Cool Cores, Conduction, and Virial Shocks: Sculpting Cosmic Gas into Clusters

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Outline

"Sculpting Cosmic Gas into Clusters"

Central Density

- * Thermal Instability
- * Implications for Cool Cores
- * Non-Self-Similarity

(McCourt, Sharma, et al. 2012) (Sharma, McCourt, et al. 2012) (Sharma, McCourt, et al., submitted)

Outer Temperature *McCourt, Quataert, & Parrish, in prep.*

- * Accretion History $\rightarrow T(r)$
- * Conduction + MTI

What determines the **density** of the gas in the **centers** of clusters?

Motivation: Non-Self-Similarity

Assume that the gas properties scale with the dark matter:

* $\rho \sim M^0$

- * $T \sim M/r \sim M^{2/3}$
- * $L \sim \rho^2 T^{1/2} r^3 \sim T^2$



Gas in the centers of clusters has *lower density* and *higher entropy* than gravitational self-similar models predict.

Background: Thermal Instability in Clusters



Is the ICM Thermally Unstable?

- * Thermal instability suppressed in cooling-flows (Balbus & Soker 1989)
- * Even some equilibria may be thermally stable (Kunz et al. 2010)
- * Many heating mechanisms are thermally unstable (Gaspari et al. 2012)
- * Multi-phase gas seen in many clusters (McDonald et al. 2010, 2011)

Assume local thermal instability

Background: The Face of Thermal Instability (Assuming it exists...)



Thermal Instability does not necessarily imply Multi-phase gas.

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Thermal Instability does **not necessarily** imply Multi-phase gas. See cold gas when $t_{cool}/t_{ff} \lesssim 10$

Background: Physics of the Saturation



- * Perturbations initially grow exponentially...
- *Saturate when $t_{sink} \sim t_{cool}$
- * Final amplitude $\propto (t_{\rm cool}/t_{\rm ff})^{-2}$

* This is a non-linear effect

Background: Feedback Regulation

Feedback and cooling self-regulate to the critical threshold for non-linear thermal stability: $\min(t_{cool}/t_{ff}) \sim 10$

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cf. Voit et al. (2001)

Results: Core Size w/ Mass



High-Mass Halos

- * High Temperature
- * Long Cooling Time
- $* \Rightarrow$ small core

Low-Mass Halos

- * Lower Temperature
- * Shorter Cooling Time
- $* \Rightarrow$ large core

(Minimum) Core size determined by Thermal Instability

Applications



Also:

- * gas fraction
- * core size & entropy
- stellar mass
- * baryon fraction

Assuming global thermal balance, these properties are ~independent of the feedback mechanism

What determines the temperature of the gas at large radii in clusters?

Motivation



*
$$\nabla T < 0 \Rightarrow MTI$$

* $t_{cond} \sim r_{vir}^2 / \chi \sim 1 \text{ Gyr} \lesssim t_{age}$

Motivation



In simulations of **isolated** clusters, the ICM becomes isothermal after ~Gyr.

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Model: Entropy Generation at the Shock



$$\frac{1}{2}v_i^2 = \frac{GM_{\rm sh}}{r_{\rm sh}} - \frac{GM_{\rm sh}}{r_{\rm ta}}$$
$$\rho v = \frac{1}{4\pi r_{\rm sh}^2} \frac{dM}{dt}$$

+ Jump Conditions $\rightarrow K(r_{sh})$

(e.g. Voit et al. 2003)

Results: Adiabatic Evolution



In the case of **Adiabatic Evolution**, this is a simple problem.

*
$$T(0) = T_{vir}$$

* $T(out) = T_s$

Temperature Gradient set by Accretion Rate

 $(t_{\rm dyn} \times d \ln M/dt)$

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Dispersion in Accretion Histories



Accretion histories from McBride et al. 2009

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Effect of Conduction



Accretion histories from McBride et al. 2009

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Application: MTI & Non-Thermal Pressure Support



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Conclusions

Density Cores

- * Assuming that the ICM is thermally unstable, multiphase gas forms only when $t_{\rm cool}/t_{\rm ff} \lesssim 10$.
- * Cooling and feedback selfregulate to the critical threshold for stability.
- * This sets a density ceiling (or entropy floor) for the gas
 ⇒ non-self-similarity.

Temperature Gradients

- * A cluster's accretion rate determines its temperature gradient
- * This temperature gradient persists even with thermal conduction
- * Importance of the MTI may be non-monotonic with halo mass. $(3 \times 10^{14} M_{\odot})$ is the sweet spot).