

# Applied Symbology

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past and in progress work with Goncharov, Spradlin, Vergu, also Prygarin



# Plan

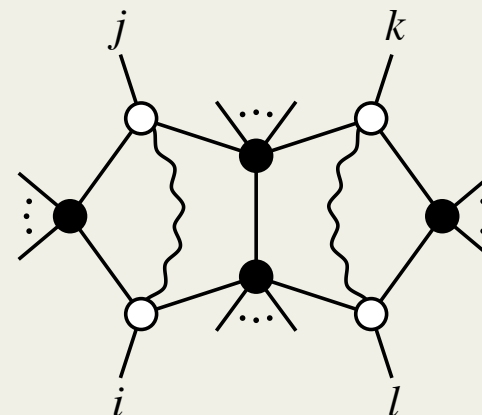
- Introduction
- Review: Symbols and Amplitudes  
Goncharov, Spradlin, Vergu, AV
- One-loop integrals from mixed Tate motives  
Spradlin, AV
- Extracting Motivic Content from the Symbol  
with Goncharov, Spradlin, Vergu
- Regge limit  
with Prygarin, Spradlin, Vergu
- Outlook and Conclusions

We have learned a lot about N=4 Yang-Mills scattering amplitudes.

Many have contributed to these recent exciting developments including:

Alday, Arkani-Hamed, Bargheer, Beisert, Belitski, Bern, Berkovits, Boels, Bourjaily, Brandhuber, Broedel, Bjerrum-Bohr, Bullimore, Britto, Cachazo, Caron-Huot, Carrasco, Cheung, Damgaard, Del Duca, Duhr, Dixon, Dolan, Duhr, Drummons, Eden, Elvang, Fend, Ferro, Forde, Freedman, Gaiotto, Goddard, Goncharov, Green, Henn, Heslop, Hodges, Huang, Ita, Johannson, Kallosh, Kaplan, Khoze, Kiermaier, Korchemsky, Kosower, Lipatov, Loebbert, Maitre, Mafra, Maldacena, Mason, Naculich, Nastase, Plefka, Prygarin, Roiban, Sabio Vera, Schnizer, Sever, Skinner, Smirnov, Sokatchev, Spence, Spradlin, Staudacher, Stelle, Steiberger, Svrcek, Taylor, Travaglini, Trnka, Tye, Vanhove, Vergu, Vieira, Wen, Witten

# Amplitudes and Integrand



- Standard Lore

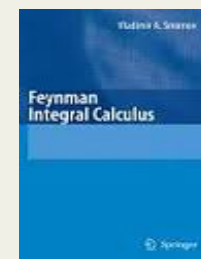
$$\textit{Amplitude} = \int d(\textit{loop momenta}) (\textit{Integrand})$$

- Integrand = solved for all loops and all legs in N=4

Bern, Dixon, Kosower; Arkani-Hamed, Cachazo, et al

- Integral = **still very hard to evaluate!** This step requires a lot of blood and tears begging for some new technology!

e.g. Smirnov's book



- Symbol = “half way between integrand and amplitude”

Goncharov, Spradlin, Vergu, AV

# Symbol of an Amplitude

- Symbol of an amplitude is **a new very powerful tool** for analysis of multiloop amplitudes which comes from modern mathematics called **theory of motives**.

Goncharov, Spradlin, Vergu, AV

- Symbol of an amplitude
  - leads to the compact expression for the amplitude
  - captures the essential physical (and “motivic”) content of the amplitude

# Motivations

Simplest nontrivial multi-loop amplitude in N=4 Yang-Mills is the **2-loop 6-particle MHV remainder function** = what is left after subtracting IR divergences, leaving a finite and dual conformal invariant quantity

$$R(u_1, u_2, u_2)$$

$$u_1 = \frac{s_{12}s_{45}}{s_{123}s_{345}} \quad u_2 = \frac{s_{23}s_{56}}{s_{234}s_{456}} \quad u_3 = \frac{s_{34}s_{61}}{s_{345}s_{234}} \quad s_{123} = (k_1 + k_2 + k_3)^2$$

Numerically known since '08 Bern, Dixon, Kosower, Roiban, Vergu, Spradlin, AV Henn, Drummond, Korchemsky, Sokatchev

In a heroic effort, Del Duca, Duhr, Smirnov '09 found a manageable way to evaluate the relevant diagrams and obtained **an analytic formula** for it:















$$\begin{aligned}
& \frac{3}{4}\mathcal{G}\left(v_{231}, 1, \frac{1}{1-u_2}; 1\right)H(0; u_3) + \frac{3}{4}\mathcal{G}\left(v_{231}, \frac{1}{1-u_2}, 1; 1\right)H(0; u_3) + \\
& \frac{1}{4}\mathcal{G}\left(v_{312}, 1, \frac{1}{1-u_3}; 1\right)H(0; u_3) + \frac{1}{4}\mathcal{G}\left(v_{312}, \frac{1}{1-u_3}, 1; 1\right)H(0; u_3) + \\
& \frac{1}{4}\mathcal{G}\left(v_{321}, 1, \frac{1}{1-u_3}; 1\right)H(0; u_3) + \frac{1}{4}\mathcal{G}\left(v_{321}, \frac{1}{1-u_3}, 1; 1\right)H(0; u_3) + \\
& \frac{1}{4}G\left(\frac{1}{u_1}, \frac{1}{u_1+u_3}; 1\right)H(0; u_1)H(0; u_3) + \\
& \frac{1}{4}G\left(\frac{1}{1-u_2}, \frac{u_3-1}{u_2+u_3-1}; 1\right)H(0; u_1)H(0; u_3) + \\
& \frac{1}{4}G\left(\frac{1}{u_3}, \frac{1}{u_1+u_3}; 1\right)H(0; u_1)H(0; u_3) - \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_2}, u_{231}; 1\right)H(0; u_1)H(0; u_3) - \\
& \frac{1}{4}G\left(\frac{1}{1-u_2}, v_{213}; 1\right)H(0; u_1)H(0; u_3) - \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_2}, v_{231}; 1\right)H(0; u_1)H(0; u_3) + \\
& \frac{5}{24}\pi^2H(0; u_1)H(0; u_3) + \frac{1}{4}G\left(\frac{1}{1-u_1}, \frac{u_2-1}{u_1+u_2-1}; 1\right)H(0; u_2)H(0; u_3) + \\
& \frac{1}{4}G\left(\frac{1}{u_2}, \frac{1}{u_2+u_3}; 1\right)H(0; u_2)H(0; u_3) + \frac{1}{4}G\left(\frac{1}{u_3}, \frac{1}{u_2+u_3}; 1\right)H(0; u_2)H(0; u_3) - \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_1}, u_{123}; 1\right)H(0; u_2)H(0; u_3) - \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_1}, v_{123}; 1\right)H(0; u_2)H(0; u_3) - \\
& \frac{1}{4}G\left(\frac{1}{1-u_1}, v_{132}; 1\right)H(0; u_2)H(0; u_3) + \frac{5}{24}\pi^2H(0; u_2)H(0; u_3) + \\
& 3H(0; u_2)H(0; u_1)H(0; u_3) + 3H(0; u_1)H(0; u_2)H(0; u_3) + \\
& \frac{1}{4}H(0; u_2)H\left(0, 1; \frac{u_1+u_2-1}{u_2-1}\right)H(0; u_3) + \frac{1}{2}H(0; u_1)H(0, 1; (u_1+u_3))H(0; u_3) + \\
& \frac{1}{4}H(0; u_1)H\left(0, 1; \frac{u_2+u_3-1}{u_3-1}\right)H(0; u_3) + \frac{1}{2}H(0; u_2)H(0, 1; (u_2+u_3))H(0; u_3) + \\
& \frac{3}{4}H(0; u_2)H(1, 0; u_1)H(0; u_3) + \frac{3}{4}H(0; u_1)H(1, 0; u_2)H(0; u_3) + \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_2}, v_{213}; 1\right)H(0, 0; u_1) + \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_2}, v_{231}; 1\right)H(0, 0; u_1) + \\
& \frac{1}{4}G\left(\frac{1}{1-u_3}, v_{312}; 1\right)H(0, 0; u_1) + \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_3}, v_{321}; 1\right)H(0, 0; u_1) - \frac{23}{24}\pi^2H(0, 0; u_1) + \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_1}, v_{123}; 1\right)H(0, 0; u_2) + \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_1}, v_{132}; 1\right)H(0, 0; u_2) + \\
& \frac{1}{4}G\left(\frac{1}{1-u_3}, v_{312}; 1\right)H(0, 0; u_2) + \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_3}, v_{321}; 1\right)H(0, 0; u_2) - \\
& \frac{25}{4}H(0, 0; u_1)H(0, 0; u_2) - \frac{23}{24}\pi^2H(0, 0; u_2) + \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_1}, v_{123}; 1\right)H(0, 0; u_3) + \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_1}, v_{132}; 1\right)H(0, 0; u_3) + \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_2}, v_{213}; 1\right)H(0, 0; u_3) + \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_2}, v_{231}; 1\right)H(0, 0; u_3) + 3H(0; u_1)H(0; u_2)H(0, 0; u_3) - \\
& \frac{25}{4}H(0, 0; u_1)H(0, 0; u_3) - \frac{25}{4}H(0; u_2)H(0, 0; u_3) - \frac{23}{24}\pi^2H(0, 0; u_3) + \frac{1}{12}\pi^2H(0, 1; u_1) + \\
& \frac{1}{12}\pi^2H(0, 1; u_2) - \frac{1}{24}\pi^2H\left(0, 1; \frac{u_1+u_2-1}{u_2-1}\right) + \frac{1}{2}H(0; u_1)H(0; u_2)H(0, 1; (u_1+u_2)) +
\end{aligned}$$

$$\begin{aligned}
& \frac{1}{12}\pi^2H(0, 1; (u_1+u_2)) + \frac{1}{12}\pi^2H(0, 1; u_3) + \frac{1}{4}H(0; u_1)H(0; u_2)H\left(0, 1; \frac{u_1+u_3-1}{u_1-1}\right) - \\
& \frac{1}{24}\pi^2H\left(0, 1; \frac{u_1+u_3-1}{u_1-1}\right) + \frac{1}{12}\pi^2H(0, 1; (u_1+u_3)) - \frac{1}{24}\pi^2H\left(0, 1; \frac{u_2+u_3-1}{u_3-1}\right) + \\
& \frac{1}{12}\pi^2H(0, 1; (u_2+u_3)) - \frac{1}{2}G\left(0, \frac{1}{u_1+u_2}; 1\right)H(1, 0; u_1) - \\
& \frac{1}{2}G\left(0, \frac{1}{u_1+u_3}; 1\right)H(1, 0; u_1) + \frac{1}{4}G\left(\frac{1}{u_1}, \frac{1}{u_1+u_2}; 1\right)H(1, 0; u_1) + \\
& \frac{1}{4}G\left(\frac{1}{u_1}, \frac{1}{u_1+u_3}; 1\right)H(1, 0; u_1) + \frac{1}{4}G\left(\frac{1}{u_2}, \frac{1}{u_1+u_2}; 1\right)H(1, 0; u_1) + \\
& \frac{1}{4}G\left(\frac{1}{1-u_3}, \frac{u_1-1}{u_1+u_3-1}; 1\right)H(1, 0; u_1) + \frac{1}{4}G\left(\frac{1}{u_3}, \frac{1}{u_1+u_3}; 1\right)H(1, 0; u_1) - \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_3}, u_{312}; 1\right)H(1, 0; u_1) - \frac{3}{4}H(0, 0; u_2)H(1, 0; u_1) - \frac{3}{4}H(0, 0; u_3)H(1, 0; u_1) + \\
& \frac{1}{4}H\left(0, 1; \frac{u_1+u_3-1}{u_1-1}\right)H(1, 0; u_1) - \frac{1}{3}\pi^2H(1, 0; u_1) - \frac{1}{2}G\left(0, \frac{1}{u_1+u_2}; 1\right)H(1, 0; u_2) - \\
& \frac{1}{2}G\left(0, \frac{1}{u_2+u_3}; 1\right)H(1, 0; u_2) + \frac{1}{4}G\left(\frac{1}{1-u_1}, \frac{u_2-1}{u_1+u_2-1}; 1\right)H(1, 0; u_2) + \\
& \frac{1}{4}G\left(\frac{1}{u_1}, \frac{1}{u_1+u_2}; 1\right)H(1, 0; u_2) + \frac{1}{4}G\left(\frac{1}{u_2}, \frac{1}{u_1+u_2}; 1\right)H(1, 0; u_2) + \\
& \frac{1}{4}G\left(\frac{1}{u_2}, \frac{1}{u_2+u_3}; 1\right)H(1, 0; u_2) + \frac{1}{4}G\left(\frac{1}{u_3}, \frac{1}{u_2+u_3}; 1\right)H(1, 0; u_2) - \\
& \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_1}, u_{123}; 1\right)H(1, 0; u_2) - \frac{3}{4}H(0, 0; u_1)H(1, 0; u_2) - \frac{3}{4}H(0, 0; u_3)H(1, 0; u_2) + \\
& \frac{1}{4}H\left(0, 1; \frac{u_1+u_2-1}{u_2-1}\right)H(1, 0; u_2) - \frac{1}{4}H(1, 0; u_1)H(1, 0; u_2) - \frac{1}{3}\pi^2H(1, 0; u_2) - \\
& \frac{1}{2}G\left(0, \frac{1}{u_1+u_3}; 1\right)H(1, 0; u_3) - \frac{1}{2}G\left(0, \frac{1}{u_2+u_3}; 1\right)H(1, 0; u_3) + \\
& \frac{1}{4}G\left(\frac{1}{u_1}, \frac{1}{u_1+u_3}; 1\right)H(1, 0; u_3) + \frac{1}{4}G\left(\frac{1}{1-u_2}, \frac{u_3-1}{u_2+u_3-1}; 1\right)H(1, 0; u_3) + \\
& \frac{1}{4}G\left(\frac{1}{u_2}, \frac{1}{u_2+u_3}; 1\right)H(1, 0; u_3) - \frac{1}{3}\pi^2H(1, 0; u_3) + \frac{1}{4}G\left(\frac{1}{u_3}, \frac{1}{u_1+u_3}; 1\right)H(1, 0; u_3) + \\
& \frac{1}{4}G\left(\frac{1}{u_3}, \frac{1}{u_2+u_3}; 1\right)H(1, 0; u_3) - \frac{1}{4}\mathcal{G}\left(\frac{1}{1-u_2}, u_{231}; 1\right)H(1, 0; u_3) + \\
& \frac{3}{4}H(0; u_1)H(0; u_2)H(1, 0; u_3) - \frac{3}{4}H(0, 0; u_1)H(1, 0; u_3) - \frac{3}{4}H(0, 0; u_2)H(1, 0; u_3) + \\
& \frac{1}{4}H\left(0, 1; \frac{u_2+u_3-1}{u_3-1}\right)H(1, 0; u_3) - \frac{1}{4}H(1, 0; u_1)H(1, 0; u_3) - \frac{1}{4}H(1, 0; u_2)H(1, 0; u_3) + \\
& \frac{1}{24}\pi^2H(1, 1; u_1) + \frac{1}{24}\pi^2H(1, 1; u_2) + \frac{1}{24}\pi^2H(1, 1; u_3) + \frac{1}{2}H(0; u_2)H(0, 0, 0; u_1) + \\
& \frac{1}{2}H(0; u_3)H(0, 0, 0; u_2) + \frac{1}{2}H(0; u_1)H(0, 0, 0; u_3) - \frac{1}{2}H(0; u_2)H\left(0, 0, 1; \frac{u_1+u_2-1}{u_2-1}\right) - \\
& \frac{1}{2}H(0; u_3)H\left(0, 0, 1; \frac{u_1+u_2-1}{u_2-1}\right) - H(0; u_1)H(0, 0, 1; (u_1+u_2)) - \\
& H(0; u_2)H(0, 0, 1; (u_1+u_2)) - \frac{1}{2}H(0; u_1)H\left(0, 0, 1; \frac{u_1+u_3-1}{u_1-1}\right) -
\end{aligned}$$



$$\begin{aligned}
& \frac{1}{4}H(0; u_2) \mathcal{H}\left(0, 1, 1; \frac{1}{u_{312}}\right) + \frac{1}{4}H(0; u_2) \mathcal{H}\left(0, 1, 1; \frac{1}{v_{123}}\right) - \frac{1}{4}H(0; u_3) \mathcal{H}\left(0, 1, 1; \frac{1}{v_{123}}\right) - \\
& \frac{1}{4}H(0; u_2) \mathcal{H}\left(0, 1, 1; \frac{1}{v_{132}}\right) + \frac{1}{4}H(0; u_3) \mathcal{H}\left(0, 1, 1; \frac{1}{v_{132}}\right) + \frac{1}{4}H(0; u_1) \mathcal{H}\left(0, 1, 1; \frac{1}{v_{213}}\right) - \\
& \frac{1}{4}H(0; u_3) \mathcal{H}\left(0, 1, 1; \frac{1}{v_{213}}\right) - \frac{1}{4}H(0; u_1) \mathcal{H}\left(0, 1, 1; \frac{1}{v_{231}}\right) + \frac{1}{4}H(0; u_3) \mathcal{H}\left(0, 1, 1; \frac{1}{v_{231}}\right) + \\
& \frac{1}{4}H(0; u_1) \mathcal{H}\left(0, 1, 1; \frac{1}{v_{312}}\right) - \frac{1}{4}H(0; u_2) \mathcal{H}\left(0, 1, 1; \frac{1}{v_{312}}\right) - \frac{1}{4}H(0; u_1) \mathcal{H}\left(0, 1, 1; \frac{1}{v_{321}}\right) + \\
& \frac{1}{4}H(0; u_2) \mathcal{H}\left(0, 1, 1; \frac{1}{v_{321}}\right) + \frac{1}{4}H(0; u_3) \mathcal{H}\left(1, 0, 1; \frac{1}{u_{123}}\right) + \frac{1}{4}H(0; u_1) \mathcal{H}\left(1, 0, 1; \frac{1}{u_{231}}\right) + \\
& \frac{1}{4}H(0; u_2) \mathcal{H}\left(1, 0, 1; \frac{1}{u_{312}}\right) + \frac{1}{4}H(0; u_2) \mathcal{H}\left(1, 0, 1; \frac{1}{v_{123}}\right) - \frac{1}{4}H(0; u_3) \mathcal{H}\left(1, 0, 1; \frac{1}{v_{123}}\right) - \\
& \frac{1}{4}H(0; u_2) \mathcal{H}\left(1, 0, 1; \frac{1}{v_{132}}\right) + \frac{1}{4}H(0; u_3) \mathcal{H}\left(1, 0, 1; \frac{1}{v_{132}}\right) + \frac{1}{4}H(0; u_1) \mathcal{H}\left(1, 0, 1; \frac{1}{v_{213}}\right) - \\
& \frac{1}{4}H(0; u_3) \mathcal{H}\left(1, 0, 1; \frac{1}{v_{213}}\right) - \frac{1}{4}H(0; u_1) \mathcal{H}\left(1, 0, 1; \frac{1}{v_{231}}\right) + \frac{1}{4}H(0; u_3) \mathcal{H}\left(1, 0, 1; \frac{1}{v_{231}}\right) + \\
& \frac{1}{4}H(0; u_1) \mathcal{H}\left(1, 0, 1; \frac{1}{v_{312}}\right) - \frac{1}{4}H(0; u_2) \mathcal{H}\left(1, 0, 1; \frac{1}{v_{312}}\right) - \frac{1}{4}H(0; u_1) \mathcal{H}\left(1, 0, 1; \frac{1}{v_{321}}\right) + \\
& \frac{1}{4}H(0; u_2) \mathcal{H}\left(1, 0, 1; \frac{1}{v_{321}}\right) + H(0; u_2) \mathcal{H}\left(1, 1, 1; \frac{1}{v_{123}}\right) - H(0; u_3) \mathcal{H}\left(1, 1, 1; \frac{1}{v_{123}}\right) - \\
& H(0; u_1) \mathcal{H}\left(1, 1, 1; \frac{1}{v_{231}}\right) + H(0; u_3) \mathcal{H}\left(1, 1, 1; \frac{1}{v_{231}}\right) + H(0; u_1) \mathcal{H}\left(1, 1, 1; \frac{1}{v_{312}}\right) - \\
& H(0; u_2) \mathcal{H}\left(1, 1, 1; \frac{1}{v_{312}}\right) - \frac{3}{2}\mathcal{H}\left(0, 0, 0, 1; \frac{1}{u_{123}}\right) - \frac{3}{2}\mathcal{H}\left(0, 0, 0, 1; \frac{1}{u_{231}}\right) - \\
& \frac{3}{2}\mathcal{H}\left(0, 0, 0, 1; \frac{1}{u_{312}}\right) - 3\mathcal{H}\left(0, 0, 0, 1; \frac{1}{v_{132}}\right) - 3\mathcal{H}\left(0, 0, 0, 1; \frac{1}{v_{213}}\right) - 3\mathcal{H}\left(0, 0, 0, 1; \frac{1}{v_{321}}\right) - \\
& \frac{1}{2}\mathcal{H}\left(0, 0, 1, 1; \frac{1}{u_{123}}\right) - \frac{1}{2}\mathcal{H}\left(0, 0, 1, 1; \frac{1}{u_{231}}\right) - \frac{1}{2}\mathcal{H}\left(0, 0, 1, 1; \frac{1}{u_{312}}\right) - \\
& \frac{1}{2}\mathcal{H}\left(0, 1, 0, 1; \frac{1}{u_{123}}\right) - \frac{1}{2}\mathcal{H}\left(0, 1, 0, 1; \frac{1}{u_{231}}\right) - \frac{1}{2}\mathcal{H}\left(0, 1, 0, 1; \frac{1}{u_{312}}\right) + \\
& \frac{1}{4}\mathcal{H}\left(0, 1, 1, 1; \frac{1}{v_{123}}\right) + \frac{1}{4}\mathcal{H}\left(0, 1, 1, 1; \frac{1}{v_{132}}\right) + \zeta_3 H(0; u_1) + \zeta_3 H(0; u_2) + \zeta_3 H(0; u_3) + \\
& \frac{5}{2}\zeta_3 H(1; u_1) + \frac{5}{2}\zeta_3 H(1; u_2) + \frac{5}{2}\zeta_3 H(1; u_3) + \frac{1}{2}\zeta_3 \mathcal{H}\left(1; \frac{1}{u_{123}}\right) + \frac{1}{2}\zeta_3 \mathcal{H}\left(1; \frac{1}{u_{231}}\right) + \\
& \frac{1}{2}\zeta_3 \mathcal{H}\left(1; \frac{1}{u_{312}}\right) - \frac{1}{2}\mathcal{H}\left(1, 0, 0, 1; \frac{1}{u_{123}}\right) - \frac{1}{2}\mathcal{H}\left(1, 0, 0, 1; \frac{1}{u_{231}}\right) - \frac{1}{2}\mathcal{H}\left(1, 0, 0, 1; \frac{1}{u_{312}}\right) + \\
& \frac{1}{4}\zeta_3 \mathcal{H}\left(1; \frac{1}{v_{123}}\right) + \frac{1}{4}\zeta_3 \mathcal{H}\left(1; \frac{1}{v_{132}}\right) + \frac{1}{4}\zeta_3 \mathcal{H}\left(1; \frac{1}{v_{213}}\right) + \frac{1}{4}\zeta_3 \mathcal{H}\left(1; \frac{1}{v_{231}}\right) + \frac{1}{4}\zeta_3 \mathcal{H}\left(1; \frac{1}{v_{312}}\right) + \\
& \frac{1}{4}\zeta_3 \mathcal{H}\left(1; \frac{1}{v_{321}}\right) + \frac{1}{4}\mathcal{H}\left(0, 1, 1, 1; \frac{1}{v_{213}}\right) + \frac{1}{4}\mathcal{H}\left(0, 1, 1, 1; \frac{1}{v_{231}}\right) + \frac{1}{4}\mathcal{H}\left(0, 1, 1, 1; \frac{1}{v_{312}}\right) + \\
& \frac{1}{4}\mathcal{H}\left(0, 1, 1, 1; \frac{1}{v_{321}}\right) + \frac{1}{4}\mathcal{H}\left(1, 0, 1, 1; \frac{1}{v_{123}}\right) + \frac{1}{4}\mathcal{H}\left(1, 0, 1, 1; \frac{1}{v_{132}}\right) + \frac{1}{4}\mathcal{H}\left(1, 0, 1, 1; \frac{1}{v_{213}}\right) + \\
& \frac{1}{4}\mathcal{H}\left(1, 0, 1, 1; \frac{1}{v_{231}}\right) + \frac{1}{4}\mathcal{H}\left(1, 0, 1, 1; \frac{1}{v_{312}}\right) + \frac{1}{4}\mathcal{H}\left(1, 1, 0, 1; \frac{1}{v_{123}}\right) + \frac{1}{4}\mathcal{H}\left(1, 1, 0, 1; \frac{1}{v_{132}}\right) + \\
& \frac{1}{4}\mathcal{H}\left(1, 1, 0, 1; \frac{1}{v_{213}}\right) + \frac{1}{4}\mathcal{H}\left(1, 1, 0, 1; \frac{1}{v_{231}}\right) + \frac{1}{4}\mathcal{H}\left(1, 1, 0, 1; \frac{1}{v_{312}}\right) + \\
& \frac{1}{4}\mathcal{H}\left(1, 1, 0, 1; \frac{1}{v_{321}}\right)
\end{aligned}$$

There should be a better presentation!

Otherwise we should abandon N=4 Yang Mills!



FullSimplify[ ] is not enough, we need a serious weapon to battle this 17-page





# How do we proceed?

Classical Polylogs

$$Li_k(z) = \int_0^z Li_{k-1}(t) d \log t \quad Li_1(z) = -\log(1-z)$$

Satisfy various identities

$$-Li_2\left(1 - \frac{1}{x}\right) = Li_2(1-x) + \frac{1}{2} \log(x)^2$$

Goncharov polylogs satisfy huge number of identities

$$G(a_k, a_{k-1}, \dots; z) = \int_0^z G(a_{k-1}, \dots; t) \frac{dt}{t - a_k}, \quad G(z) \equiv 1$$

How can we bring the identities under control?

**We will use motivic high tech!!!**

# Symbol of Transcendental Function

Goncharov

$$T_k \rightarrow S(T_k) = R_1 \otimes \cdots \otimes R_k$$

Symbol is an element of the k-fold tensor product of the multiplicative group of rational functions, it can be defined recursively

$$dT_k = \sum_i T_{k-1}^i d \log R_i \rightarrow S(T_k) = \sum_i S(T_{k-1}^i) \otimes R_i$$

$$R_1 \otimes (R_2 R_3) = R_1 \otimes R_2 + R_1 \otimes R_3$$

Properties

$$R_1 \otimes (cR_2) = R_1 \otimes R_2 \quad c = \text{const}$$

# Examples

Function	Differential	Symbol
$\log R$	$d \log R$	$R$
$\log R_1 \log R_2$	$\log R_1 d \log R_2 + \log R_2 d \log R_1$	$R_1 \otimes R_2 + R_2 \otimes R_1$
$Li_2(R)$	$-\log(1 - R) d \log R$	$-(1 - R) \otimes R$
$Li_k(R)$	$Li_{k-1}(R) d \log R$	$-(1 - R) \otimes R \cdots \otimes R$

# Why is the Symbol Very Useful?

Symbol converts polylog functional equations into rational function identities



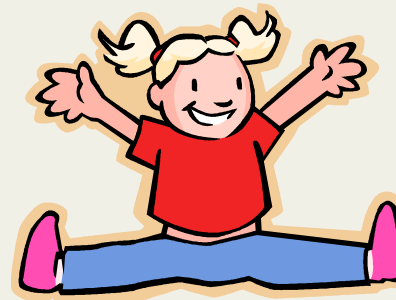
◆  $Li_2(x) + Li_2(-x) = \frac{1}{2}Li_2(x^2)$

$-(1-x) \otimes x - (1+x) \otimes (-x) = -(1-x^2) \otimes x = -\frac{1}{2}(1-x^2) \otimes x^2$

◆  $Li_2(z) + Li_2(1-z) + \log(z) \log(1-z) = \frac{\pi^2}{6}$   
 $-z \otimes (1-z) - (1-z) \otimes z + z \otimes (1-z) + (1-z) \otimes z = 0$

Symbol fixes only the leading functional transcendental piece

# Final Result



Goncharov, Spradlin, Vergu, AV '10

Two-loop six-point MHV remainder function is

$$R_6^{(2)}(u_1, u_2, u_3) = \sum_{i=1}^3 (L_4(x_i^+, x_i^-) - \frac{1}{2} Li_4(1-1/u_i)) - \frac{1}{8} \left( \sum_{i=1}^3 Li_2(1-1/u_i) \right)^2 + \frac{1}{24} J^4 + \frac{\pi^2}{12} J^2 + \frac{\pi^4}{72}.$$

$$L_4(x^+, x^-) = \sum_{m=0}^3 \frac{(-1)^m}{(2m)!!} \log(x^+ x^-)^m (\ell_{4-m}(x^+) + \ell_{4-m}(x^-)) + \frac{1}{8!!} \log(x^+ x^-)^4$$

$$\ell_n(x) = \frac{1}{2} (Li_n(x) - (-1)^n Li_n(1/x))$$

$$J = \sum_{i=1}^3 (\ell_1(x_i^+) - \ell_1(x_i^-))$$

$$x_i^\pm = u_i x^\pm$$

$$x^\pm = \frac{u_1 + u_2 + u_3 - 1 \pm \sqrt{\Delta}}{2u_1 u_2 u_3}$$

$$\Delta = (u_1 + u_2 + u_3 - 1)^2 - 4u_1 u_2 u_3$$

Bartels, Prygarin, Lipatov '10: analyzed Regge limit  
Gaiotto, Maldacena, Sever, Veira '11: OPE approach

# What's next?

- Our formula provides hope to the idea that we might be able to really unlock the secrets of multi-loop SYM amplitudes, and connect to strong coupling **Alday, Gaiotto, Maldacena, Sever, Veira**
- We need a technology for writing down symbols without first evaluating integrals at all
- It would be nice to have a derivation from integrals: work on two-loops is in progress, but solved for one-loop higher dimensional polygons

# More Motivic Magic

Spradlin, AV 05/11

- Consider **one-loop  $2m$ -gon in  $2m$ -dimensions**:  
appear in dim regulated one-loop MHV & can also be related to higher-loop four-dimensional integrals
- One-loop hexagon in six dimensions have been evaluated for massless & three-mass case:  
surprisingly the same  $L_3$  function appears.....

Dixon, Drummond, Henn; Del Duca, Duhr, Smirnov 05/11

- It turns out its symbol for **any  $m$**  can be read off from theorem by **Goncharov'96** on mixed Tate motives

# Recursions for Symbols

- After Feynman parameterization, any one-loop  $2m$ -gon integral in  $D=2m$  takes the form

$$F(Q) = \int_{\mathbb{CP}^{2m-1}} \frac{D^{2m-1} W}{(W \cdot Q \cdot W)^m} \quad Q_{ij} = (p_i + \cdots + p_j)^2$$

in terms of a quadric  $Q$  in  $\mathbb{CP}^{2m-1}$

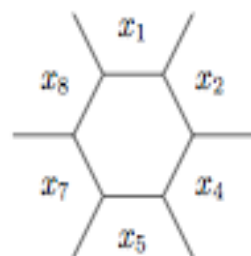
- The symbol of the integral is given recursively by application of a result on mixed Tate motives [Goncharov 96](#)

$$S_m(Q) = \sum_{i < j} S_{m-1}(\overline{Q}_{ij}) \otimes \frac{Q_{ij}^{-1} - \sqrt{(Q_{ij}^{-1})^2 - Q_{ii}^{-1} Q_{jj}^{-1}}}{Q_{ij}^{-1} + \sqrt{(Q_{ij}^{-1})^2 - Q_{ii}^{-1} Q_{jj}^{-1}}}$$

$(\overline{Q} = Q$  with rows and columns  $i, j$  deleted) [Spradlin, AV 05/11](#)



# Example: 1-loop 3-mass hexagon



$$Q = \begin{pmatrix} 0 & 0 & x_{14}^2 & x_{15}^2 & x_{17}^2 & x_{18}^2 \\ 0 & 0 & x_{24}^2 & x_{25}^2 & x_{27}^2 & x_{28}^2 \\ x_{14}^2 & x_{24}^2 & 0 & 0 & x_{47}^2 & x_{48}^2 \\ x_{15}^2 & x_{25}^2 & 0 & 0 & x_{57}^2 & x_{58}^2 \\ x_{17}^2 & x_{27}^2 & x_{47}^2 & x_{57}^2 & 0 & 0 \\ x_{18}^2 & x_{28}^2 & x_{48}^2 & x_{58}^2 & 0 & 0 \end{pmatrix}$$

$$(1 + g + g^2) \left[ S_2(Q_{12}) \otimes \frac{\langle 2358 \rangle \langle 12\bar{5} \cap \bar{8} \rangle}{\langle 1258 \rangle \langle 23\bar{5} \cap \bar{8} \rangle} \right. \\ + S_2(Q_{13}) \otimes \frac{\langle 1238 \rangle \langle 2568 \rangle \langle 5789 \rangle \langle 23\bar{5} \cap \bar{8} \rangle}{\langle 2358 \rangle \langle 2789 \rangle \langle 4568 \rangle \langle 56\bar{2} \cap \bar{8} \rangle} \\ + S_2(Q_{14}) \otimes \frac{\langle 2358 \rangle \langle 2789 \rangle \langle 4568 \rangle \langle 45\bar{2} \cap \bar{8} \rangle}{\langle 1238 \rangle \langle 2458 \rangle \langle 5789 \rangle \langle 23\bar{5} \cap \bar{8} \rangle} \\ + S_2(Q_{16}) \otimes \frac{\langle 2358 \rangle \langle 2456 \rangle \langle 5789 \rangle \langle 78\bar{2} \cap \bar{5} \rangle}{\langle 1235 \rangle \langle 2578 \rangle \langle 4568 \rangle \langle 23\bar{5} \cap \bar{8} \rangle} \\ \left. + S_2(Q_{24}) \otimes \frac{\langle 1238 \rangle \langle 2458 \rangle \langle 5789 \rangle \langle 12\bar{5} \cap \bar{8} \rangle}{\langle 1258 \rangle \langle 2789 \rangle \langle 4568 \rangle \langle 45\bar{2} \cap \bar{8} \rangle} \right]$$

$$S_2(Q) = x_{ii}^2 x_{jk}^2 \otimes \frac{-x_{ii}^2 x_{jk}^2 + x_{ik}^2 x_{jl}^2 + x_{ij}^2 x_{kl}^2 - \Delta}{-x_{ii}^2 x_{jk}^2 + x_{ik}^2 x_{jl}^2 + x_{ij}^2 x_{kl}^2 + \Delta} \\ + x_{ik}^2 x_{jl}^2 \otimes \frac{+x_{ii}^2 x_{jk}^2 - x_{ik}^2 x_{jl}^2 + x_{ij}^2 x_{kl}^2 - \Delta}{+x_{ii}^2 x_{jk}^2 - x_{ik}^2 x_{jl}^2 + x_{ij}^2 x_{kl}^2 + \Delta} \\ + x_{ij}^2 x_{kl}^2 \otimes \frac{+x_{ii}^2 x_{jk}^2 + x_{ik}^2 x_{jl}^2 - x_{ij}^2 x_{kl}^2 - \Delta}{+x_{ii}^2 x_{jk}^2 + x_{ik}^2 x_{jl}^2 - x_{ij}^2 x_{kl}^2 + \Delta}$$

$$\Delta = \sqrt{+\det Q}$$

$$g : i \rightarrow i + 3$$

$$\langle ij\bar{k} \cap \bar{l} \rangle = \langle i \ k-1 \ k \ k+1 \rangle \langle j \ l-1 \ l \ l+1 \rangle - \langle j \ k-1 \ k \ k+1 \rangle \langle i \ l-1 \ l \ l+1 \rangle$$

# Higher Loops and Points Symbols

- Similar story should exist for higher loops: recursions for symbols? Geometrical meaning?
- Higher points: **Caron-Huot** Wilson loop approach to symbol for all n MHV 2-loops
- Higher loops: **Dixon, Drummond, Henn** ansatz for a symbol for 6-points 3-loops from OPE and other constraints, **assuming the same cross-ratios as at 2-loops**
- Non-MHV ratio function: **Dixon, Drummond, Henn**

Deriving function from symbol requires work.

This is a very nontrivial step.

Duhr, Gangl, Rhodes

# Two-loop All-n MHV Symbol

The differential of the n-point function is expressed as

$$dR_n = \sum_{i,j} C_{i,j} d \log \langle i-1 \ i \ i+1 \ j \rangle \quad (\text{A.1})$$

where  $C_{2,i}$  is the sum of the four contributions

Caron-Huot

$$\begin{aligned} C_{2,i}^{(1)} &= \log u_{2,i-1,i,1} \times \sum_{j=2}^{i-1} \sum_{k=i}^{n+1} \left[ \text{Li}_2(1 - u_{j,k,k-1,j+1}) + \log \frac{x_{j,k}^2}{x_{j+1,k}^2} \log \frac{x_{j,k}^2}{x_{j,k-1}^2} \right], \\ C_{2,i}^{(2)} &= \sum_{j=4}^{i-2} \Delta(1, 2; j-1, j; i-1, i), \\ C_{2,i}^{(3)} &= \sum_{j=i+2}^n \Delta(2, 1; j, j-1; i, i-1), \\ C_{2,i}^{(4)} &= -2\text{Li}_3(1 - \frac{1}{u}) - \text{Li}_2(1 - \frac{1}{u}) \log u - \frac{1}{6} \log^3 u + \frac{\pi^2}{6} \log u, \end{aligned} \quad (\text{A.2})$$

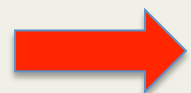
and other  $C_{i,j}$  are obtained by cyclic symmetry. In the first line,  $x_{j+1} \equiv x_2$  when  $j = i-1$ , and  $x_{k-1} \equiv x_1$  when  $k = i$ , and in the last line,  $u = u_{2,i-1,i,1}$ . The symbol of  $\Delta$  is

$$\begin{aligned} & \mathcal{S}\Delta(1, 2; j-1, j; i-1, i) \\ &= \left( \mathcal{S}[I_5(i; 1, 2; j-1, j)] \otimes \frac{\langle ii+1(\bar{2}) \cap (\bar{j}) \rangle \langle 23ij \rangle}{\langle j-1jj+1i \rangle \langle 123j \rangle \langle 23ii+1 \rangle} - ((ii+1) \rightarrow (i-1i)) \right) \\ &+ \left( \begin{aligned} & \frac{1}{2} \mathcal{S}[\text{Li}_2(1 - u_{j,2,1,i-1}) - \text{Li}_2(1 - u_{j,2,1,i})] \otimes \left( \frac{\langle 123i \rangle \langle j-1jj+12 \rangle \langle 23ij \rangle}{\langle 123j \rangle \langle j-1jj+1i \rangle \langle 23ii+1 \rangle} \right)^2 \frac{\langle jj+1(\bar{2}) \cap (\bar{i}) \rangle \langle ii+1jj+1 \rangle}{\langle 2i jj+1 \rangle \langle 13(2i-1i) \cap (2jj+1) \rangle} \\ & + \frac{1}{2} \mathcal{S}[\text{Li}_2(1 - u_{j,i-1,i,2}) - \text{Li}_2(1 - u_{j,i-1,i,1})] \otimes \left( \frac{\langle 12i-1i \rangle \langle 23ij \rangle}{\langle 123i \rangle \langle i-1ii+1j \rangle \langle 23ii-1i \rangle} \right)^2 \frac{\langle jj+1(\bar{2}) \cap (\bar{i}) \rangle \langle i-1i+1(i23) \cap (ijj+1) \rangle}{\langle 2i jj+1 \rangle \langle 12jj+1 \rangle} \\ & + \frac{1}{2} \mathcal{S}[\text{Li}_2(1 - u_{2,i-1,i,1})] \otimes \frac{\langle jj+1(\bar{2}) \cap (\bar{i}) \rangle \langle i-1i+1(i23) \cap (ijj+1) \rangle}{\langle 2i jj+1 \rangle \langle 13(2i-1i) \cap (2jj+1) \rangle} \\ & + \frac{1}{2} \mathcal{S}[\log u_{j,i-1,i,2} \log u_{j,2,1,i-1}] \otimes \left( \frac{\langle 23ij \rangle}{\langle 123j \rangle} \right)^2 \frac{\langle jj+1(\bar{2}) \cap (\bar{i}) \rangle \langle 13(2i-1i) \cap (2jj+1) \rangle}{\langle 2i jj+1 \rangle \langle 23i-1i \rangle \langle i-1i+1(i23) \cap (ijj+1) \rangle} \\ & - ((jj+1) \rightarrow (j-1j)) \end{aligned} \right) \\ &+ \mathcal{S}[I_5(1; i-1, i; j-1, j)] \otimes \frac{\langle 12ij \rangle \langle 23i-1i \rangle}{\langle 12i-1i \rangle \langle 23ij \rangle} \\ &+ \mathcal{S}[\log u_{i,j-1,j,1} \log u_{2,i-1,i,1}] \otimes \frac{\langle j-1j+1(j12) \cap (jii+1) \rangle \langle 123i \rangle \langle 23i-1i \rangle}{\langle 123j \rangle \langle j-1jj+1i \rangle \langle 12i-1i \rangle \langle 23ii+1 \rangle}. \end{aligned} \quad (\text{A.3})$$

# How can we integrate it?

- Constructing the function which has the same symbol is nontrivial. We need luck/guessing.

$$(1 + x + x^2) \otimes x = \frac{1-x^3}{1-x} \otimes x = \frac{1}{3}(1 - x^3) \otimes x^3 - (1 - x) \otimes x$$


$$-\frac{1}{3}Li_2(x^3) + Li_2(x)$$

- In 2-loop 6-point case we had Goncharov's theorem to guide us  $A \wedge A = 0 \rightarrow$  classical Li's only
- More generally we can have for degree 4:

$$Li_{2,2}(x, y) = \sum_{0 < n < m} \frac{x^n}{n^2} \frac{y^m}{m^2} \qquad Li_4(x) = \sum_{n=1}^{\infty} \frac{x^n}{n^2}$$

# Motivic Content

$$S = \sum_i c_i^L \otimes \Delta_i = \sum_i \Delta_i \otimes c_i^R$$

$$c_i^L - c_i^R = \sum_j b_{ij}^L \otimes \Delta_j = \sum_j \Delta_j \otimes b_{ij}^R$$

Motivic content is the antisymmetrization of

$$\sum_j (b_{ij}^L - b_{ij}^R) \otimes \Delta_j$$

with Goncharov, Spradlin, Vergu

This operation kills all terms except for Li<sub>4</sub> and Li<sub>2,2</sub>

# Examples

$$\Delta = x : Li_4(x) \rightarrow -(1-x) \wedge x \otimes x$$

$$\Delta = 1 - xy : Li_{2,2}(x, y) \rightarrow -2(1-y) \wedge y \otimes y + 2(1-x) \wedge x \otimes x$$

$$\Delta = x; \Delta = y : Li_{2,2}(x, y) \rightarrow -(1-R) \wedge R \otimes R$$

- Motivic content will help identify good set of cross-ratios out of 2310 for 7-points
- We also found new 40-term  $Li_3$  identities

$$\sum a_i (1 - R_i) \wedge R_i \otimes R_i = 0$$

with Goncharov, Spradlin, Vergu

# Regge limit

- While hard to evaluate the amplitude in general, it can be done in the Regge limit. For  $n=7$ , this is

$$u_{26} \sim 1 - \mathcal{O}(\epsilon^2); \quad u_{25}, u_{36} \sim 1 - \mathcal{O}(\epsilon); \quad u_{ij} \sim \mathcal{O}(\epsilon)$$

- The leading log of the imaginary part is  $\log \epsilon F_2$

$$F_2 = \sum_{i=4}^5 \log \frac{|\beta_3|^2 |\beta_i - \beta_6|^2}{|\beta_i|^2 |\beta_3 - \beta_6|^2} \log \frac{|\beta_i|^2 |\beta_3 - \beta_6|^2}{|\beta_6|^2 |\beta_3 - \beta_i|^2}$$

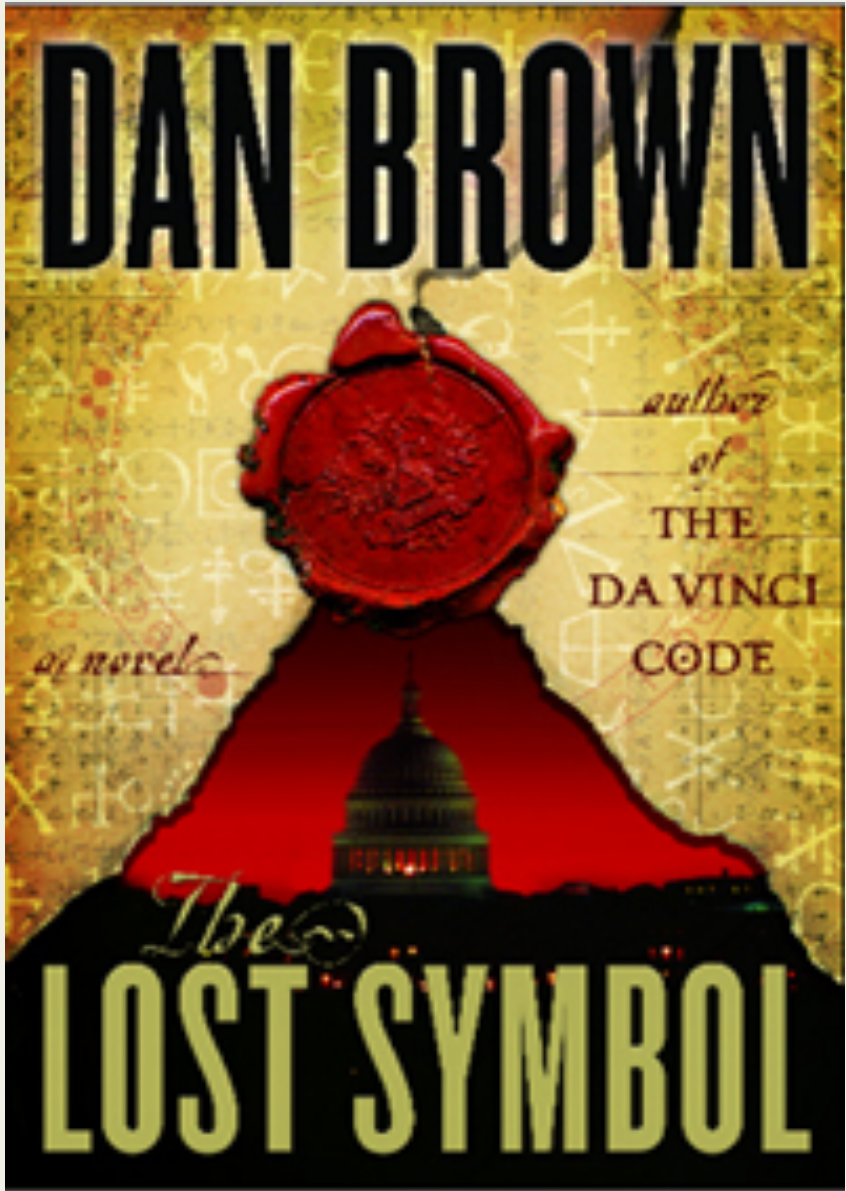
Computed by Prygarin using BFKL theory.  $\beta_i = \sum_{k=3} \vec{p}_k$

- Regge limit of the discontinuity  $u_{26} \rightarrow u_{26} e^{2\pi i}$  of Simon's symbol reproduces this result exactly.  
with Spradlin, Prygarin, Vergu
- More legs and loops in Regge limit ?

# Conclusions

- If a problem is too hard, look for something simpler. The symbol is a useful stepping stone half way between an integrand and its integral.
- We need a technology for writing down symbols without first evaluating integrals at all.
- We need better technology to find functions from symbols, **motivic content** is very helpful in addressing this question.
- Symbols are not the ultimate goal, but serve the simplest incarnation of a deeper **motivic structure of amplitudes** which we are only beginning to see.





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