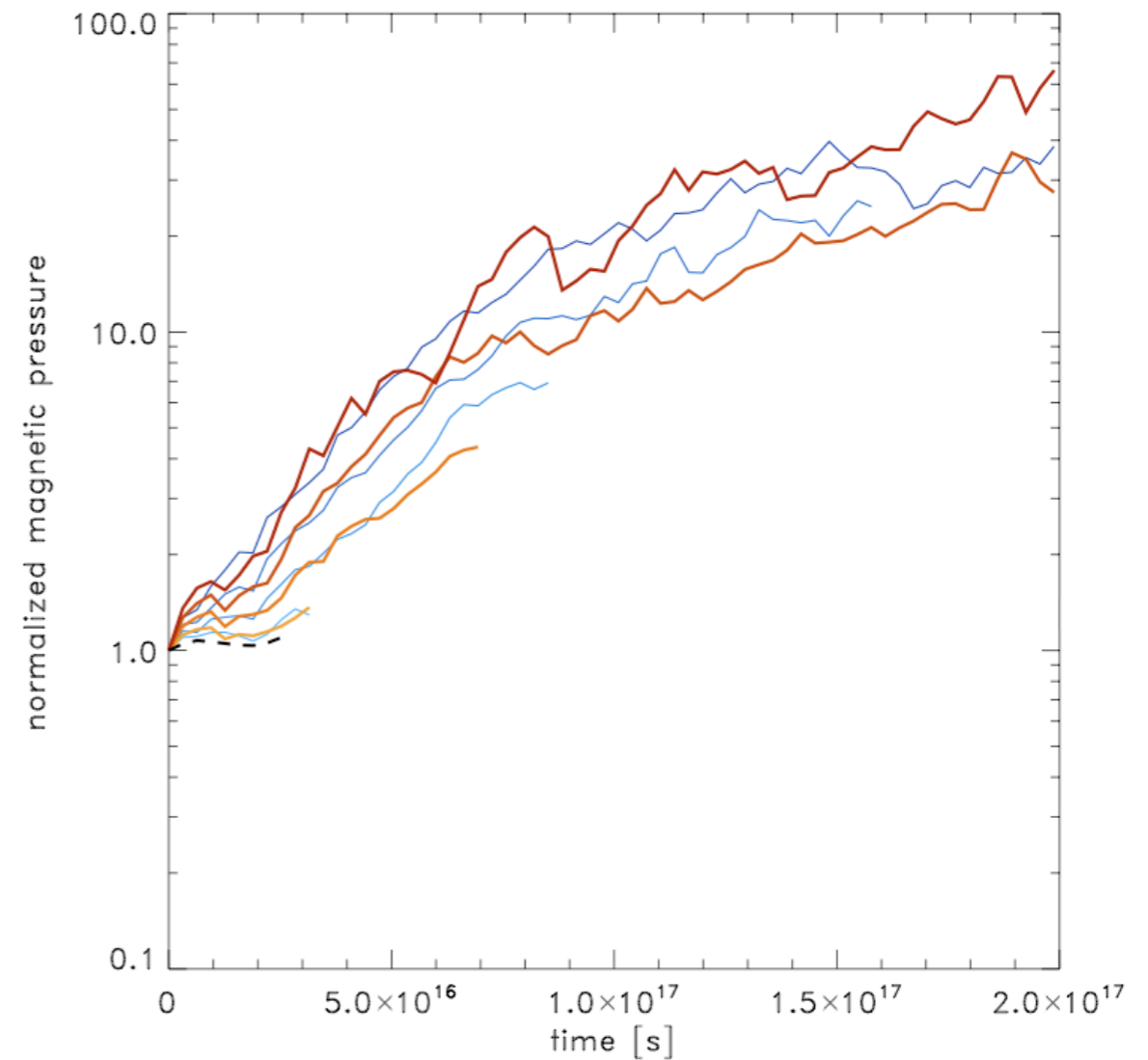
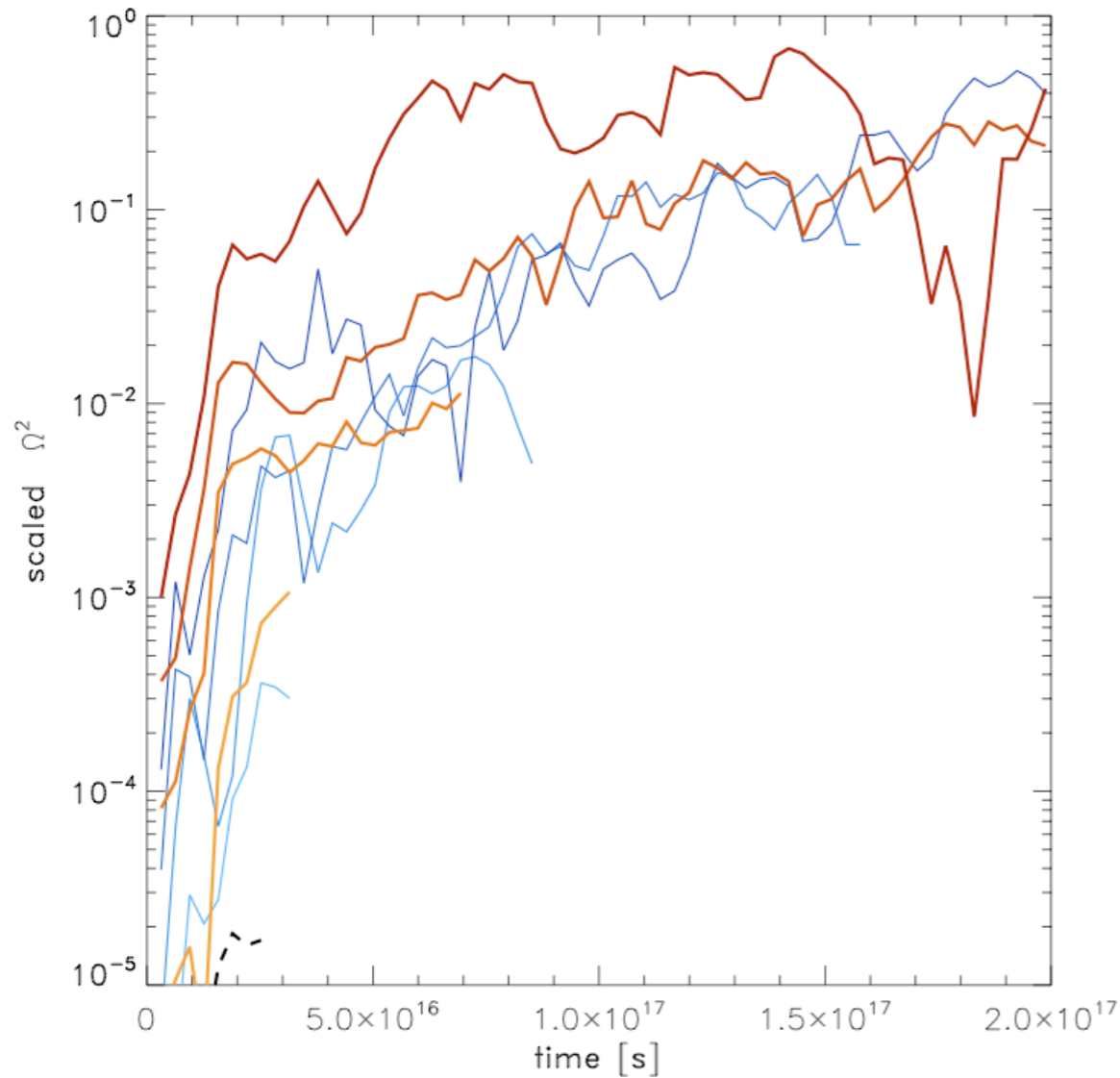


Black contours: Time to cooling catastrophe (6 Gyr=end of sim, stable)

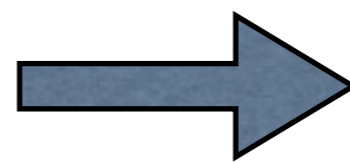
White contours: fraction of mass in substructure

**No cooling catastrophe for plausible parameters consistent with lensing**

# Magnetic fields are amplified



Vorticity growth

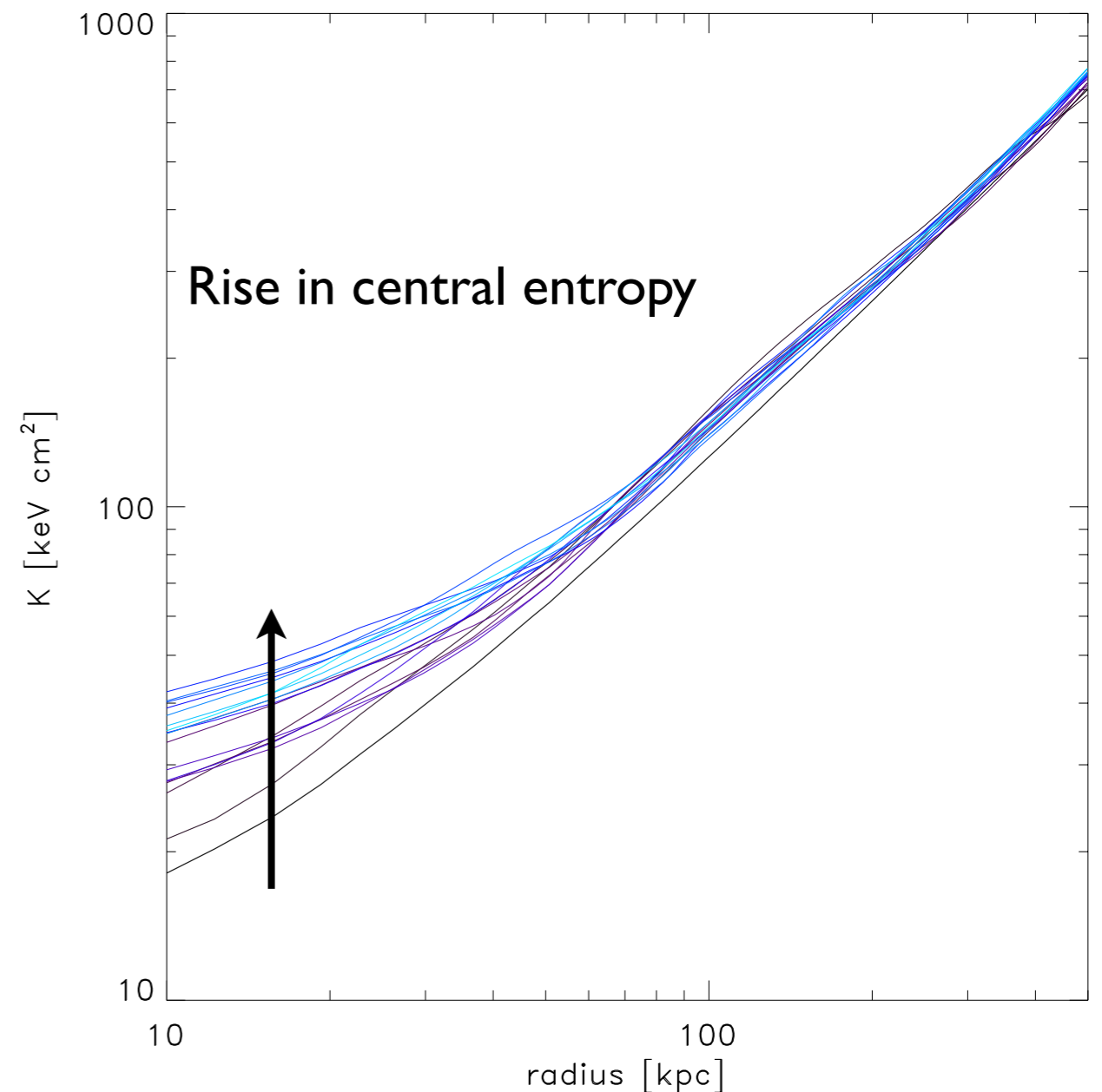
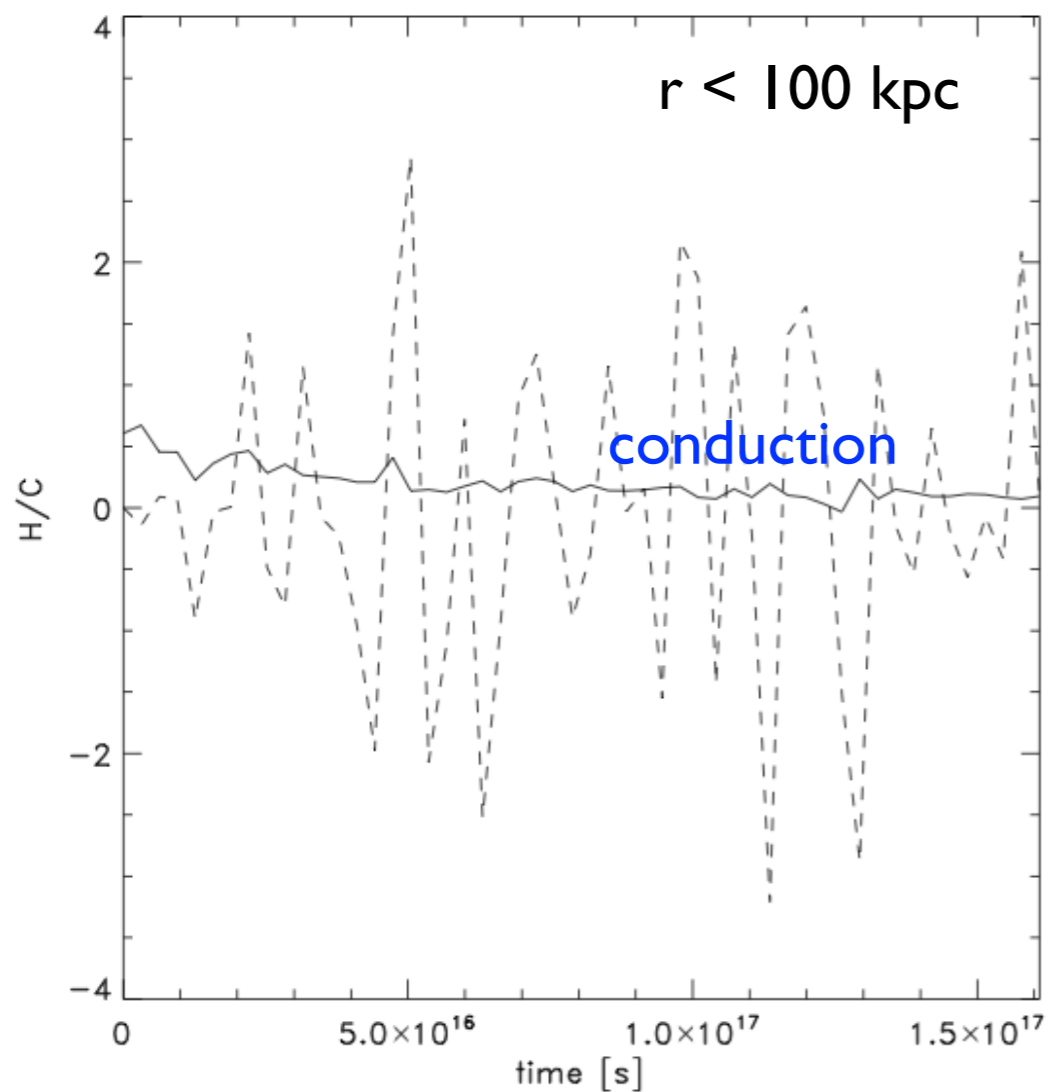


B-field growth

N.B. Magnetic tension negligible here

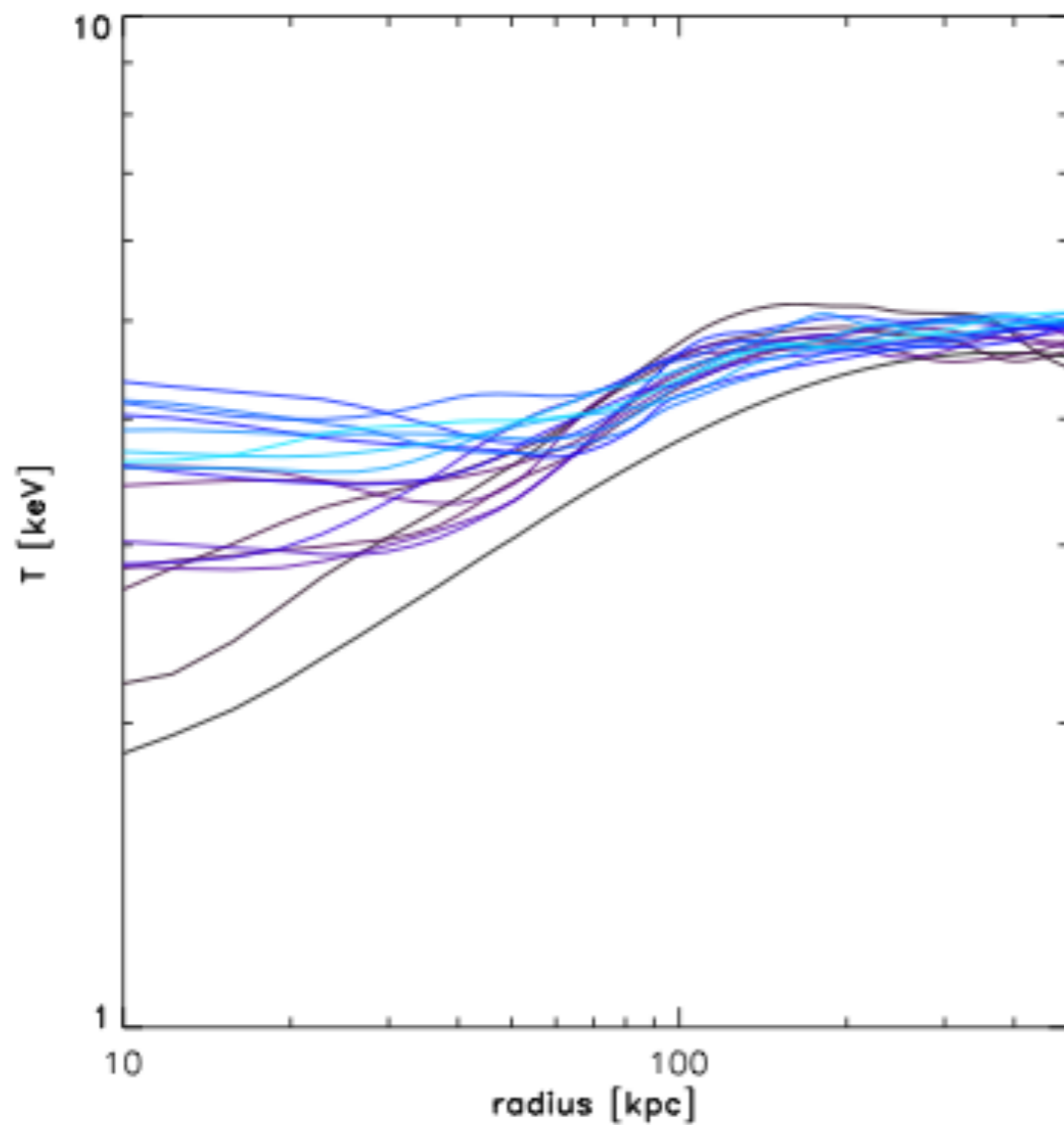


# Conduction is at most ~50% of energy budget

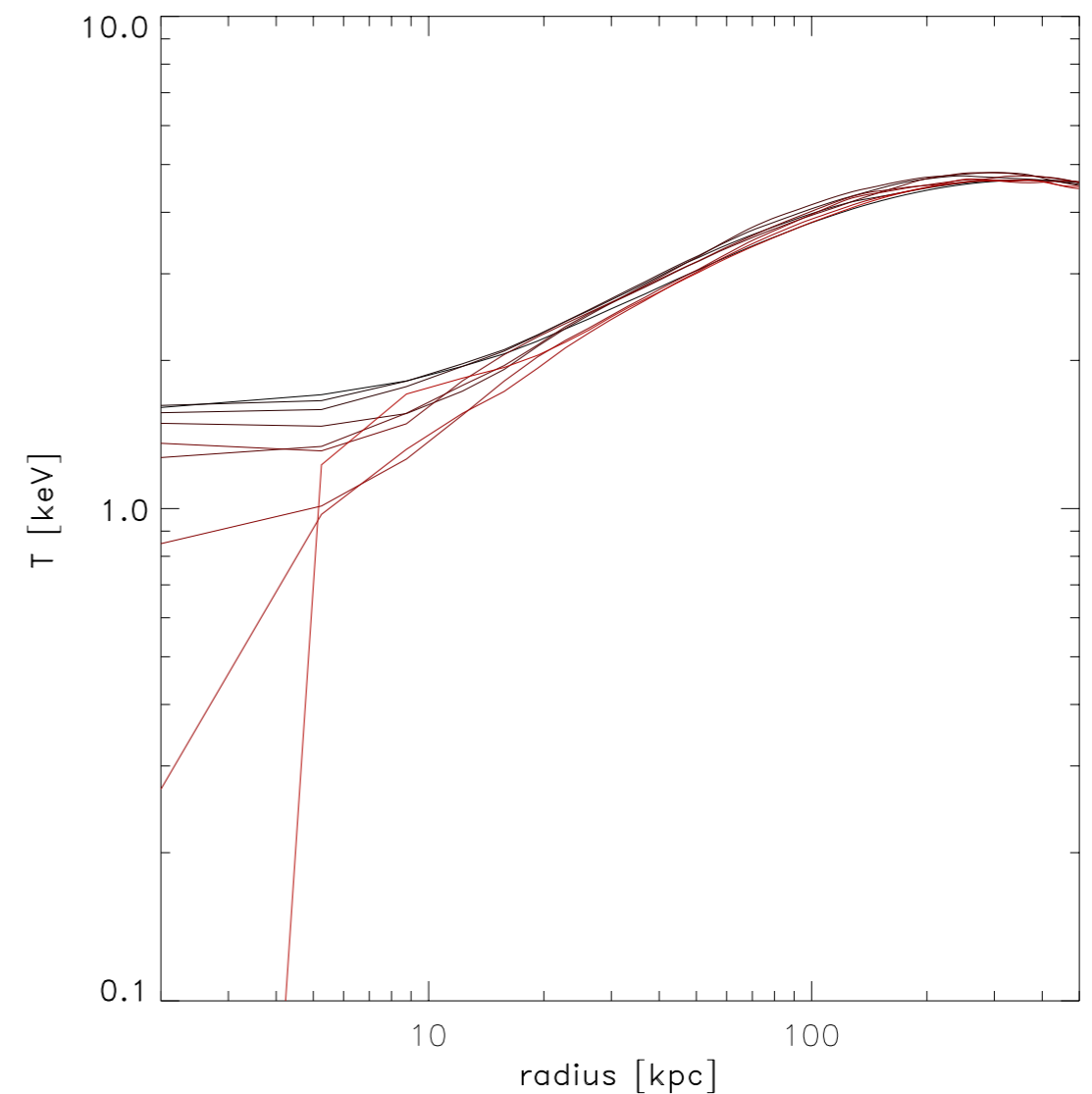


# ...but is a critical part of the story

MHD: **stable**

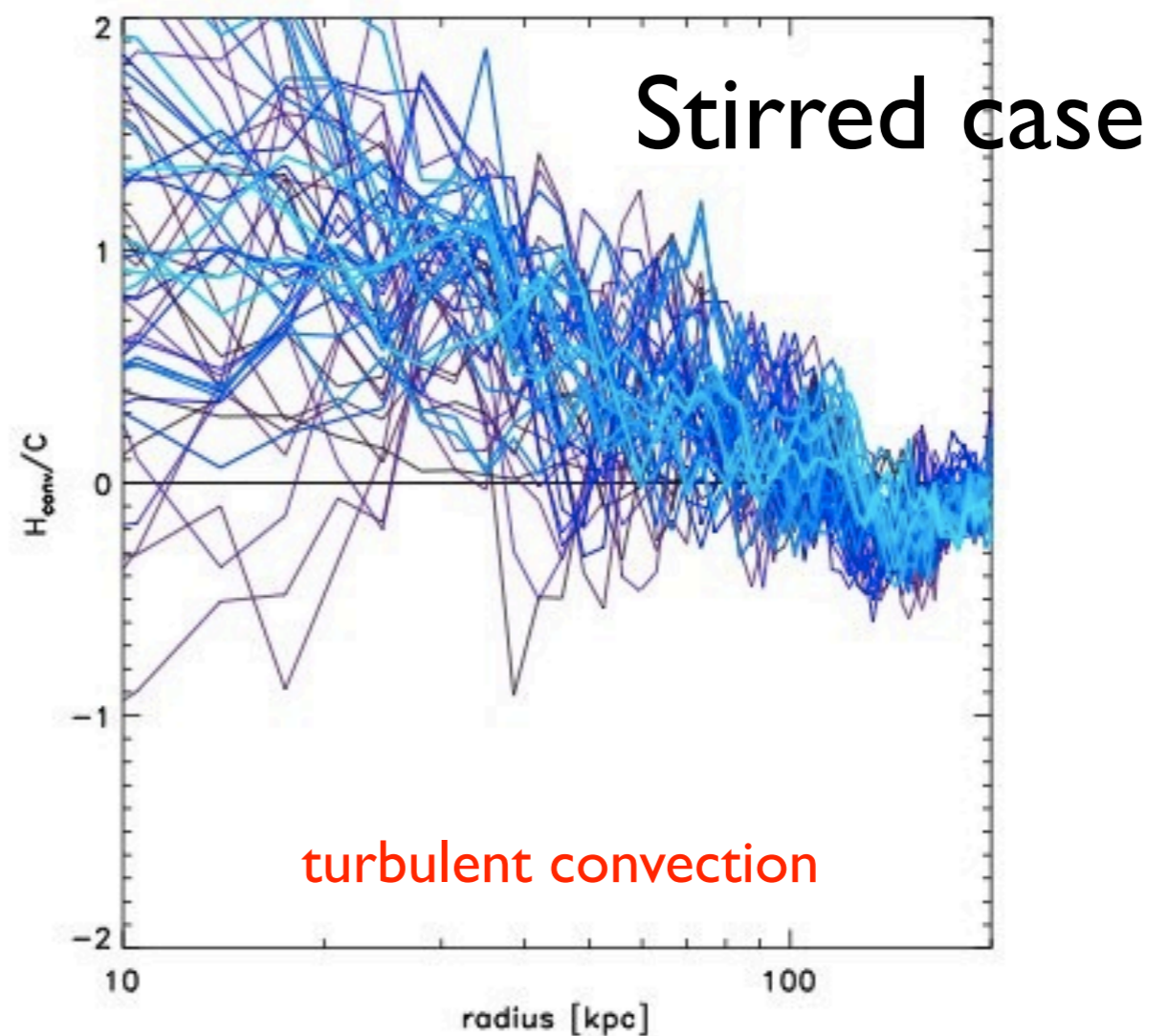
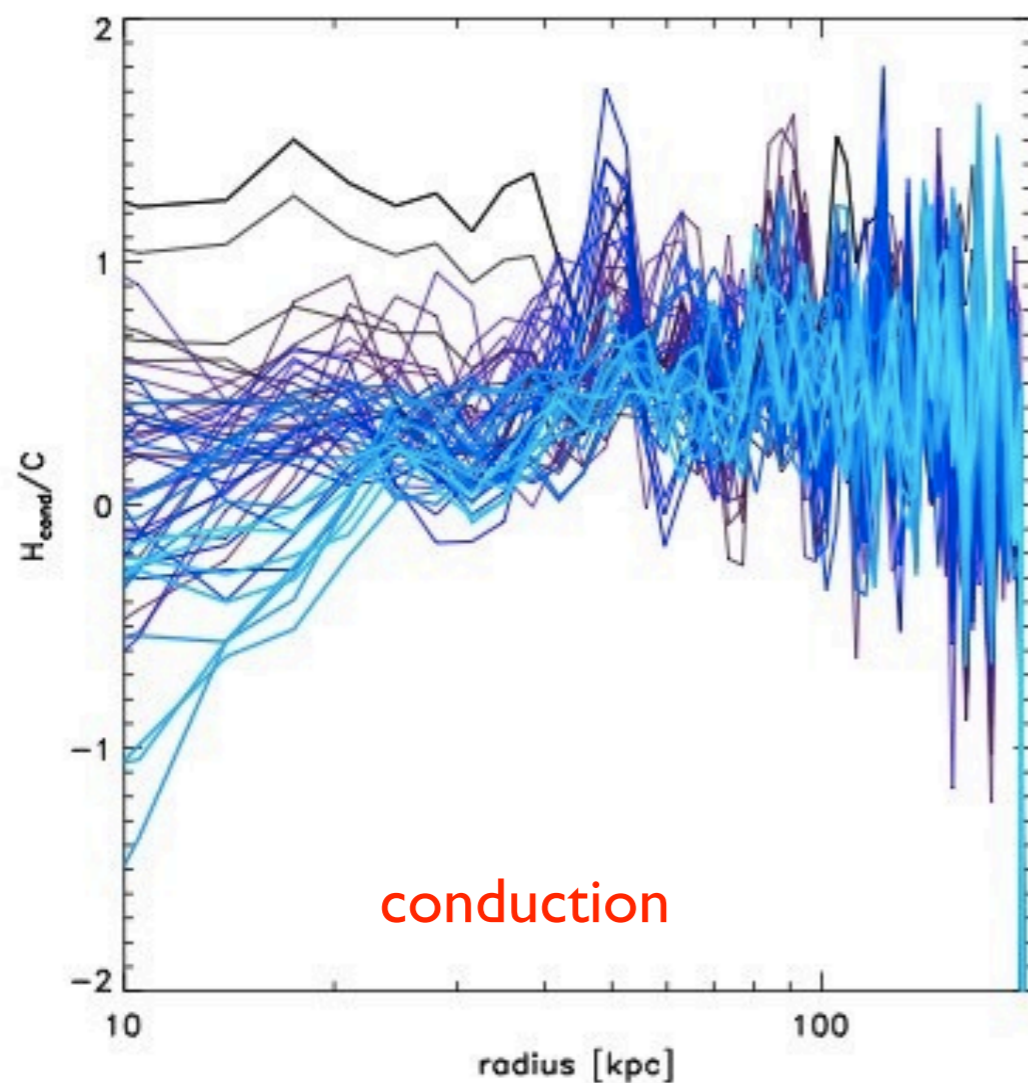


Pure hydro: **cooling catastrophe**



Likely that conduction **facilitates** turbulent diffusion

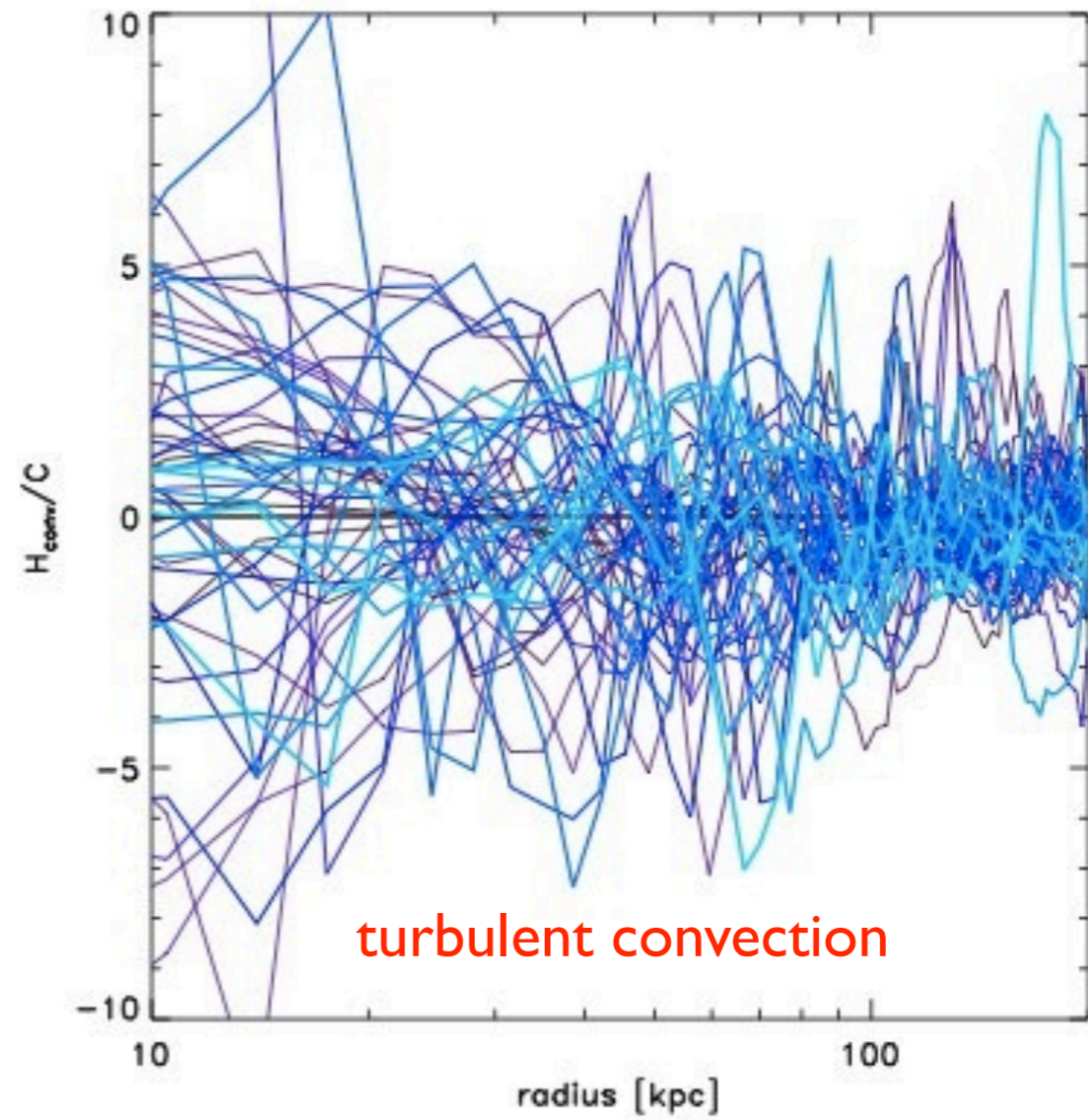
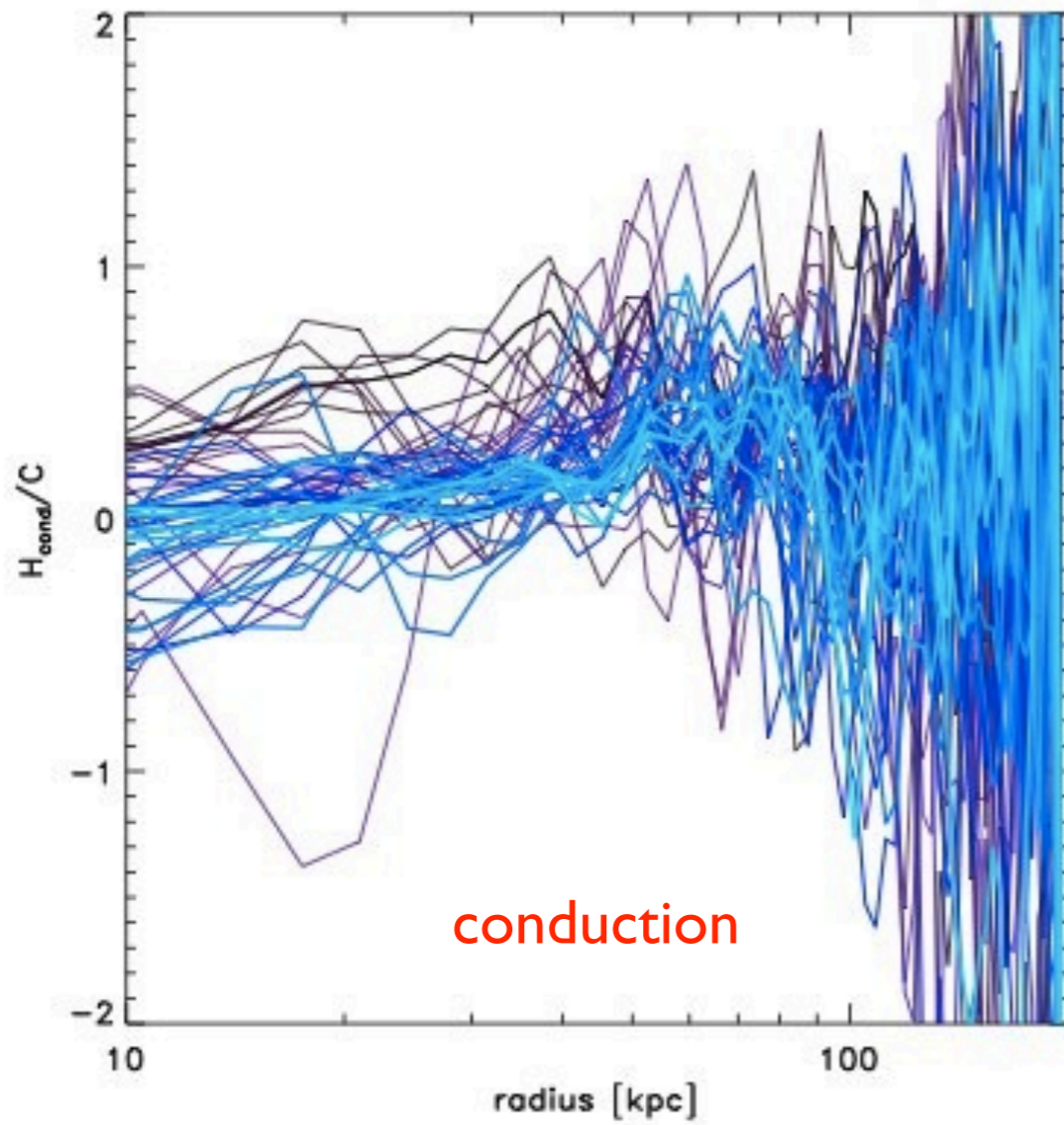
# Conduction is more important in outer regions



Turbulent heat transport more important in inner regions

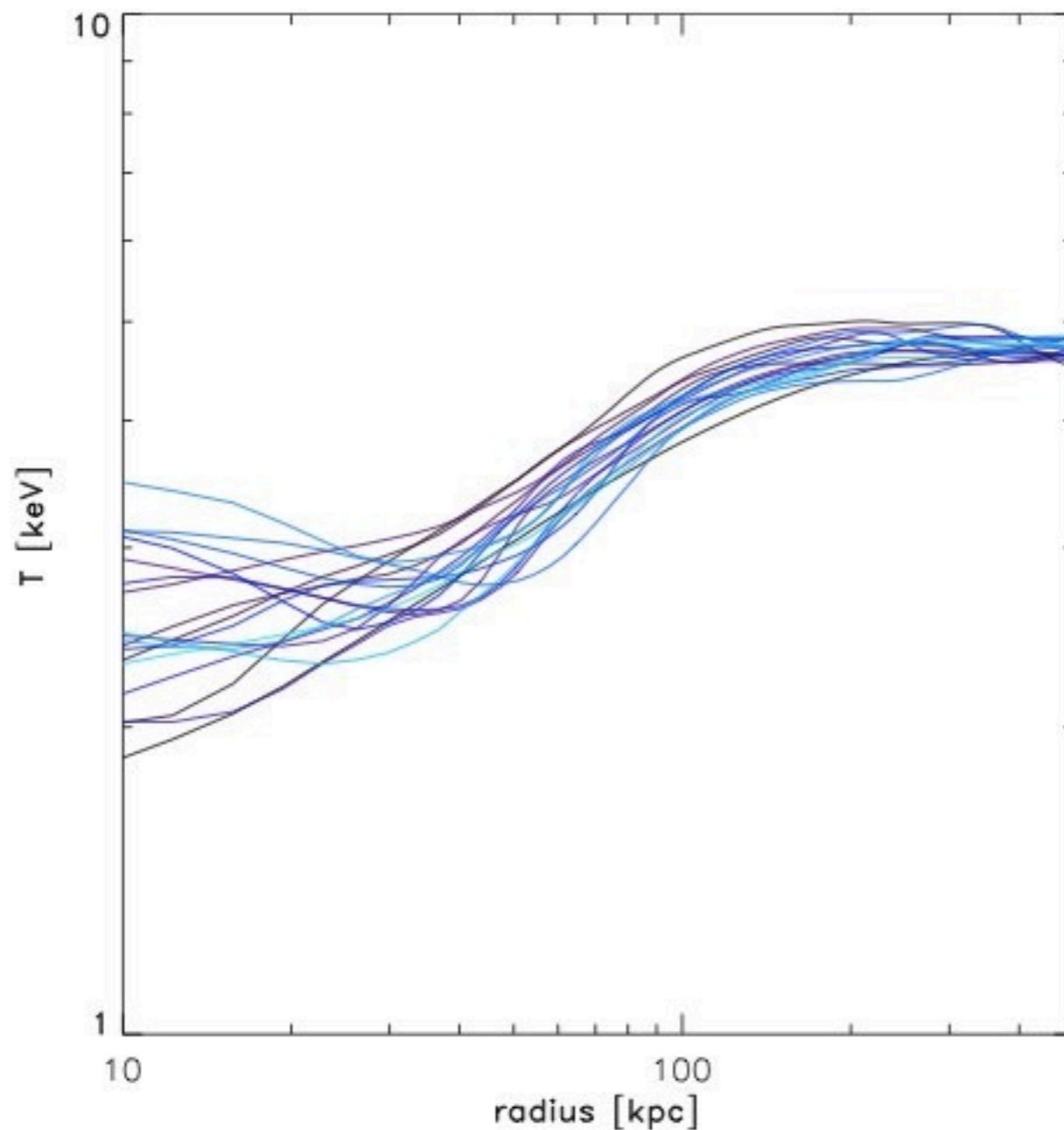


# galaxies case





# So conduction-only solutions don't apply



Stable solutions don't have to be isothermal

Can have inverted temperature profiles near the center

# Dissipation of turbulent motions is negligible

$$\Gamma = \frac{c_{\text{diss}} \rho v_t^3}{l} = \frac{c_{\text{diss}} U_t}{t_{\text{edd}}}$$

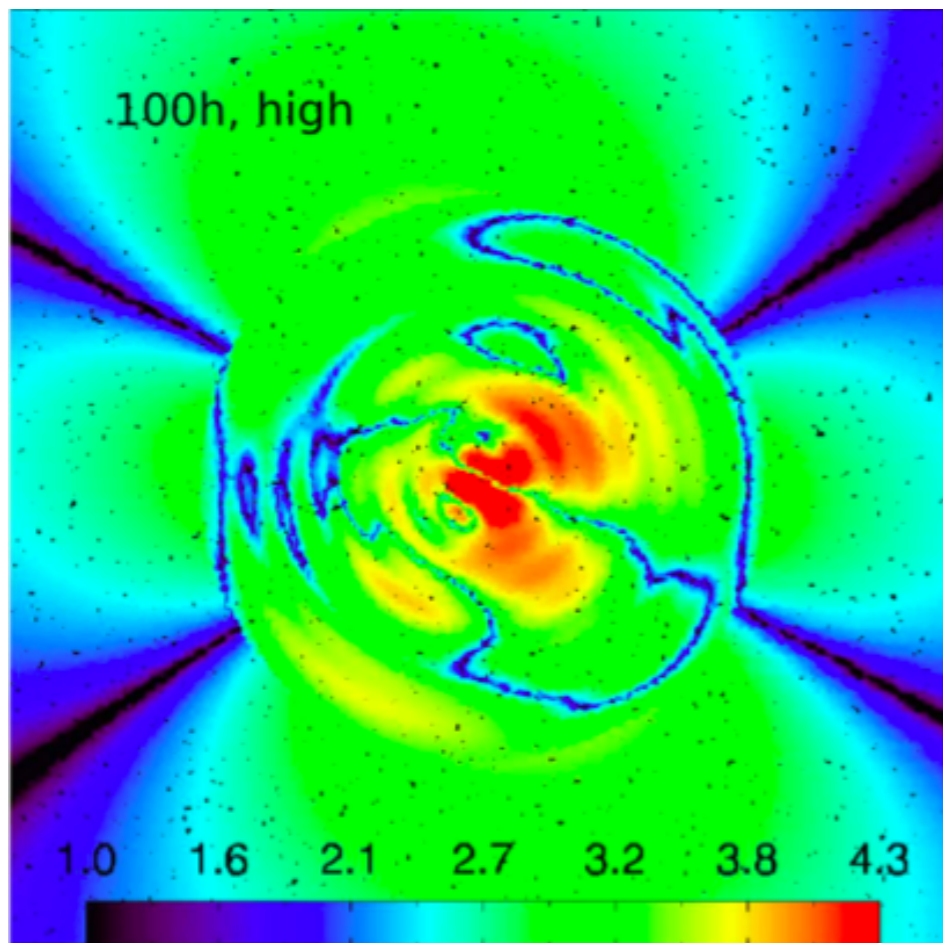
$$t_{\text{edd}} \sim t_{\text{cool}} \Rightarrow t_{\text{heat}} = \frac{U_{\text{thermal}}}{\Gamma} \approx \frac{U_{\text{thermal}}}{U_{\text{turb}}} t_{\text{cool}} \sim 100 t_{\text{cool}}$$

Dynamical friction heating unimportant in these simulations

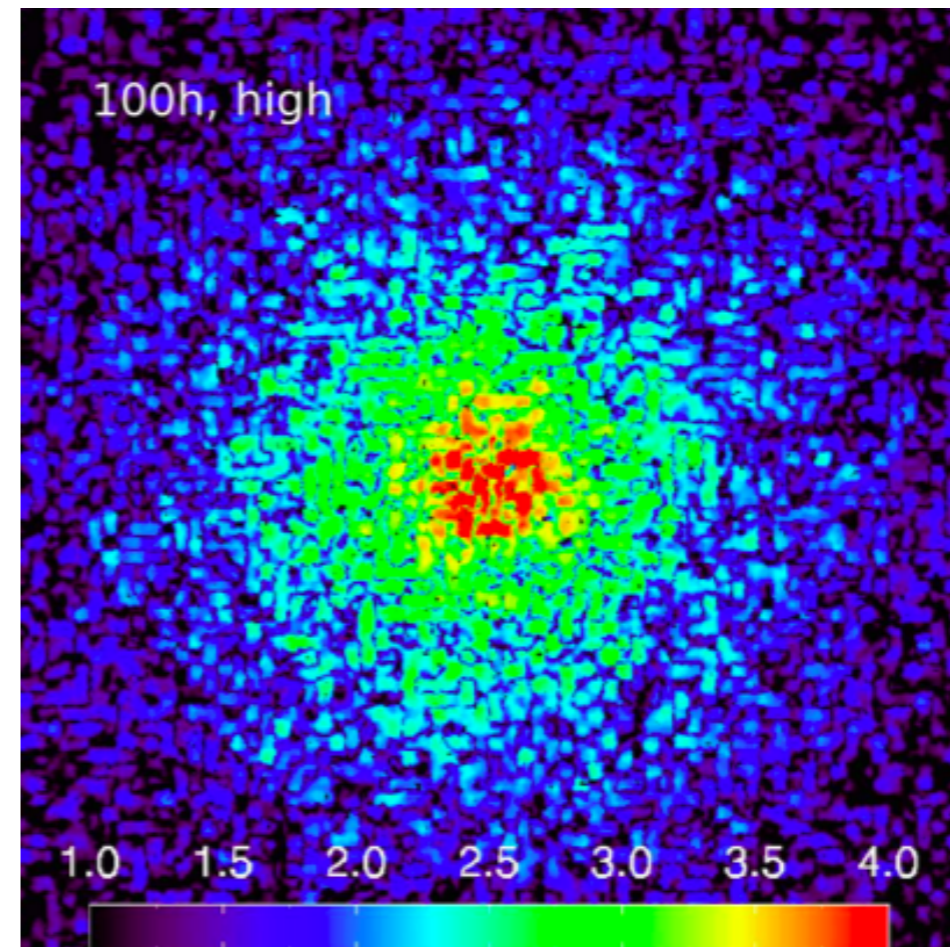
# Implications

# Faraday rotation could test field geometries

HBI/MTI dominated



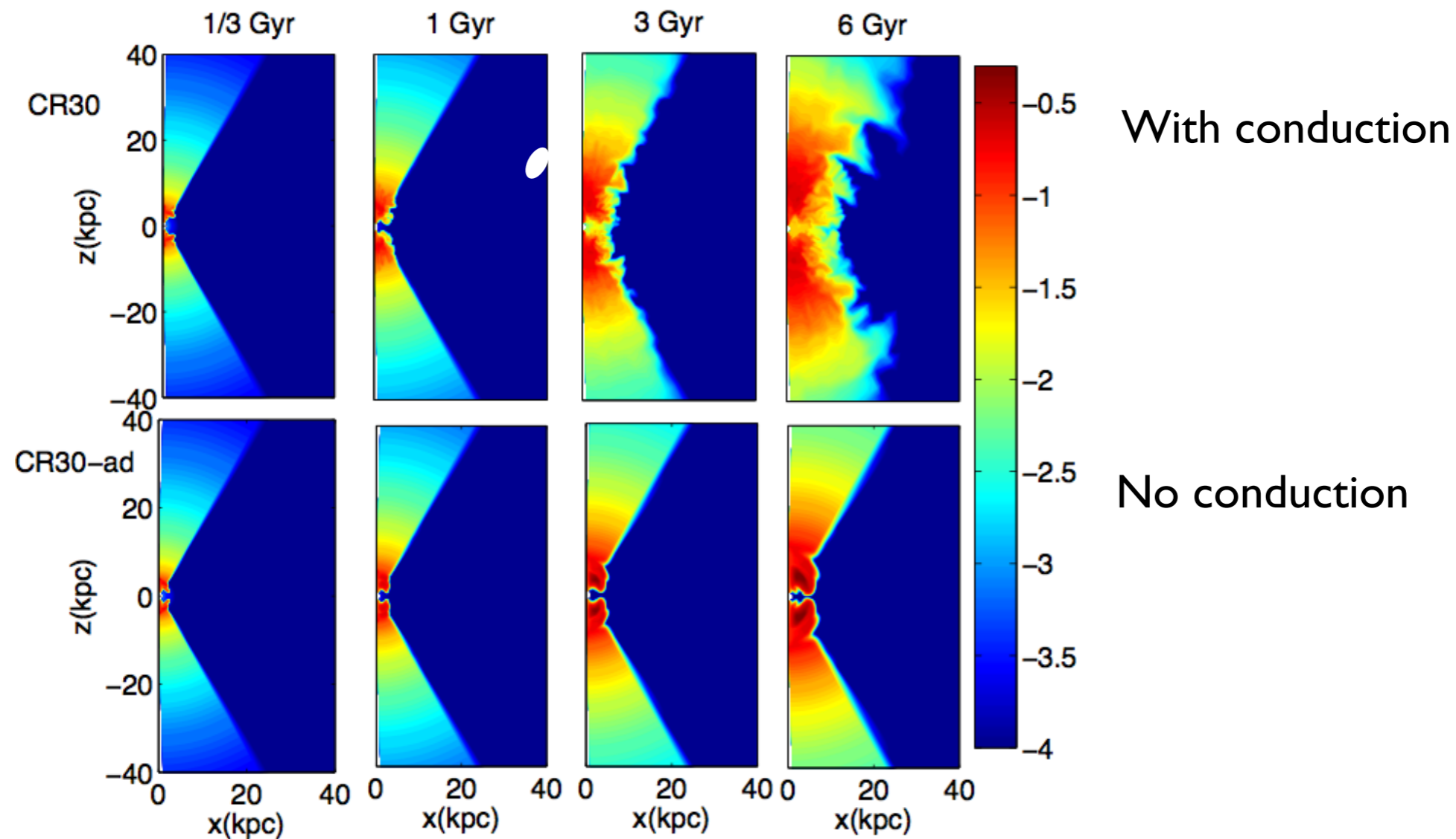
Turbulence dominated



Bogdanovic et al (2010)



# Metal diffusion much more efficient with conduction



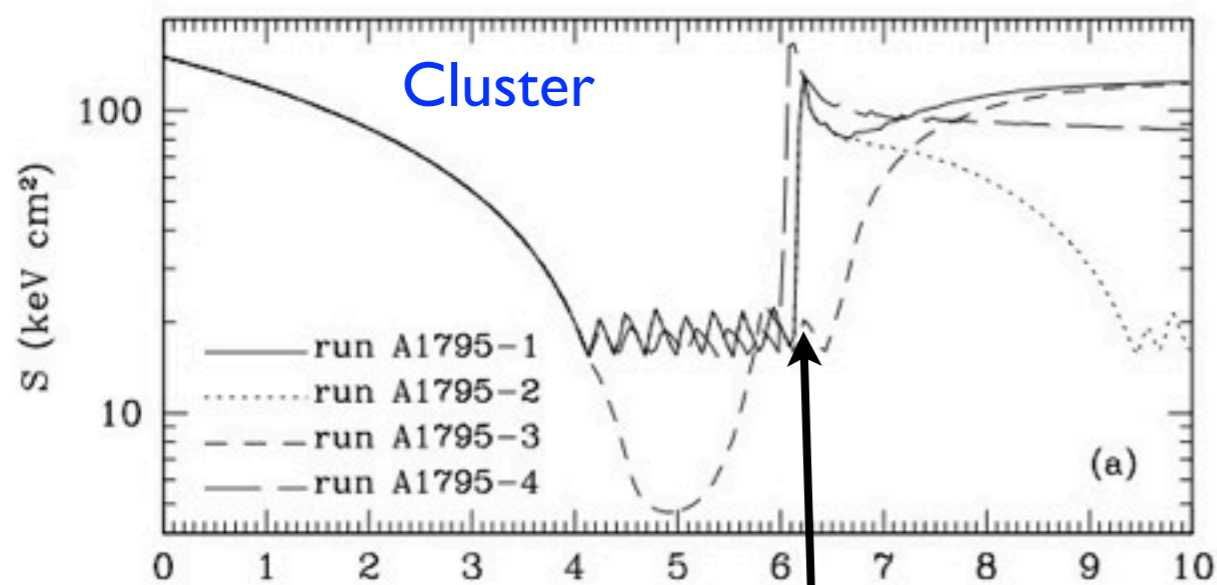
Sharma et al (2009)

# Turbulence could regulate the CC to NCC transition...

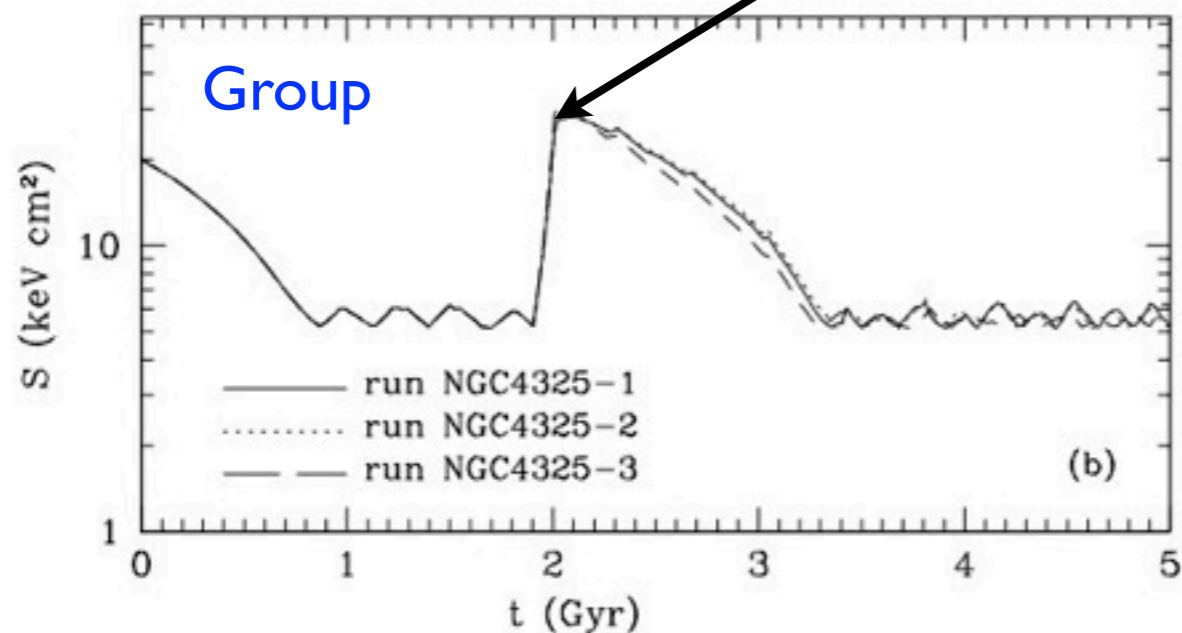
AGN or mergers could stimulate turbulence and change conductivity. **Big impact on cluster profiles, even if their energetic contribution small**

If this is right, should not see bimodality in groups

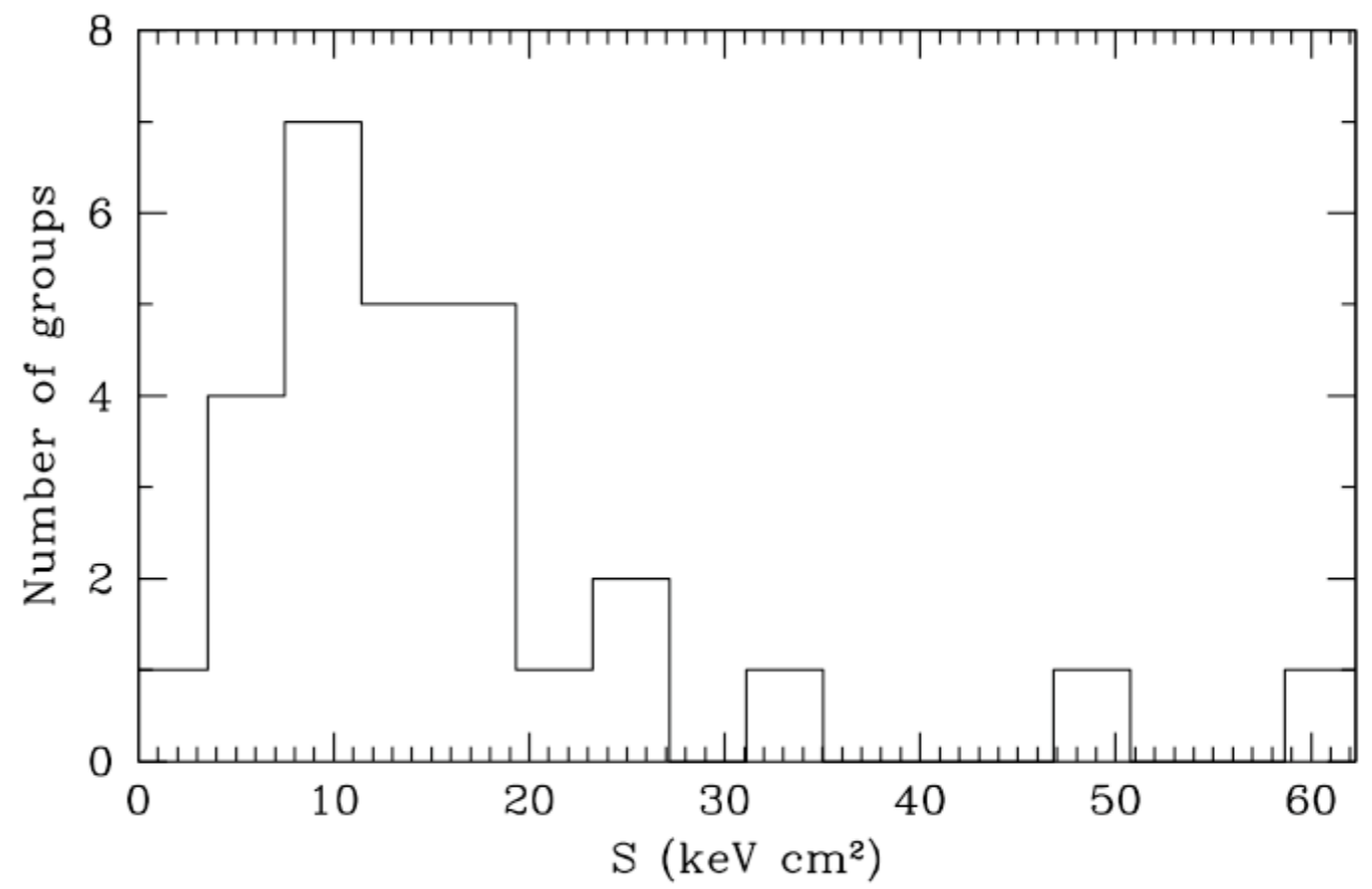
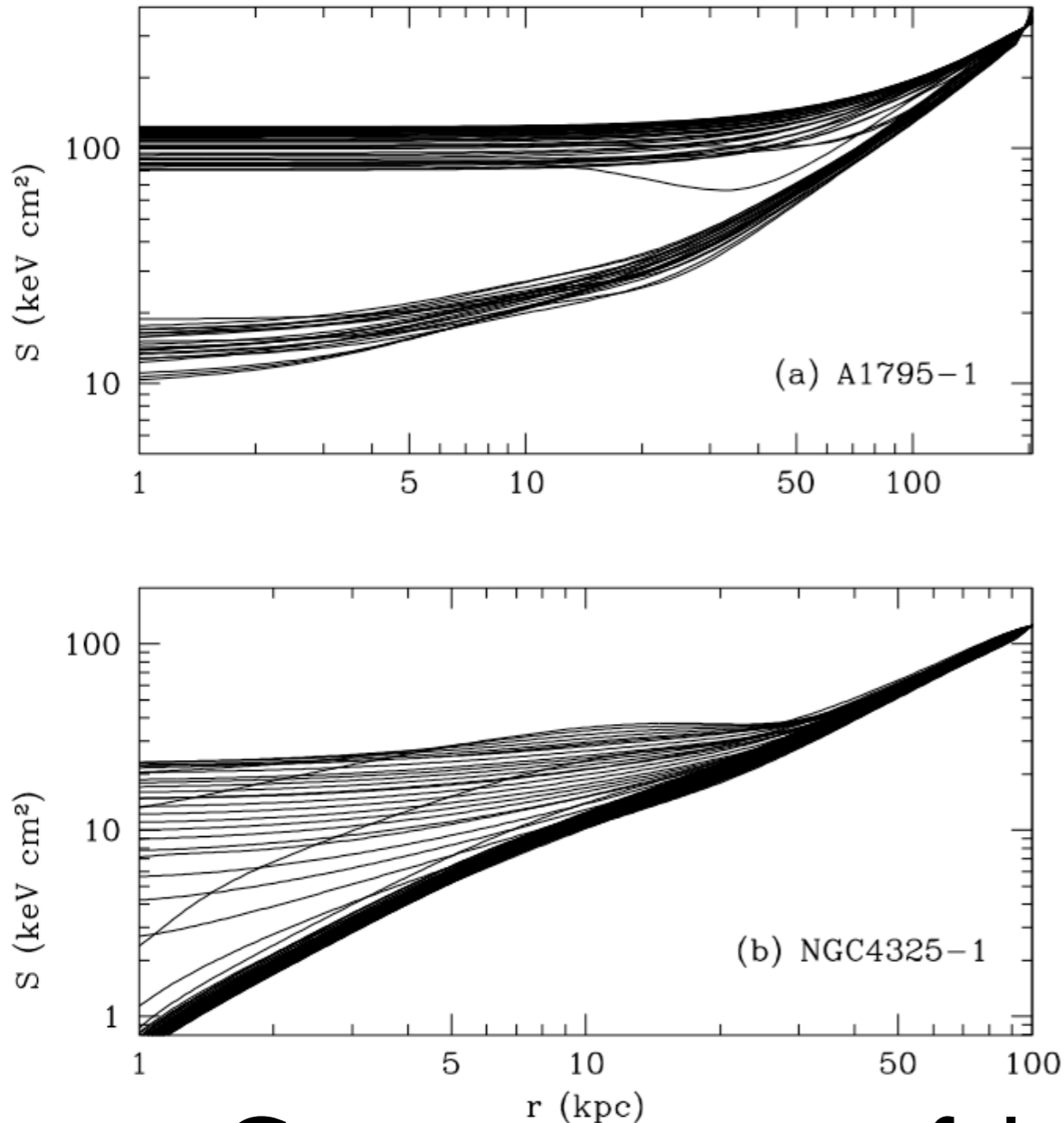
Guo & Oh (2009)



Sudden increase in conduction



Guo & Oh (2009)

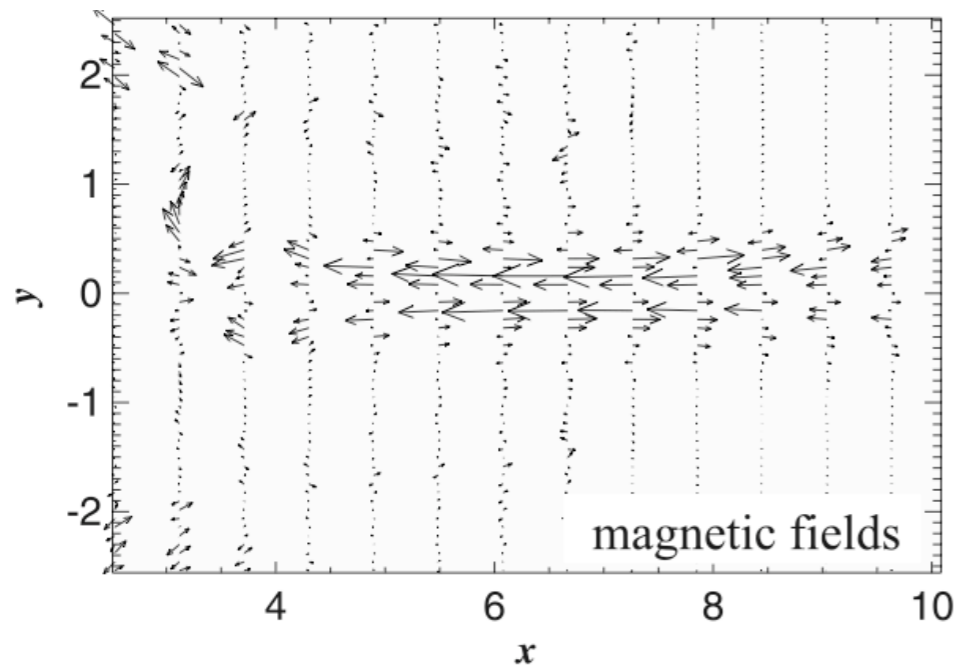


(group data from Johnson et al 2009)

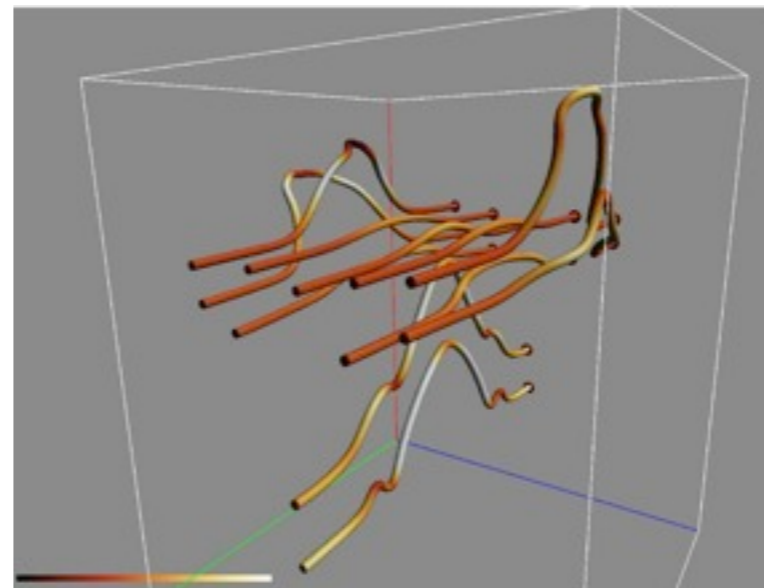
**Groups are a useful test: cannot be stabilized by conduction, we predict no bimodality**

**...although role of turbulent heat transport in groups needs to be investigated**

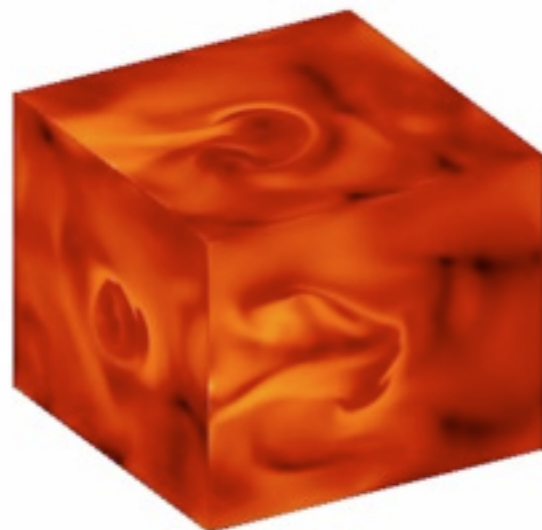
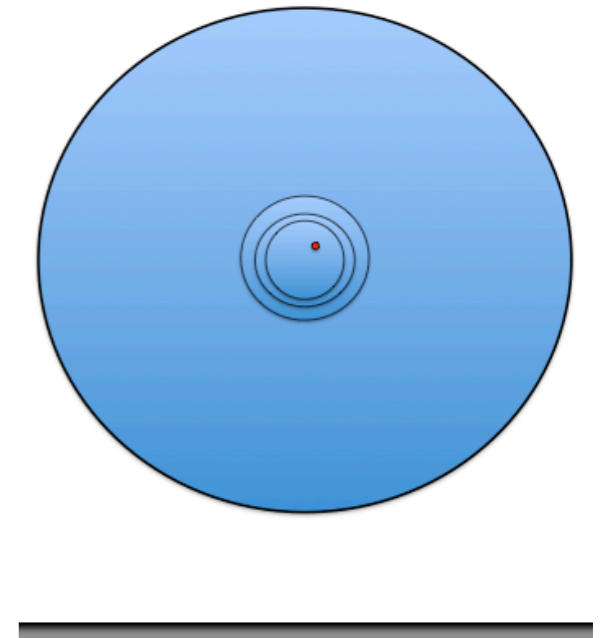
# Could rising bubbles straighten field lines?



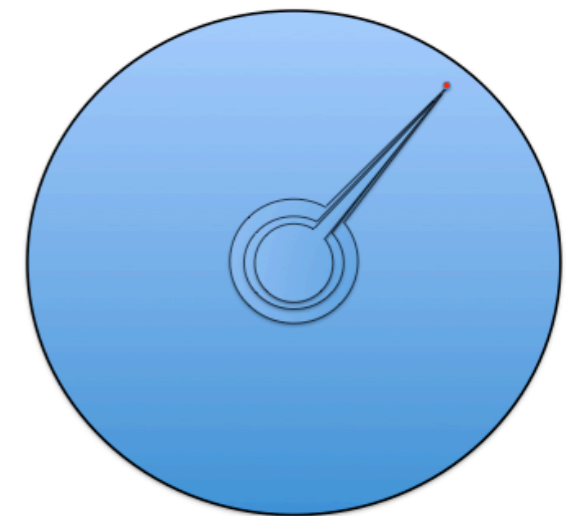
Asai et al (2007)



O'Neill et al (2009)

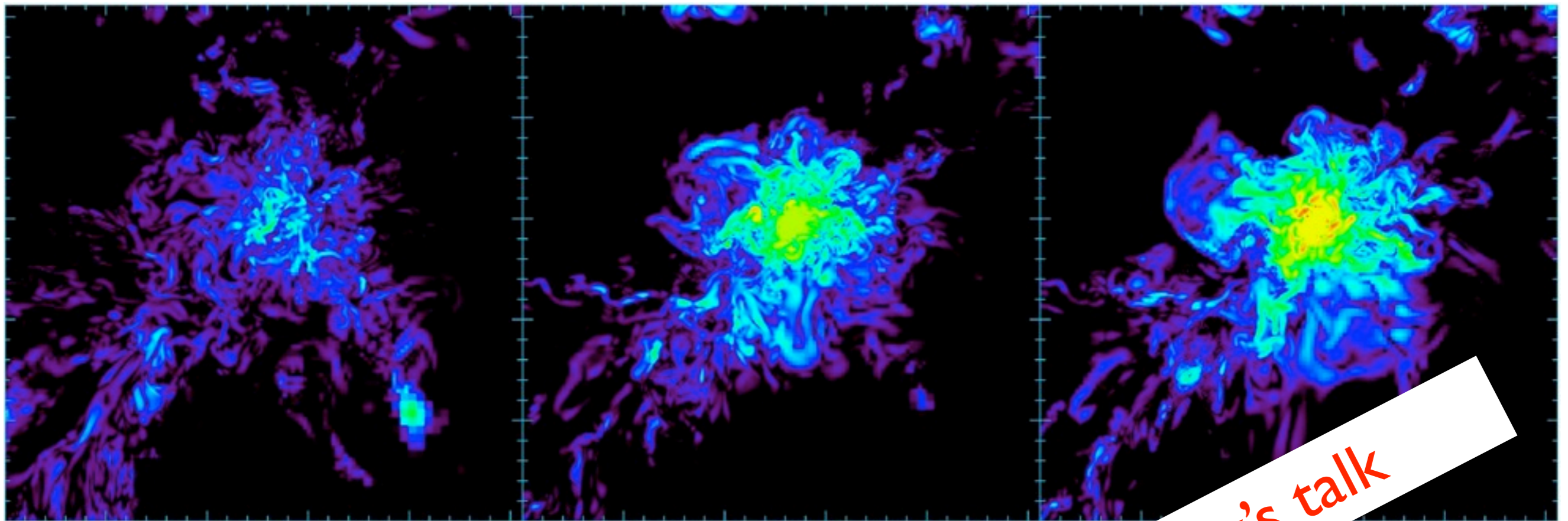


Ruszkowski et al (2007)





# Can turbulence overwhelm the MTI?

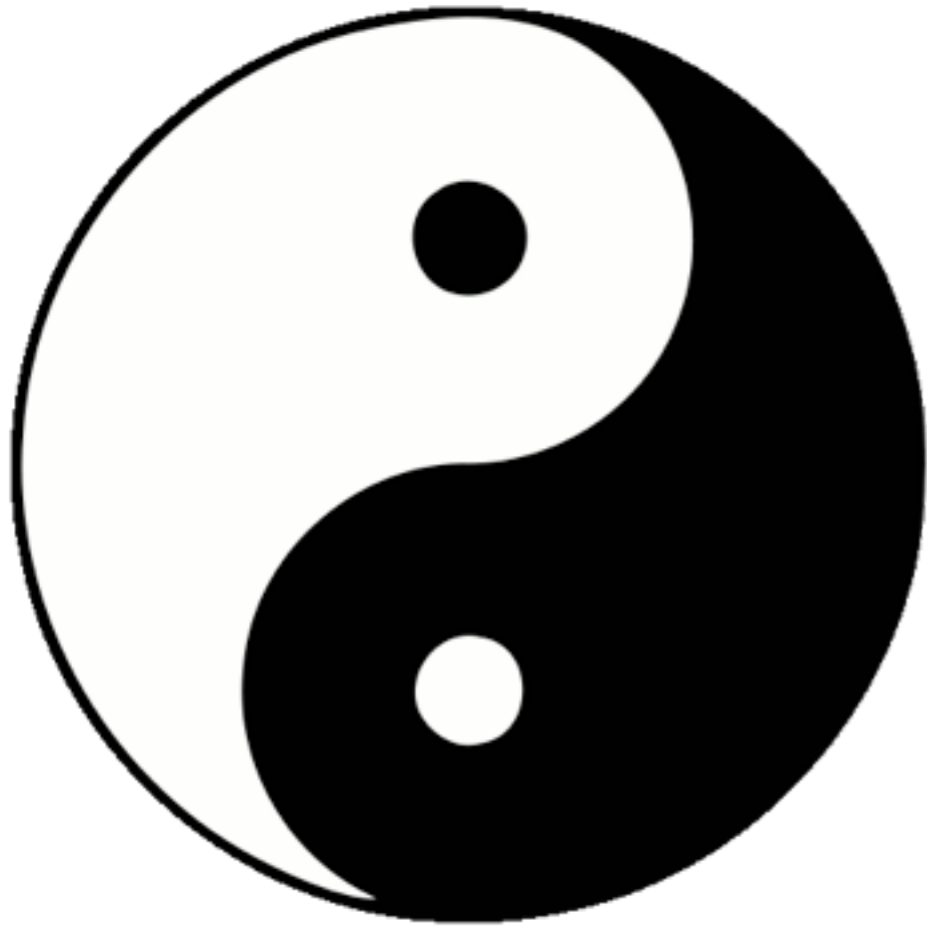


See Mateusz's talk

# Conclusions

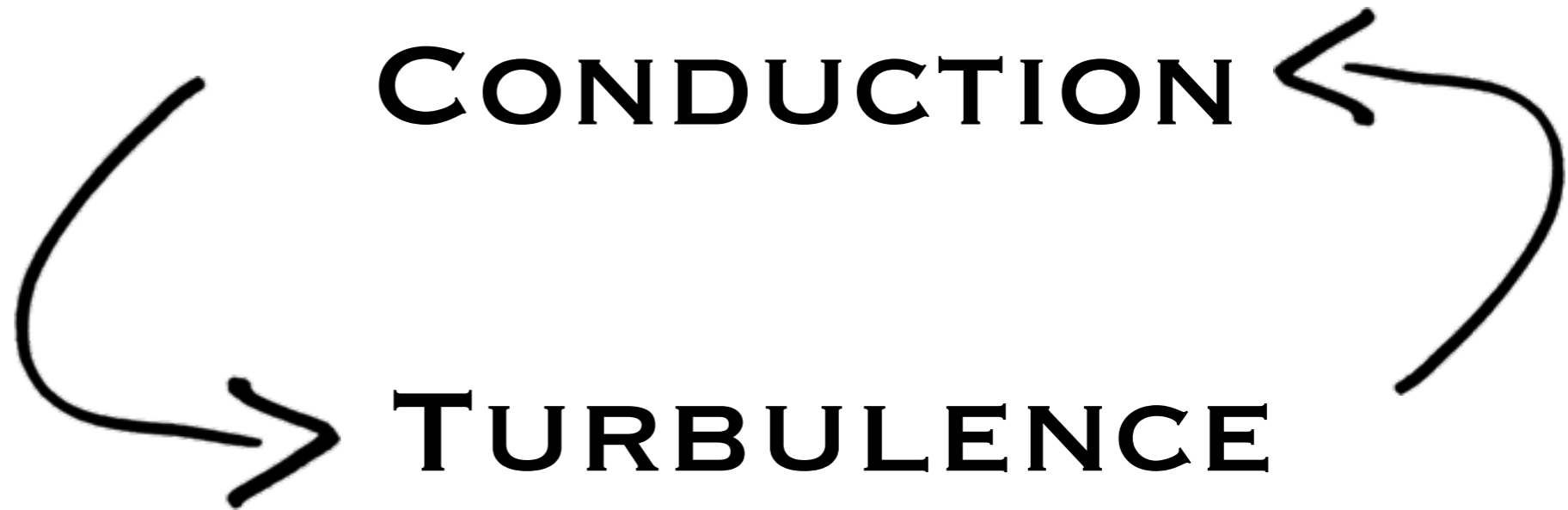
- Conduction reduces buoyancy forces
- Mild, subsonic turbulence is sufficient to overcome HBI and isotropize magnetic field, restoring conduction
- Galaxies which excite g-modes can provide this volume-filling turbulence
- Turbulent heat diffusion is important, but may be aided by conduction

# THE TAO OF HEAT TRANSPORT



CONDUCTION

TURBULENCE



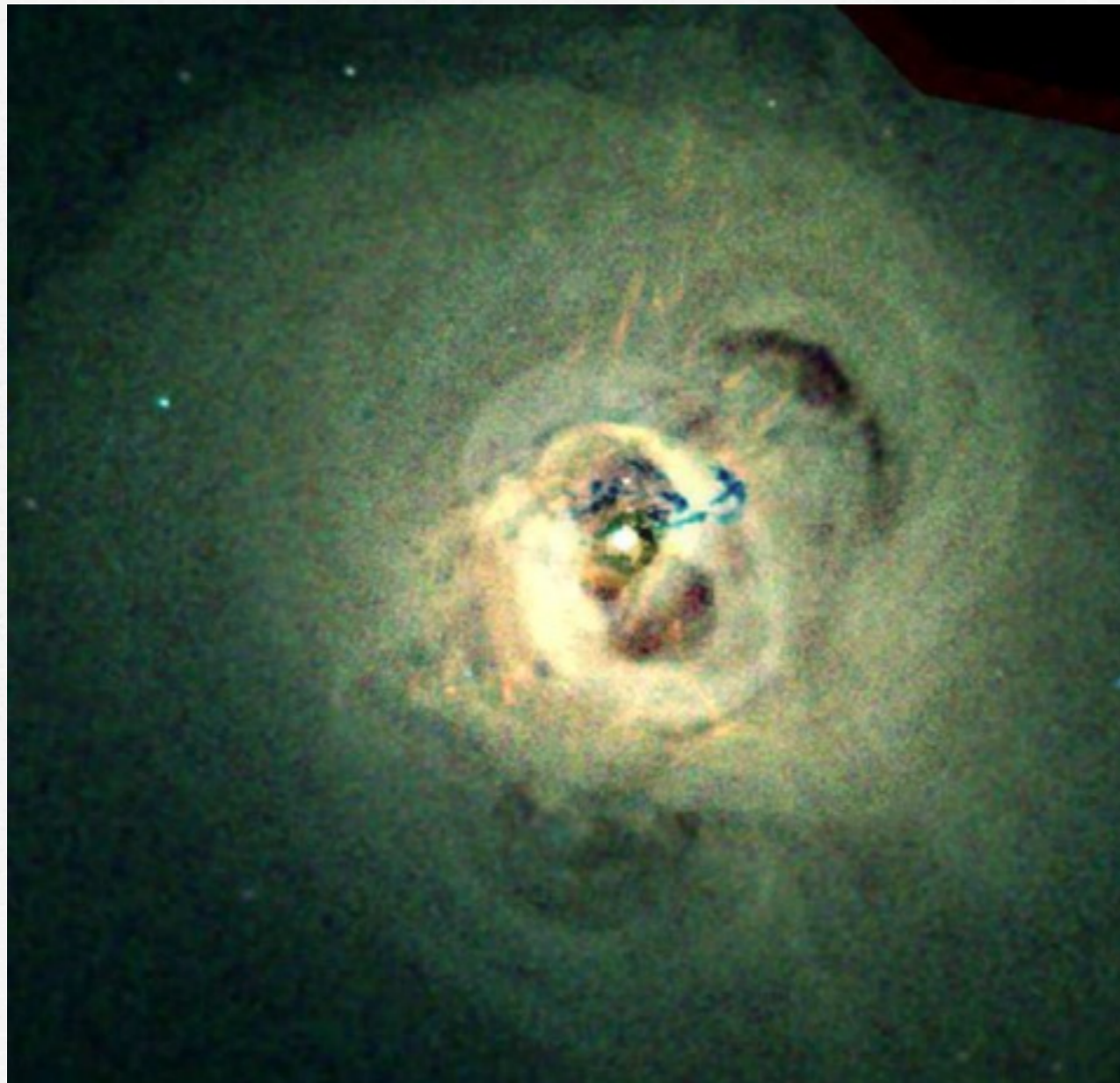




**Other stuff**

# AGN/radio galaxy heating

See review by McNamara & Nulsen (2007)



Bubbles observed in ICM,  
filled with hot/relativistic  
plasma

Maybe: entrain cold gas  
pdv work

This talk: **cosmic ray  
heating**

(Guo & Oh 2008)

Chandra image, Perseus cluster



# A key problem: CR transport is slow

$$\mathbf{F}_c = \gamma_c E_c (\mathbf{u} + \mathbf{v}_A) - n \kappa_c (n \cdot \nabla E_c), \quad (\text{A14})$$

$$\frac{\partial E_c}{\partial t} = (\gamma_c - 1) (\mathbf{u} + \mathbf{v}_A) \cdot \nabla E_c - \nabla \cdot \mathbf{F}_c + \bar{Q}. \quad (\text{A15})$$

Diffusive and other CR transport timescales are  
long

Leads to overpressured center with insufficient heating at outskirts (though may drive turbulent convection: Chandran & collaborators)



# Method

- 1D Zeus code: solve time-dependent hydrodynamic equations + CR heating & transport equations
- calculate steady steady CR spectrum, assuming Coulomb, hadronic and Alfvén-wave energy losses (latter dominates):

$$\Gamma_{wave} = v_A \frac{dP_c}{dr}$$



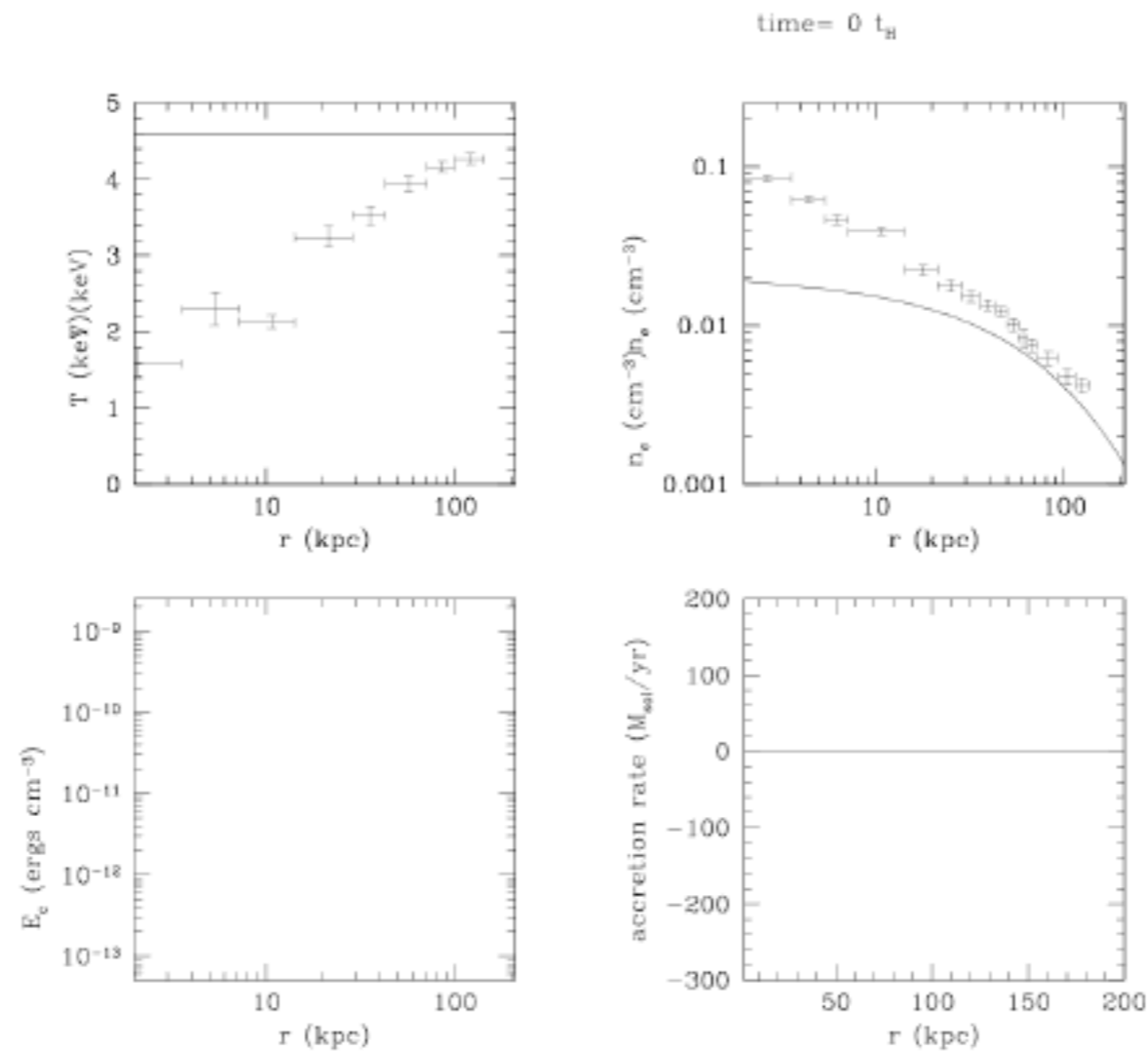
- Assume energy density in bubbles is a power-law with radius (note: CR injection rate depends on gas cooling---feedback effect)

$$L_{\text{bubble}} \sim -\epsilon \dot{M}_{\text{in}} c^2 \left( \frac{r}{r_0} \right)^{-\nu} \quad \text{for } r > r_0,$$

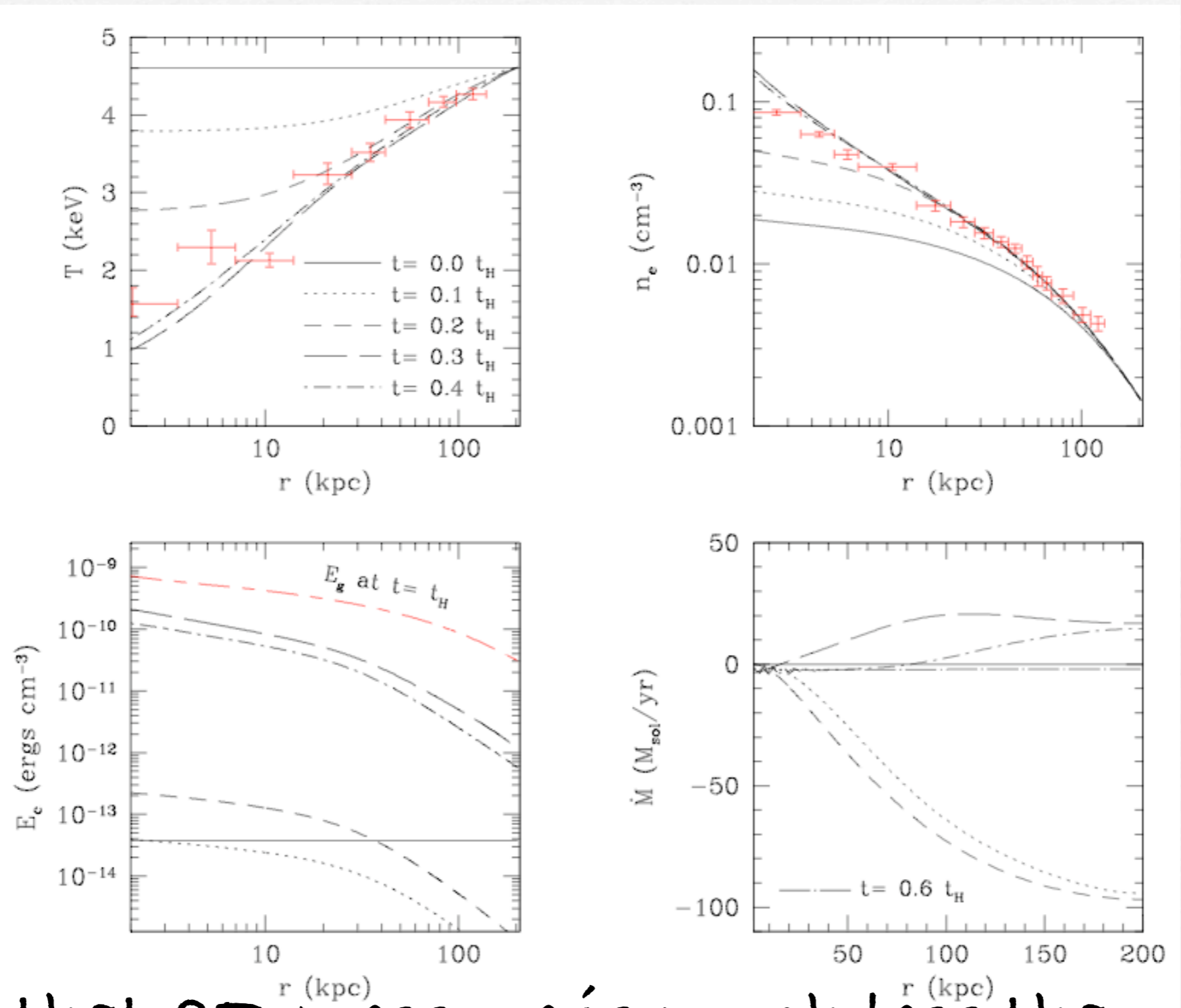
$$\begin{aligned} Q_c = \nabla \cdot \mathbf{F}_{\text{bubble}} &\sim -\frac{1}{4\pi r^2} \frac{\partial L_{\text{bubble}}}{\partial r} \left[ 1 - e^{-(r/r_0)^2} \right] \\ &\sim -\frac{\nu \epsilon \dot{M}_{\text{in}} c^2}{4\pi r_0^3} \left( \frac{r}{r_0} \right)^{-3-\nu} \left[ 1 - e^{-(r/r_0)^2} \right], \end{aligned} \quad (19)$$

Slope is free parameter, implicitly specifies bubble disruption rate

# Bottom line: it works!



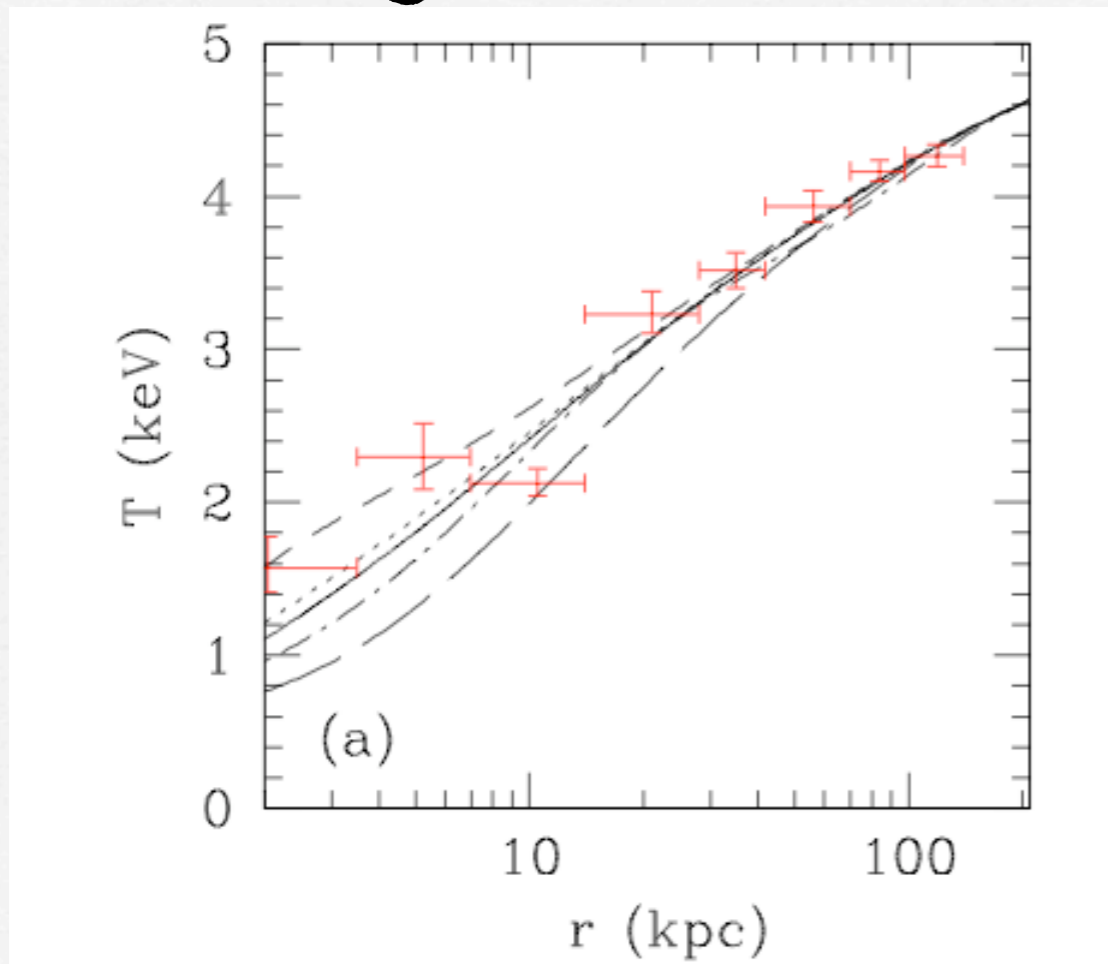




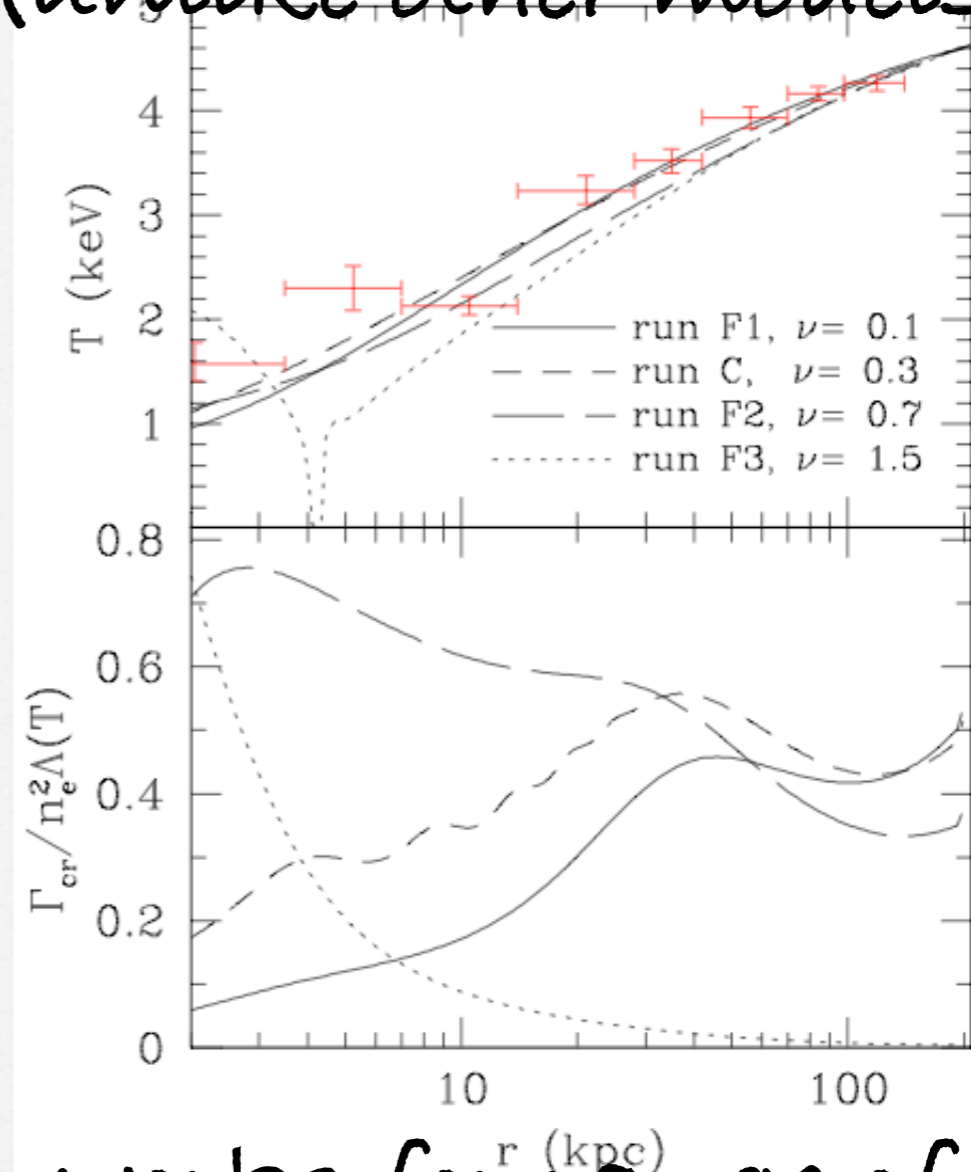
Note that CR pressure is much less than thermal pressure

# No fine tuning

works (i.e., no massive cooling flow) starting from arbitrary initial conditions (unlike other models...)



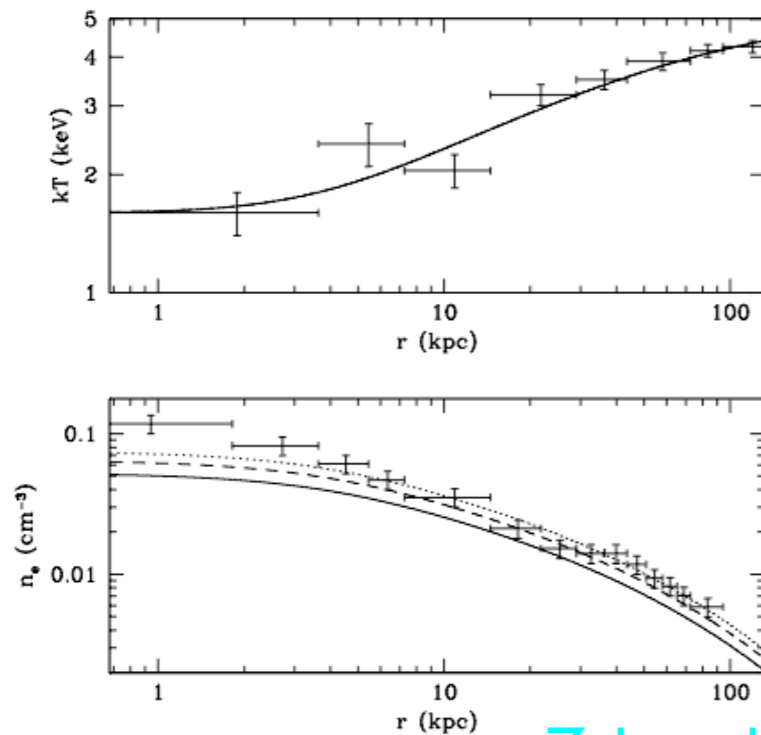
works for range of AGN + conduction parameters



works for range of CR profiles

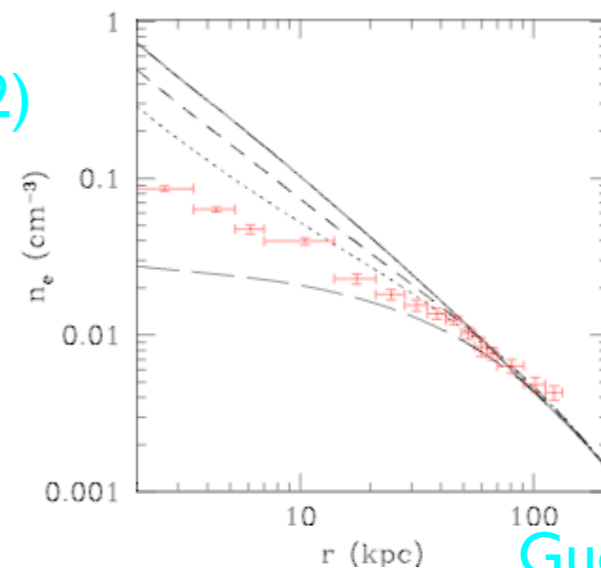
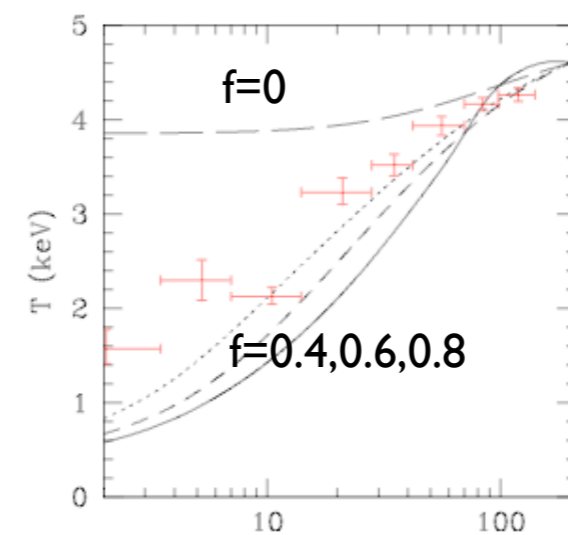


# Theoretical Puzzle: How to Guarantee Stability and Avoid Fine Tuning ?



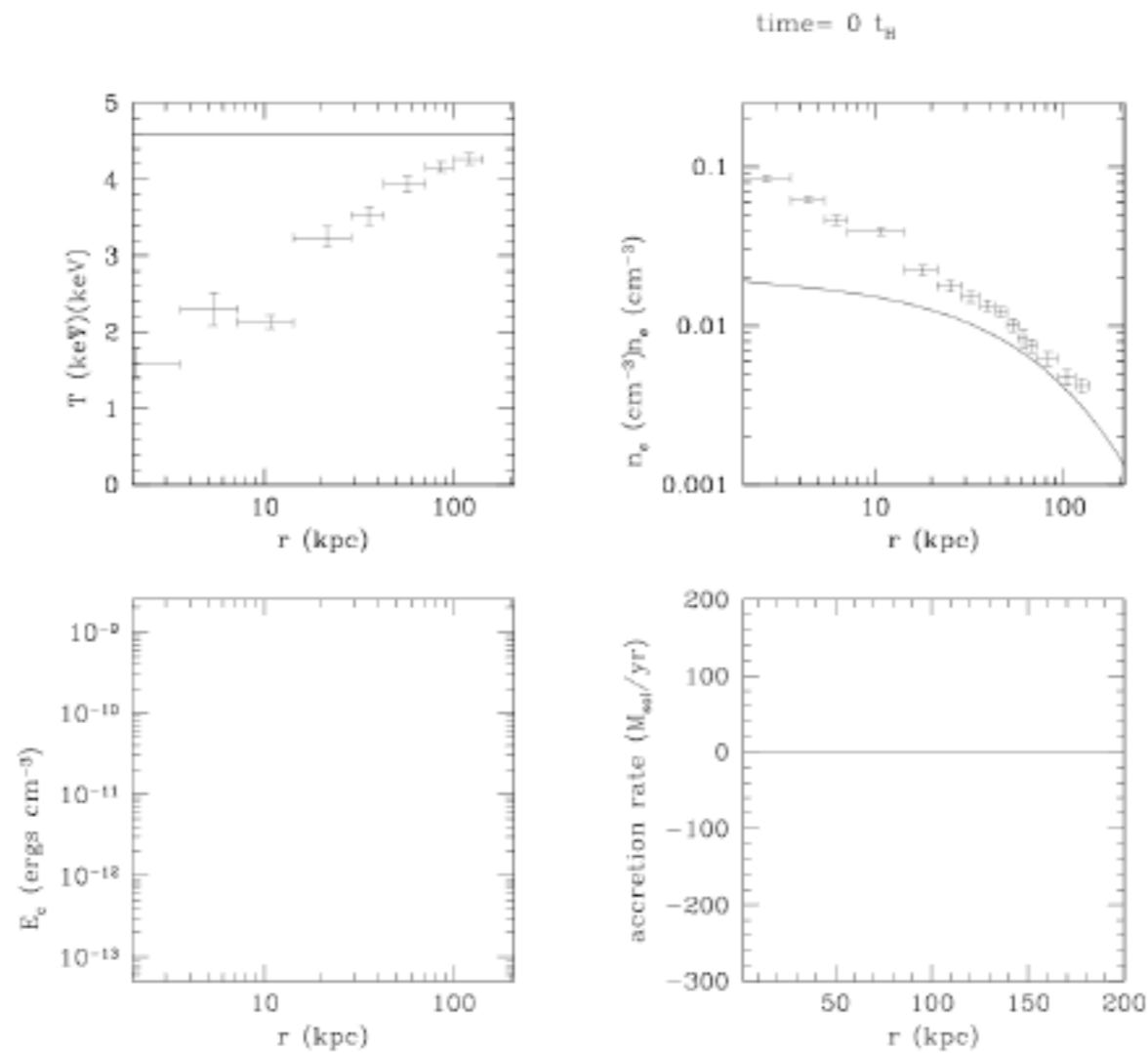
Zakamska & Narayan (2002)

Conduction only model can fit observations (solve eigenvalue problem)



But have to tune parameters, and it won't evolve toward this state in general...

Guo & Oh 2008



Guo & Oh (2008) -- see also Ruszkowski & Begelman (2002)

## 2) Observed profiles are an attractor solution

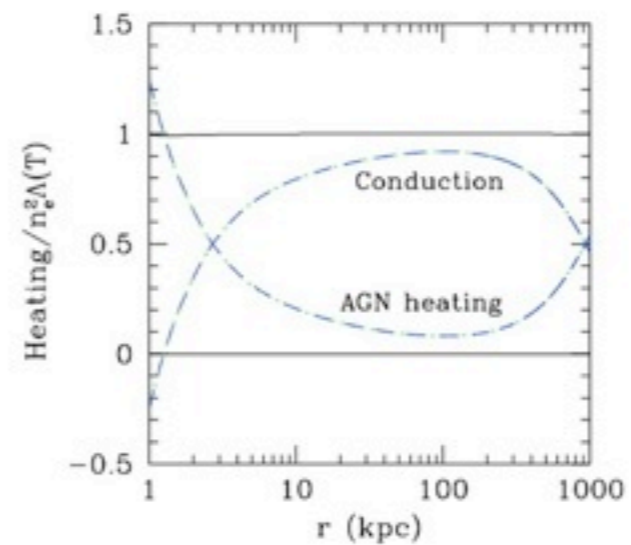
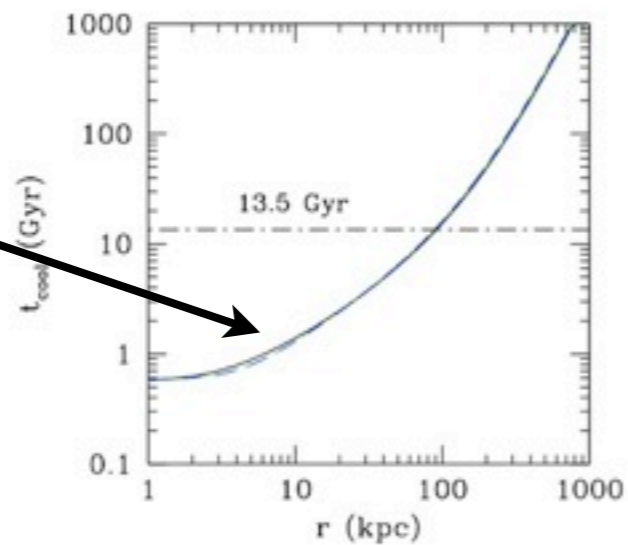
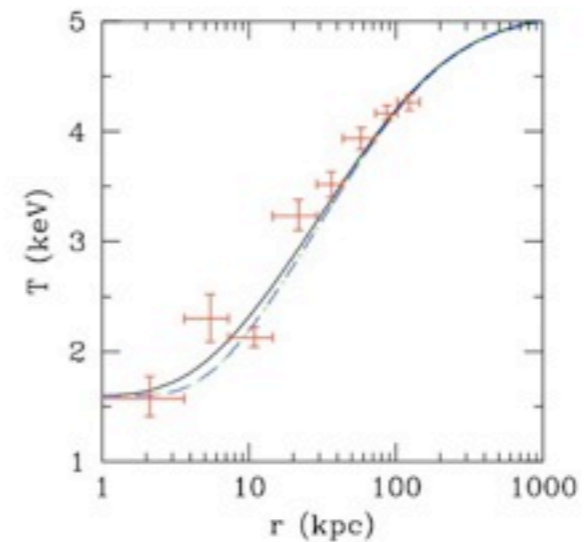
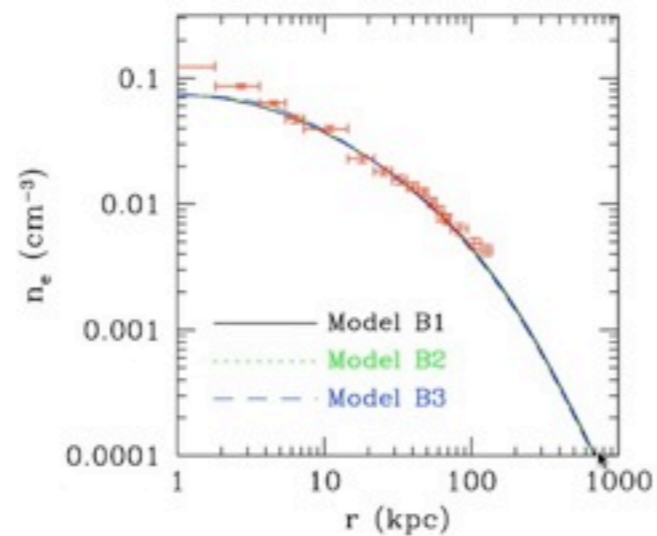
# How to understand this?



‘Pick any two’

Semi-analytic model: Explore parameter space **quickly**  
Aids in physical intuition

note short  
cooling  
time



First, build a background equilibrium solution



..and perform a **global** stability

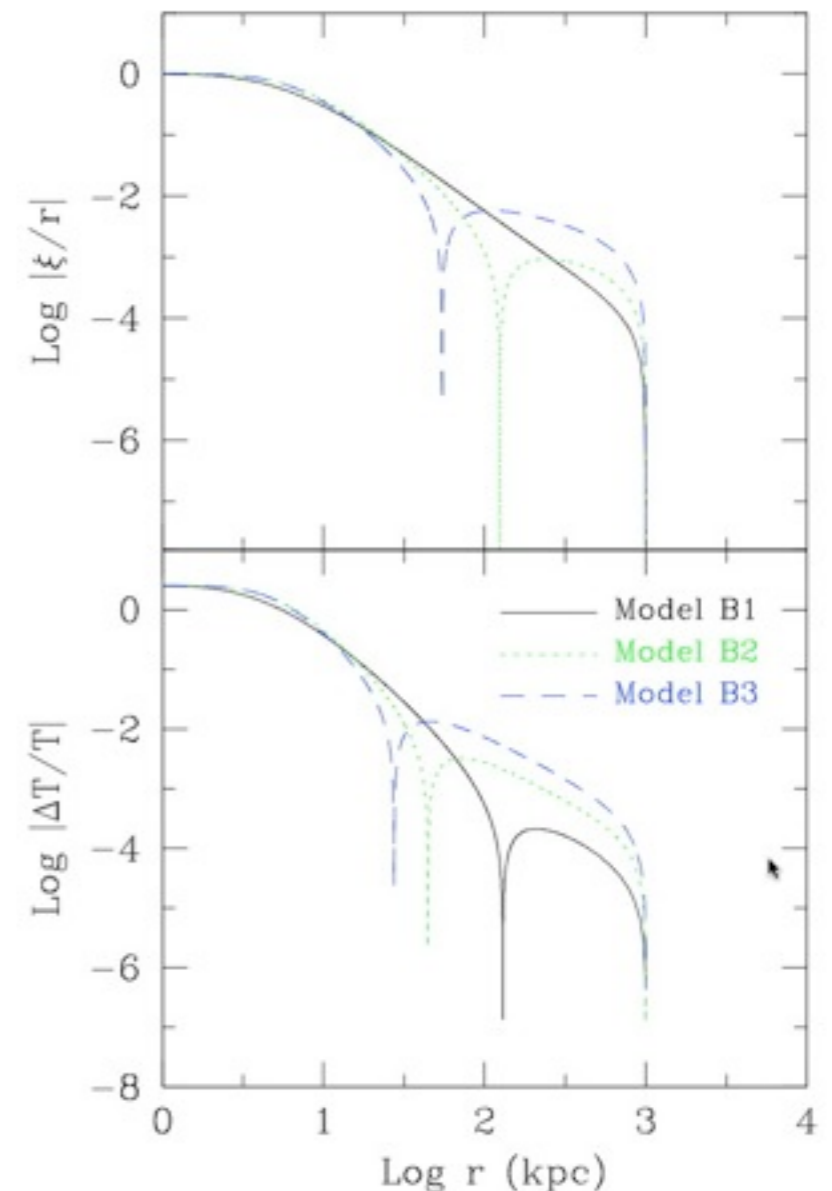
~~WKB~~ analysis

Done for conduction only case by  
Kim & Narayan (2003)

$$\left(\frac{P}{\rho} - v^2\right) \frac{d}{dr}(\nabla \cdot \xi) = \left(r\sigma^2 + r \frac{d^2\Phi}{dr^2}\right) \frac{\xi}{r} + \frac{1}{\rho} \frac{d}{dr} \left(P \frac{\Delta T}{T}\right) - 2v^2 \frac{d}{dr} \left(\frac{\xi}{r}\right) + \left(2\sigma v + v \frac{dv}{dr} - \frac{1}{\rho} \frac{dP}{dr}\right) \frac{d\xi}{dr} \quad (33)$$

$$\kappa T \frac{d}{dr} \left(\frac{\Delta T}{T}\right) = F \left[ \frac{7}{2} \frac{\Delta T}{T} - r \frac{d}{dr} \left(\frac{\xi}{r}\right) + \frac{\xi}{r} \right] + \frac{\Delta L_r}{4\pi r^2} \quad (34)$$

$$\frac{1}{4\pi r^2} \frac{d}{dr} \Delta L_r = (P\sigma - \rho^2 \mathcal{L}_\rho - \mathcal{H})(\nabla \cdot \xi) - \Delta \mathcal{H} + \left(\frac{P\sigma}{\gamma-1} + \rho T \mathcal{L}_T + \frac{v}{\gamma-1} \frac{dP}{dr} - \frac{\gamma v}{\gamma-1} \frac{P}{\rho} \frac{d\rho}{dr}\right) \frac{\Delta T}{T} + Pv \frac{d}{dr}(\nabla \cdot \xi) + \frac{Pv}{\gamma-1} \frac{d}{dr} \left(\frac{\Delta T}{T}\right) \quad (35)$$



Growth rate is an eigenvalue of analysis

Explore parameter space rapidly!

Guo, Oh & Ruszkowski 2008

# Just as in Stellar Structure calculations...

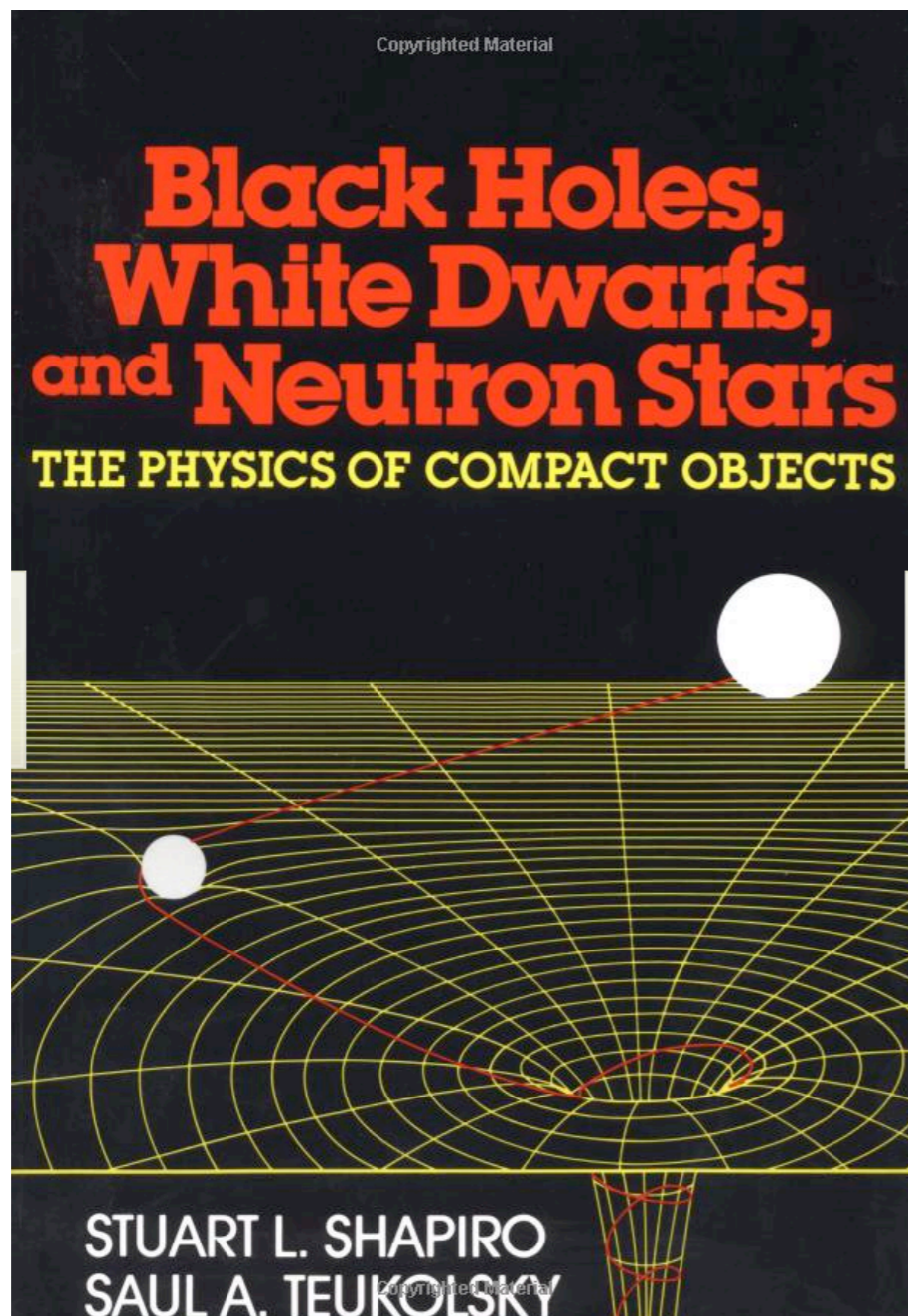


Fig from A. Piro

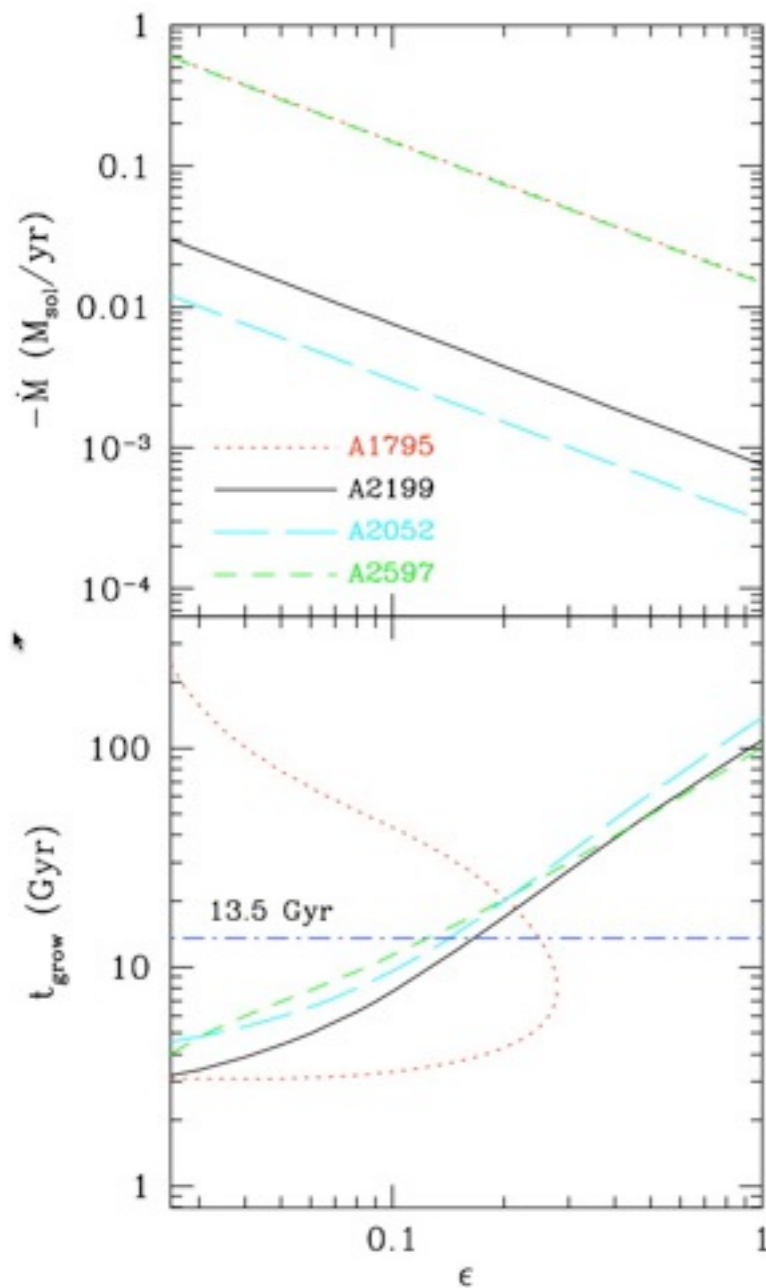
# Global unstable modes suppressed with AGN!

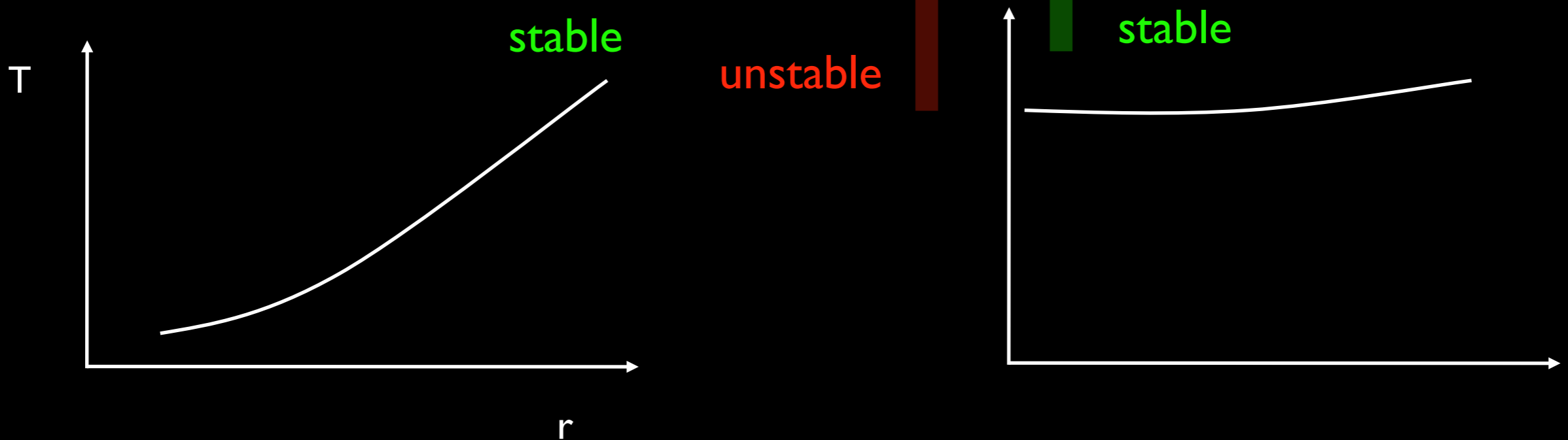
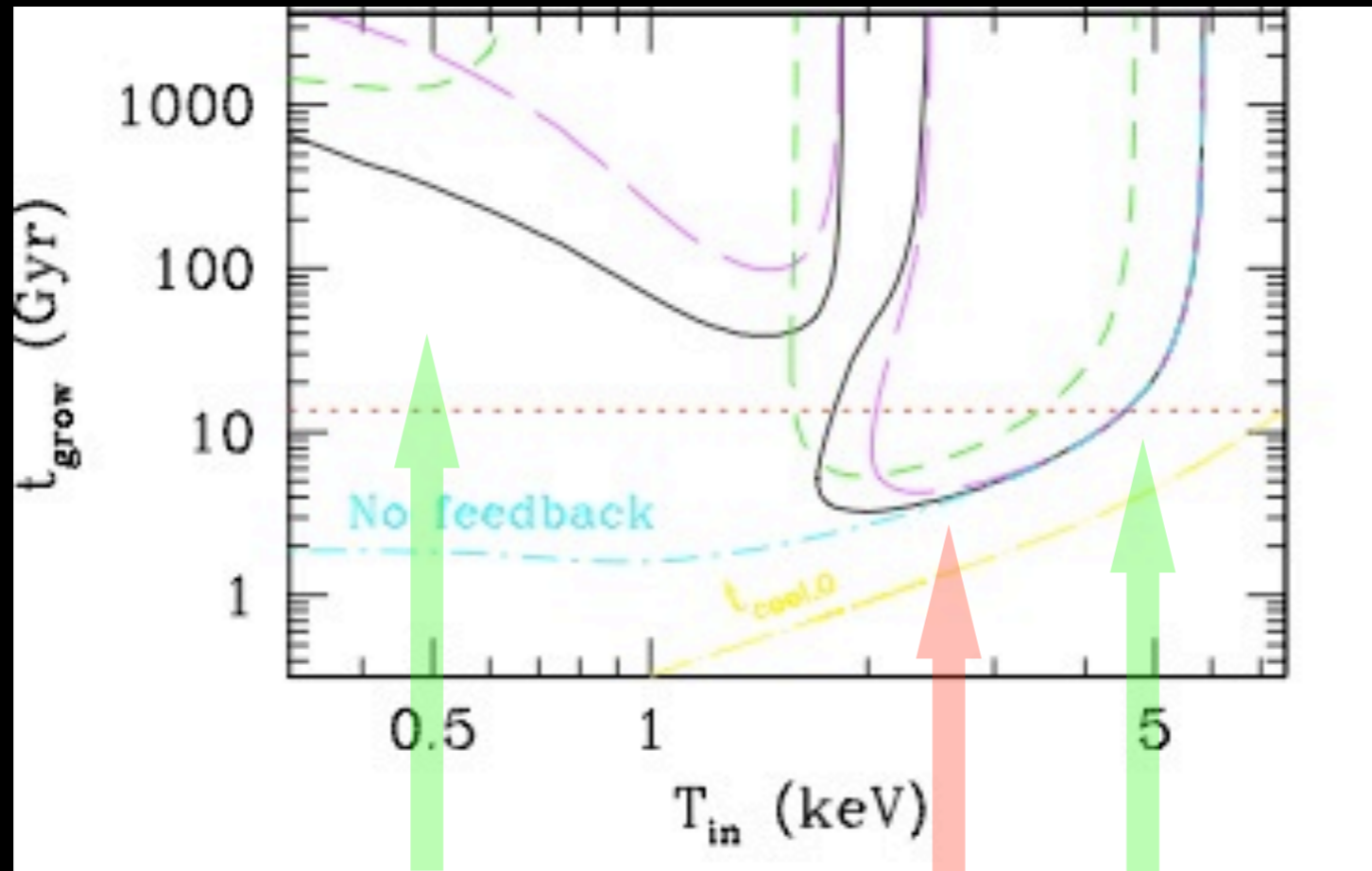
Suppression depends on efficiency

$$L_{\text{agn}} = -\epsilon \dot{M}_{\text{in}} c^2,$$

The crucial term: feedback

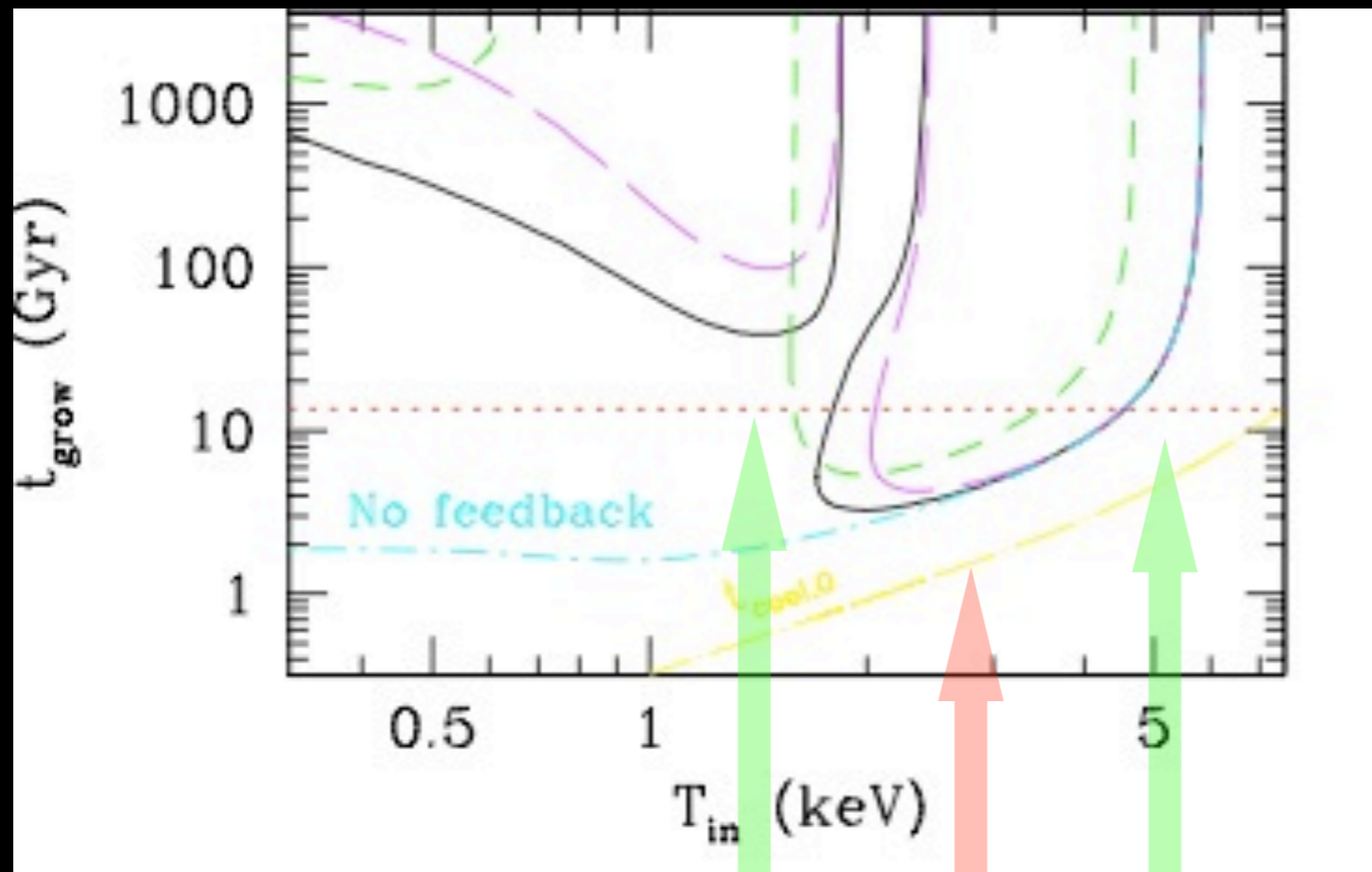
$$\Delta \mathcal{H}_{\text{feed}} \equiv \mathcal{H} \Delta \dot{M}(r_{\text{in}}) / \dot{M}_{\text{in}} = \frac{\mathcal{H} \sigma}{v_0} \xi(r_{\text{in}}),$$





Stability is bimodal!





stable

unstable

stable

## Bimodality !

cool core



stabilized by AGN + conduction

non cool core



stabilized by conduction