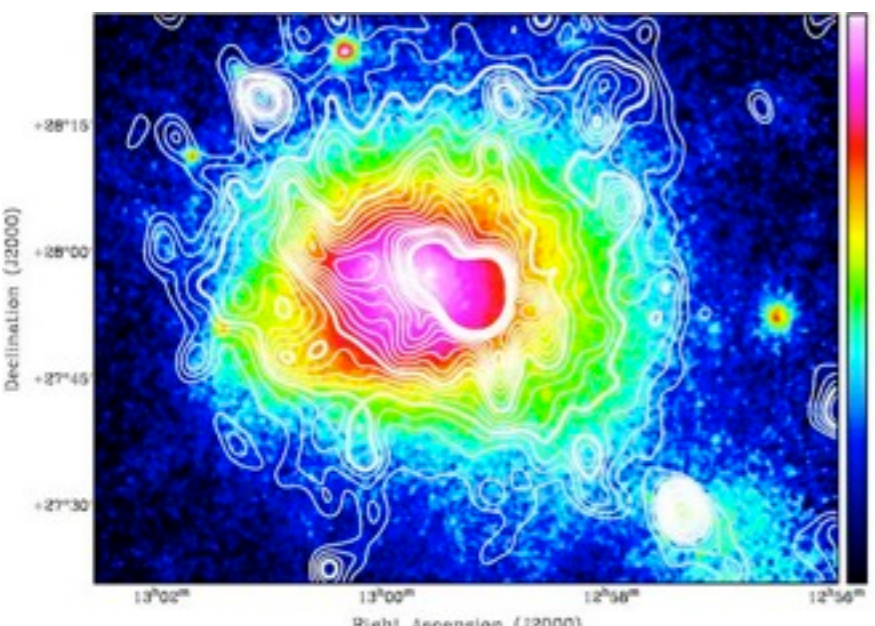


Turning Radio Relics **On**

and

Radio Halos **Off**



Peng Oh  
UCSB

# The Switches



Josh Wiener (UCSB)



Anders Pinzke  
(UCSB)



Fulai Guo  
(Santa Cruz)



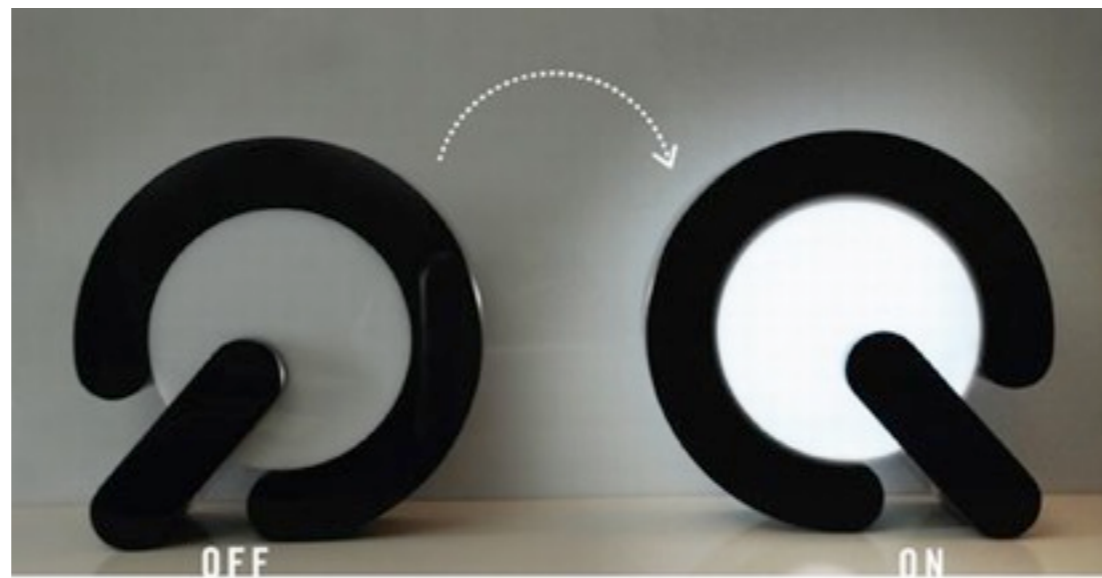
Christoph Pfrommer  
(Heidelberg)



# Why you should care



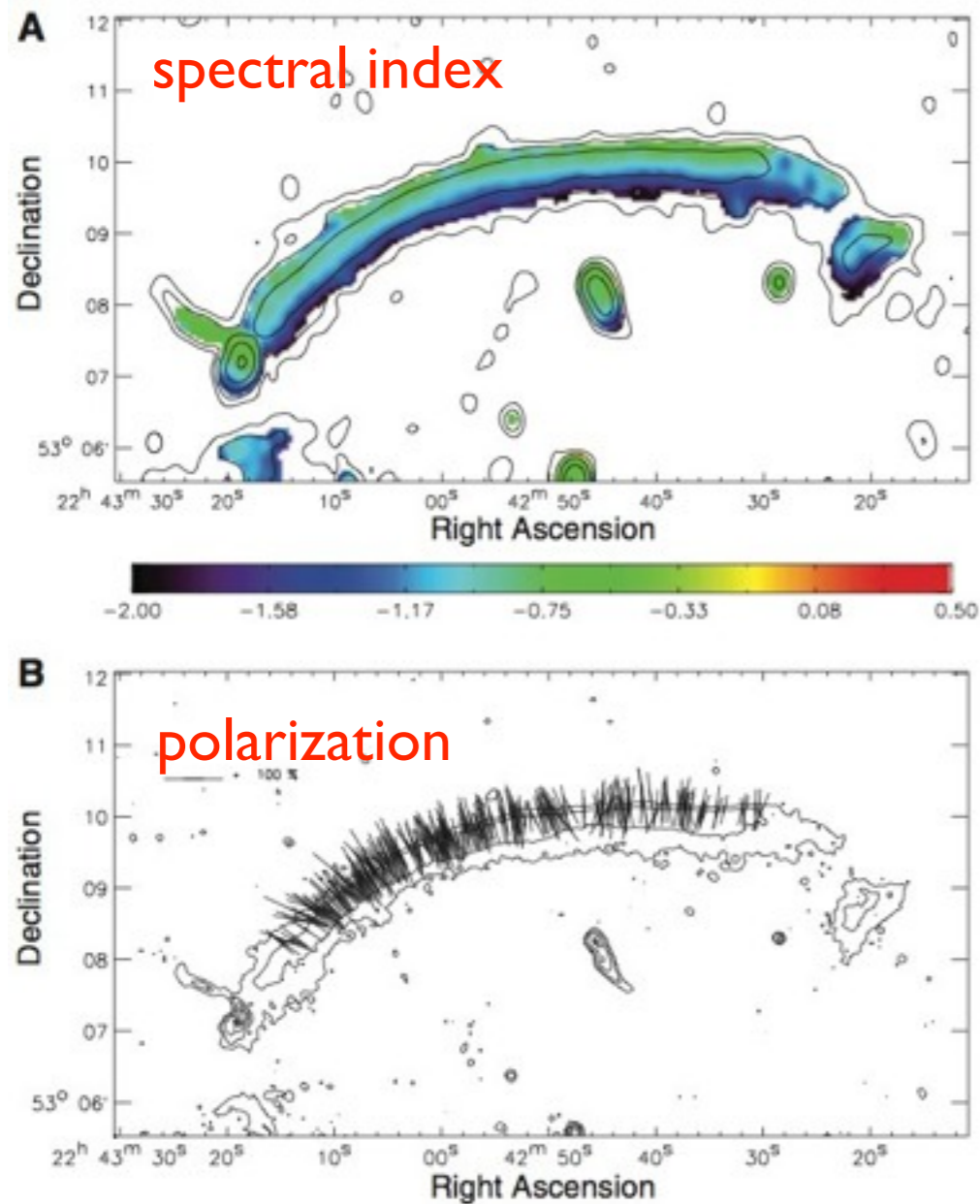
LOFAR could really shake things up!



# Turning Radio Relics On with Fossil Electrons

Pinzke, Oh, Pfrommer, 2012, in prep

# Radio Relics: Awesome!



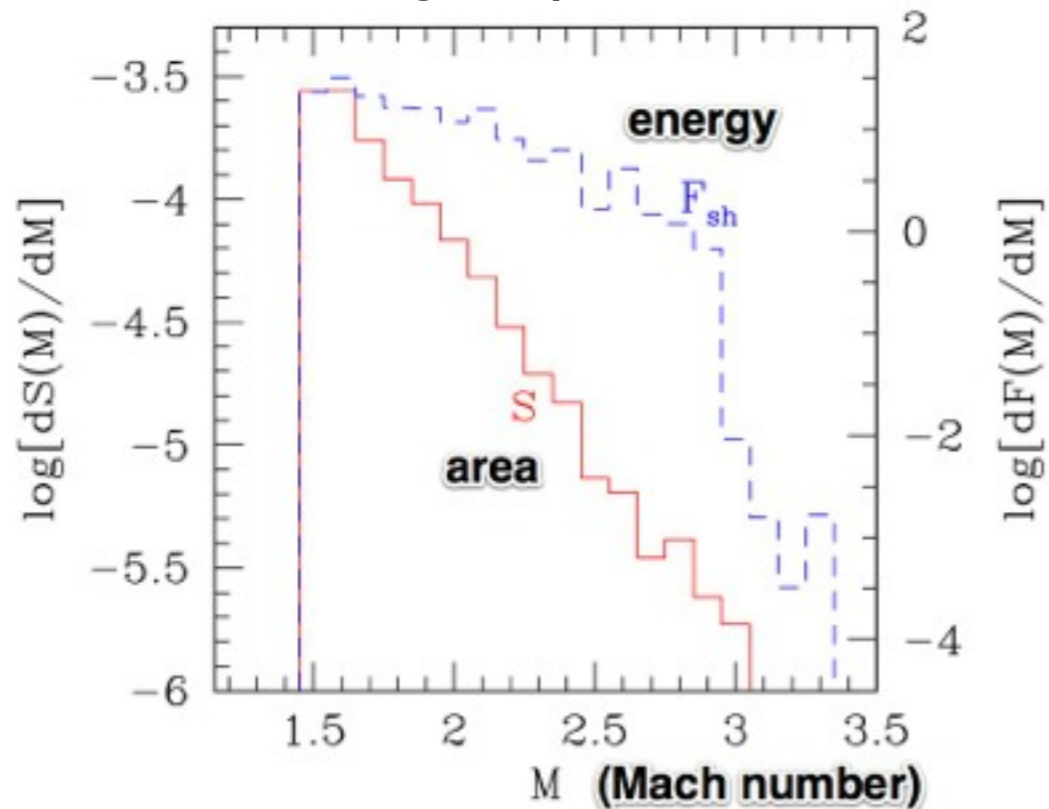
van Weeren et al, 2010

- Trace shocks in cluster outskirts
- Spectral index: shock Mach number
- Spectral ageing: B-field strength
- Polarization: B-field orientation

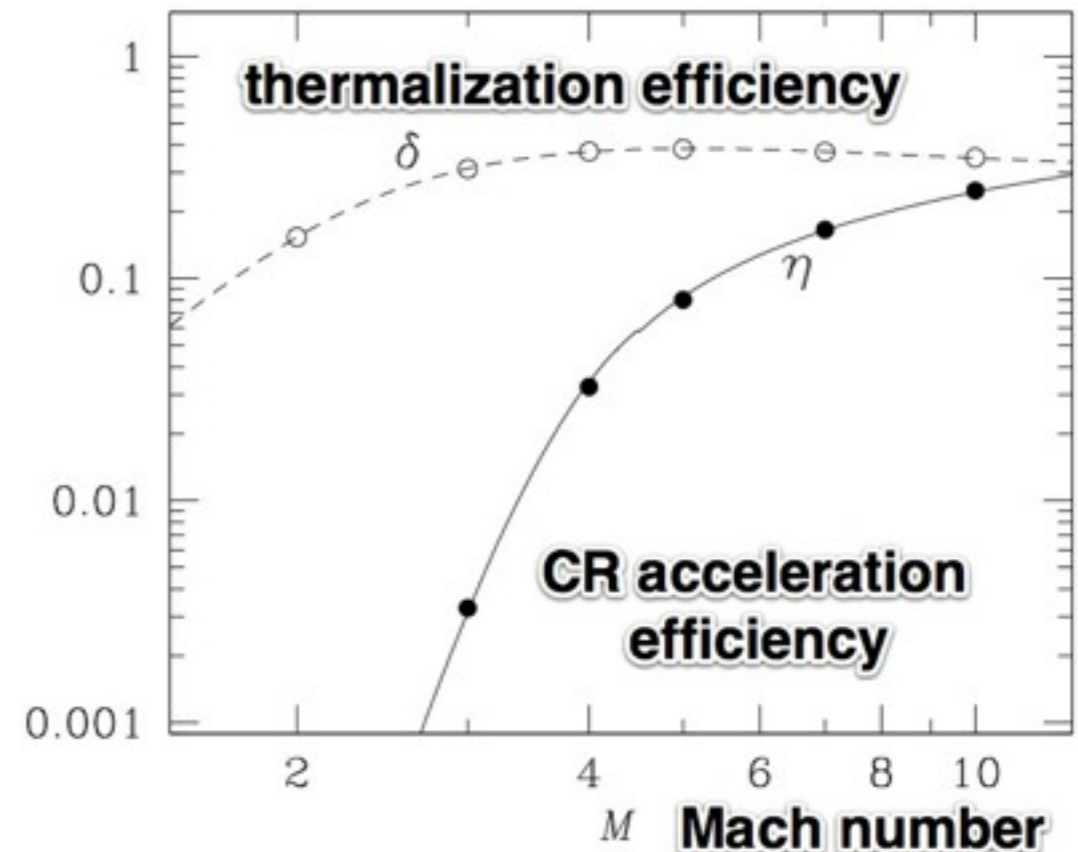


# Biggest unknown: shock acceleration efficiency

Kang & Ryu 2011



Bruggen et al 2012

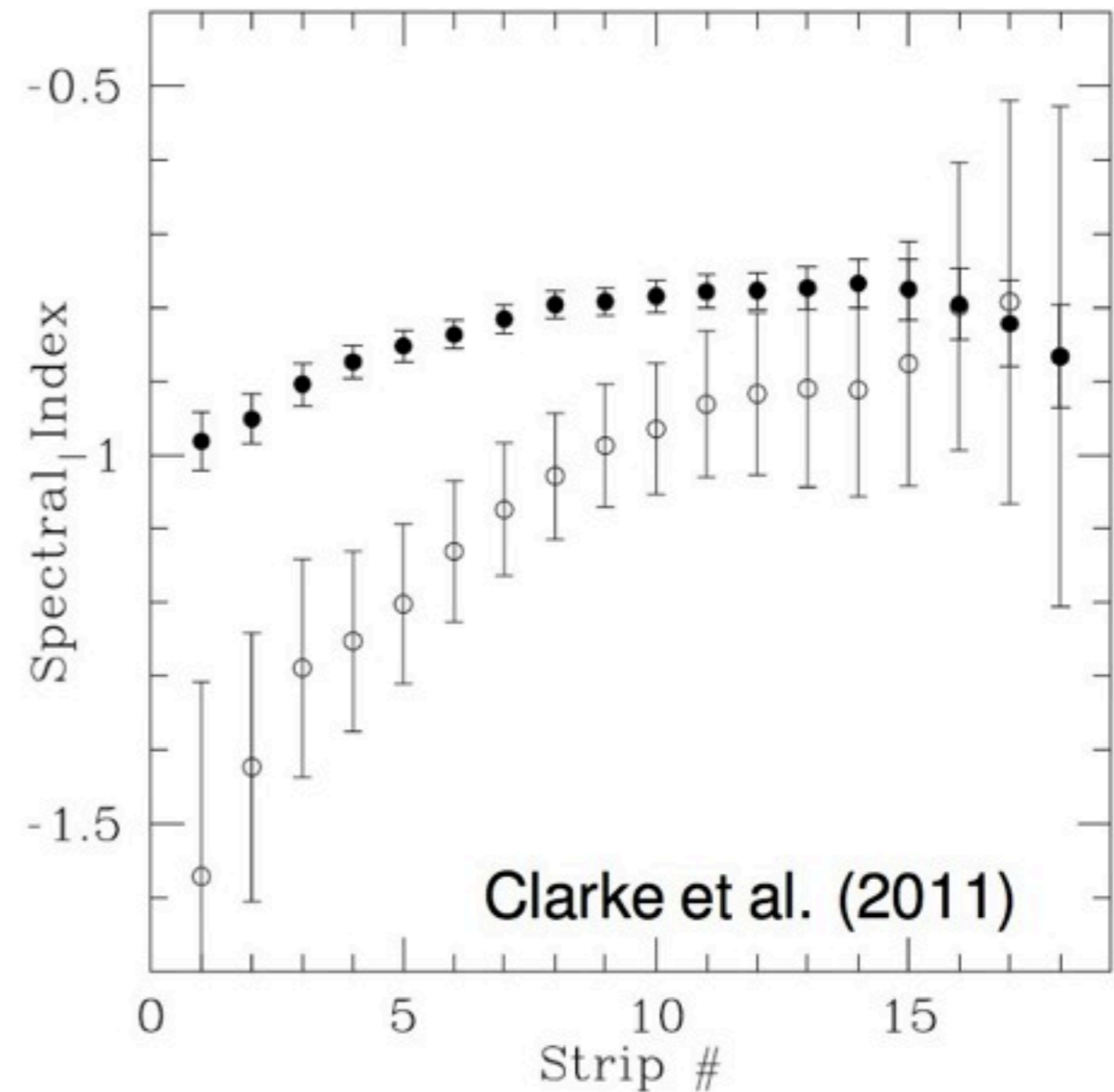
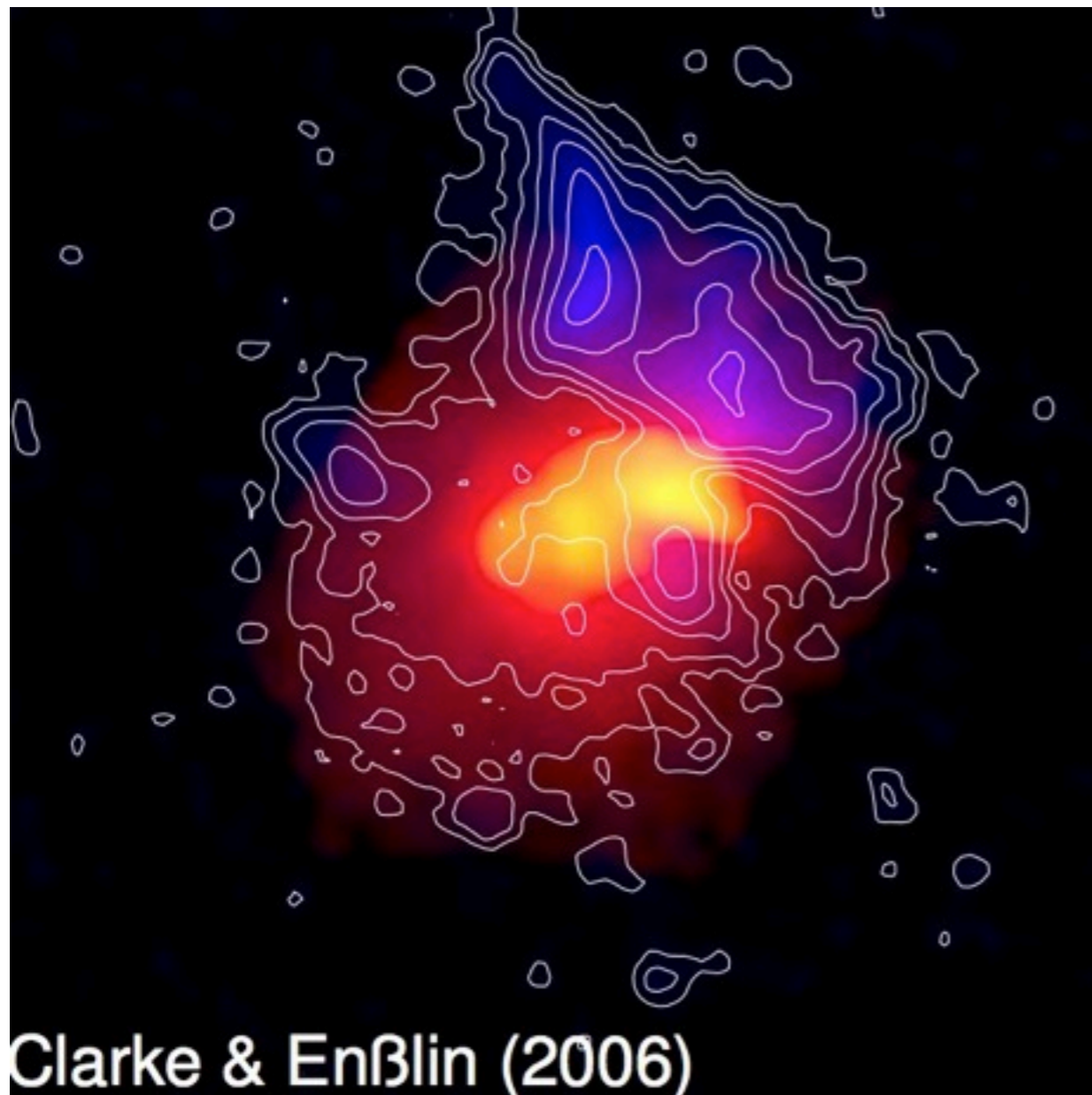


Outskirts dominated by low Mach number shocks  
These shocks have low acceleration efficiency

How many will LOFAR see?



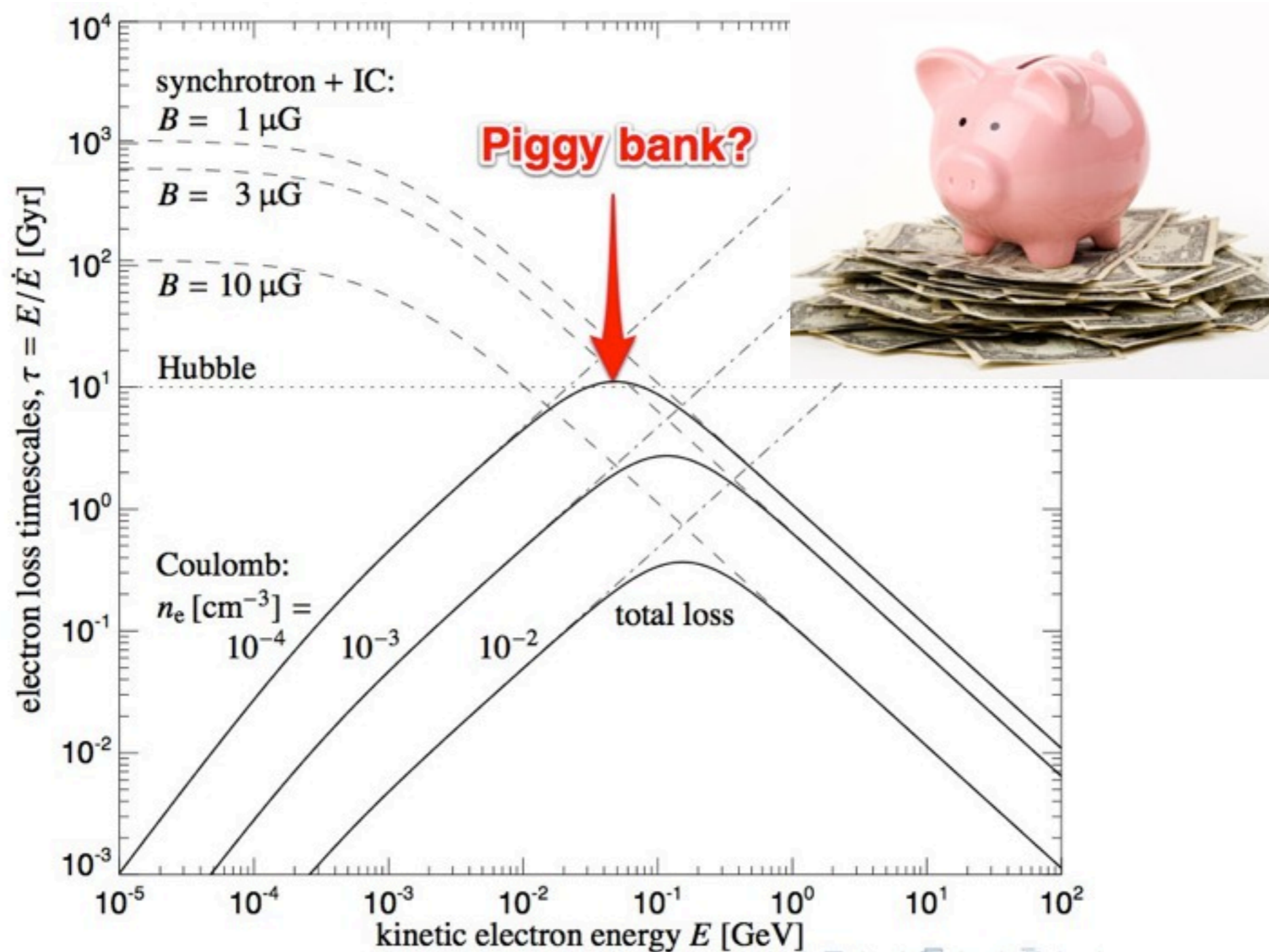
# A poster child: A2256



$$\alpha_{\nu} = 0.85 \quad \rightarrow \quad \mathcal{M} = 2.6:$$

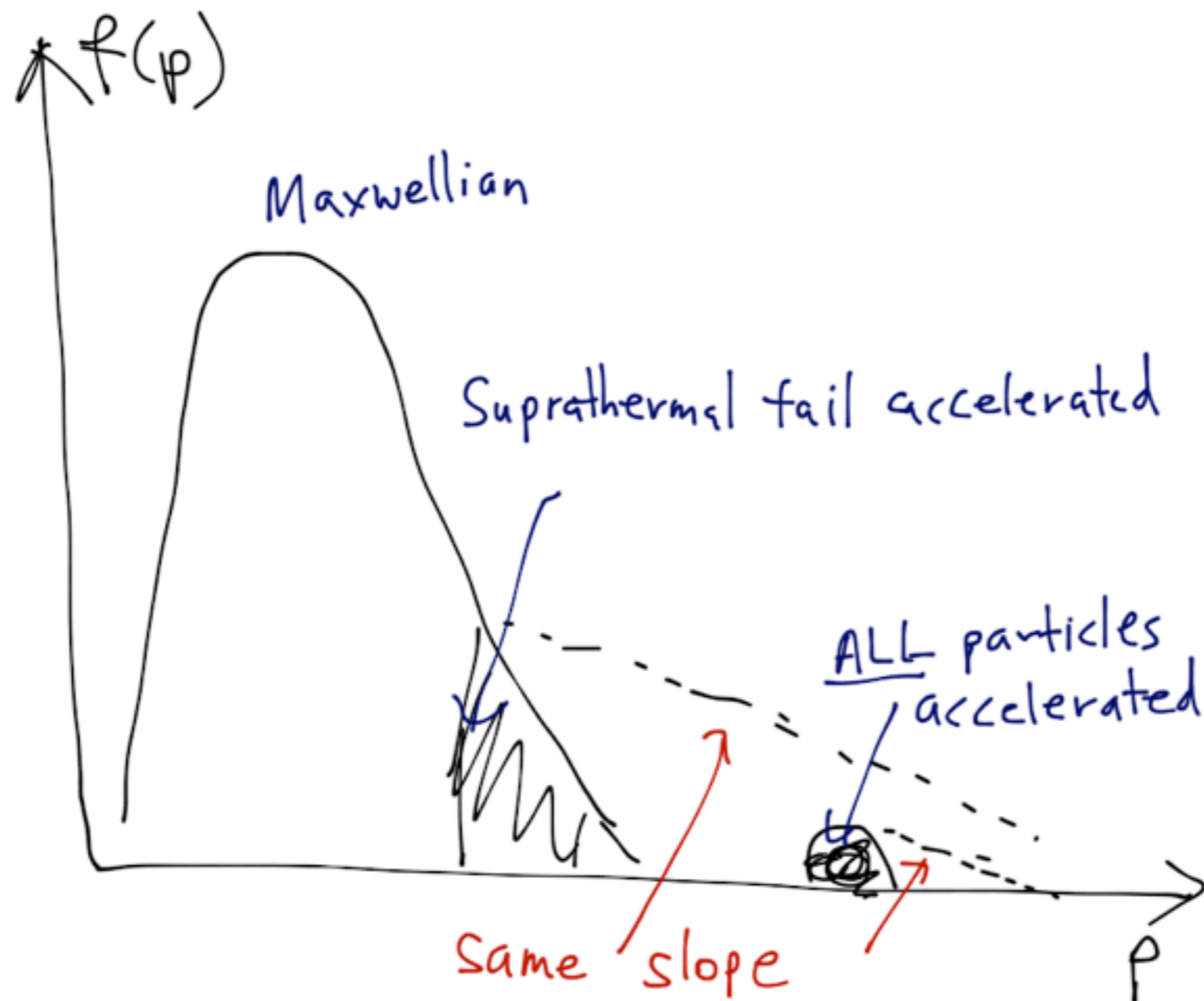
How is this possible?

# Could fossil electrons be responsible?





# How would relic reacceleration work?



# Could this beat direct injection?



Ultimately, we care about  $f(p_{\text{emit}})$  where  $p_{\text{emit}} \approx 10 \text{ GeV}$

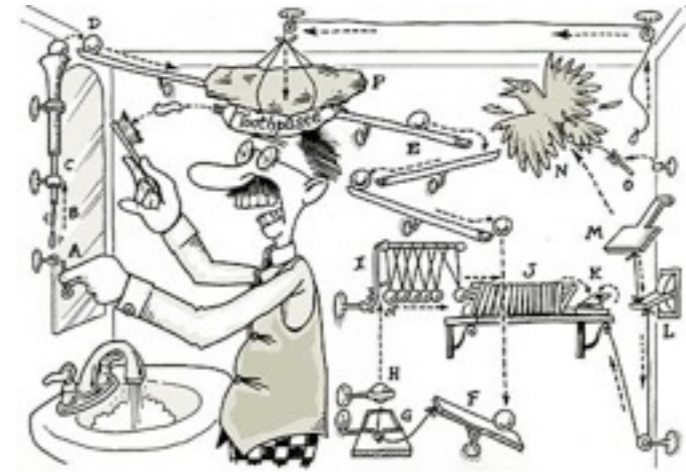
Can show that this is approximately:

$$f_2(p_{\text{emit}}) \approx q \frac{n_o(p > p_*)}{4\pi p_*^3} \left( \frac{p_{\text{emit}}}{p_*} \right)^{-q}$$

where  $p_*$  is where number density peaks

So just need to figure out  $p_*, n(p > p_*)$

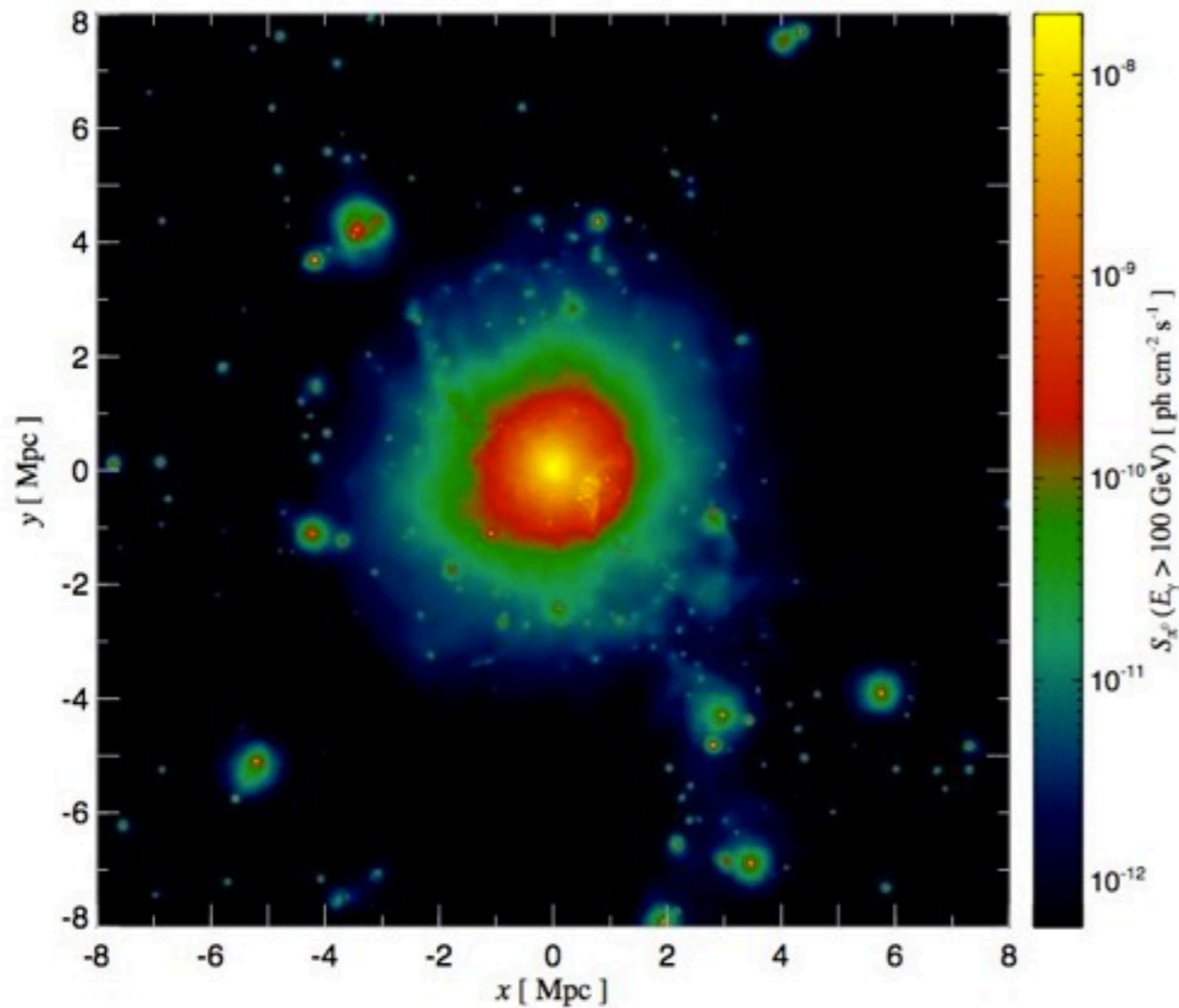
# Method



Use previous CR proton  
sims: (time resolution:  
 $\sim 100$  Myr, 1 Gyr)

Now inject CR electrons  
with  $e/p \sim 0.01$

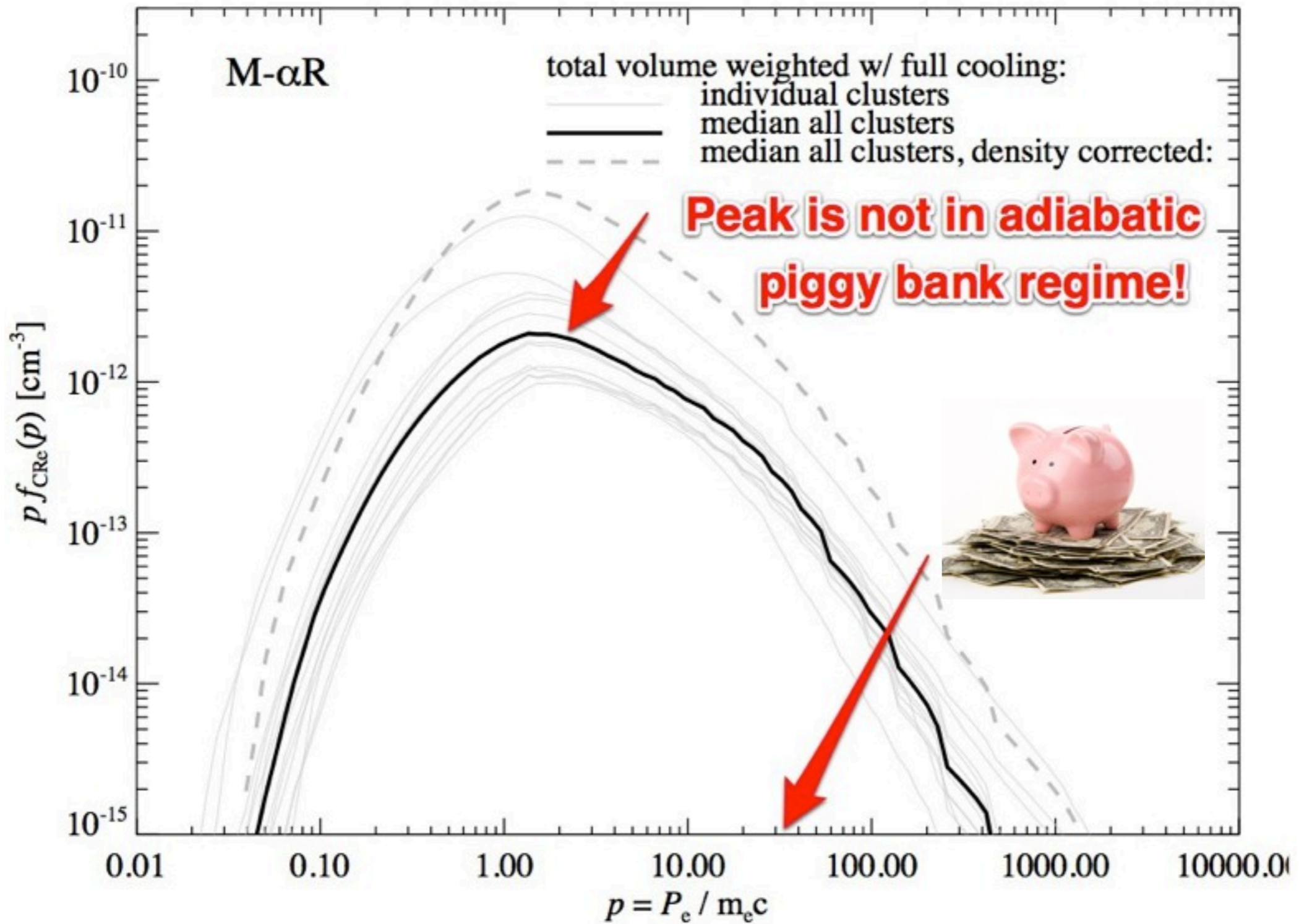
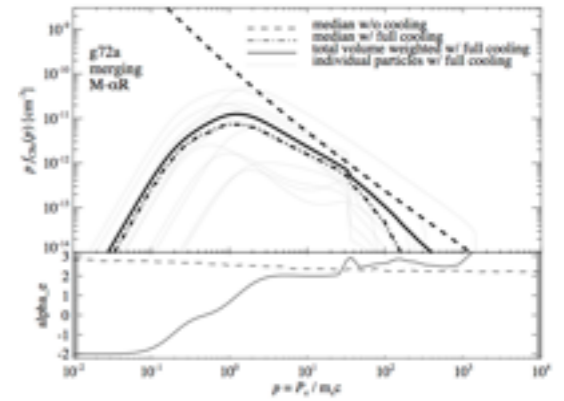
Track Coulomb +  
synchrotron/inverse  
Compton cooling







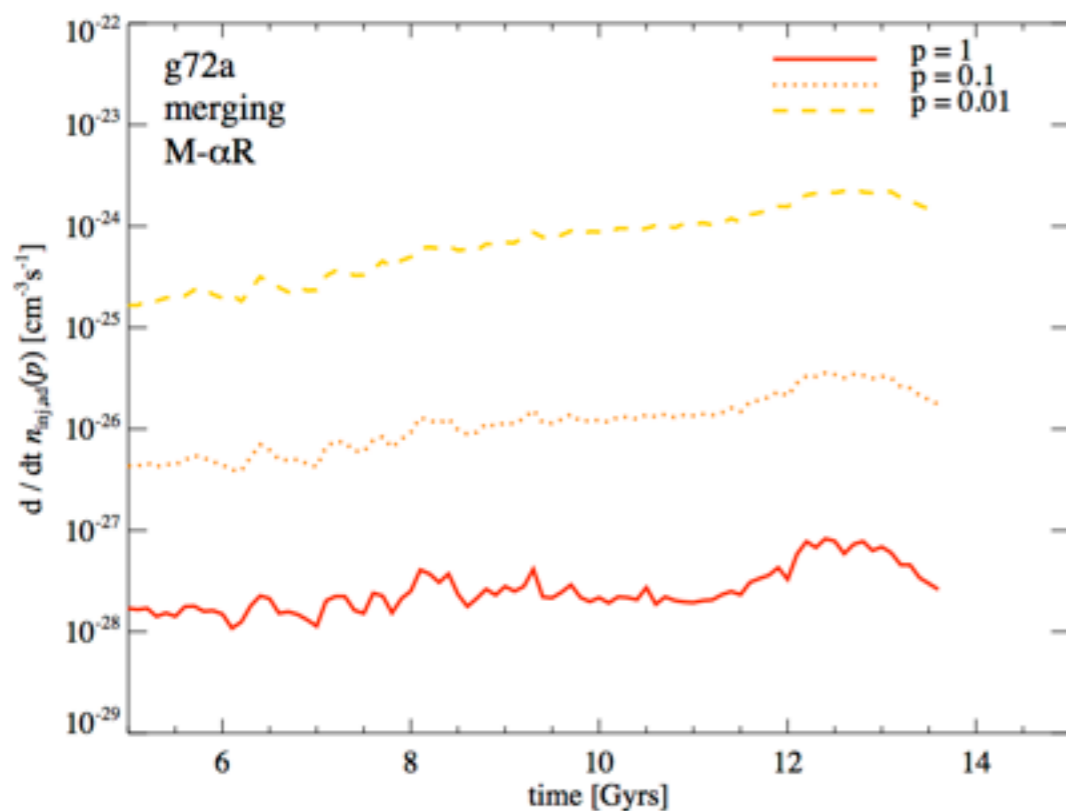
# Result



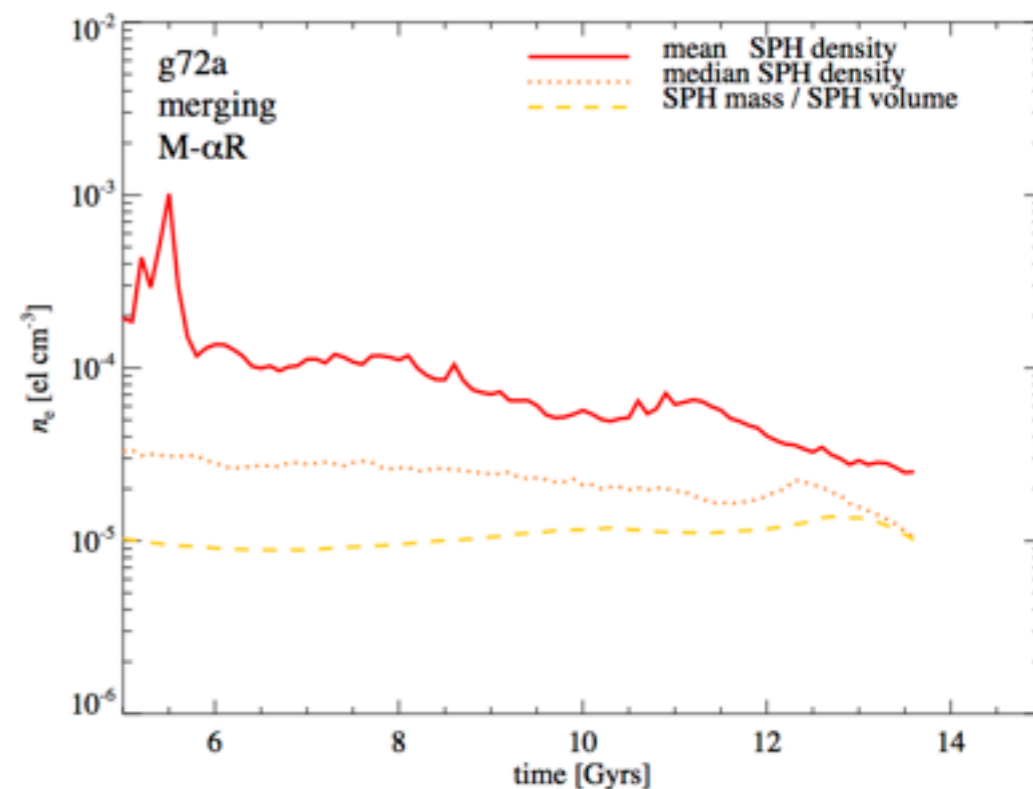


# How to understand it

We only care about **injection** and **Coulomb cooling**



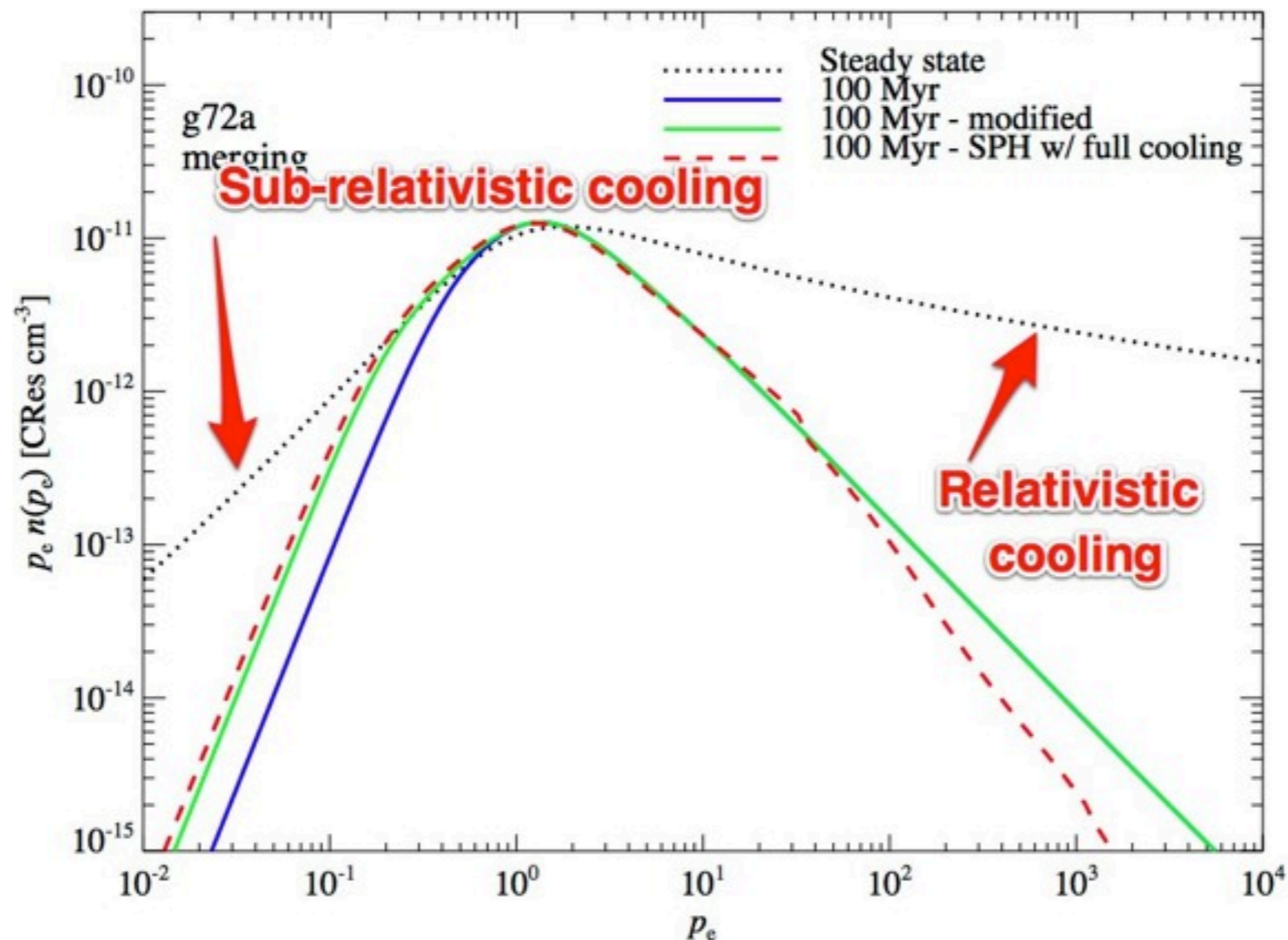
Injection is fairly steady, not impulsive



Density (and hence cooling) is fairly constant

# ...so it's in steady state

this means you can calculate **analytic** models!



The peak looks good...why is the rest a terrible fit?

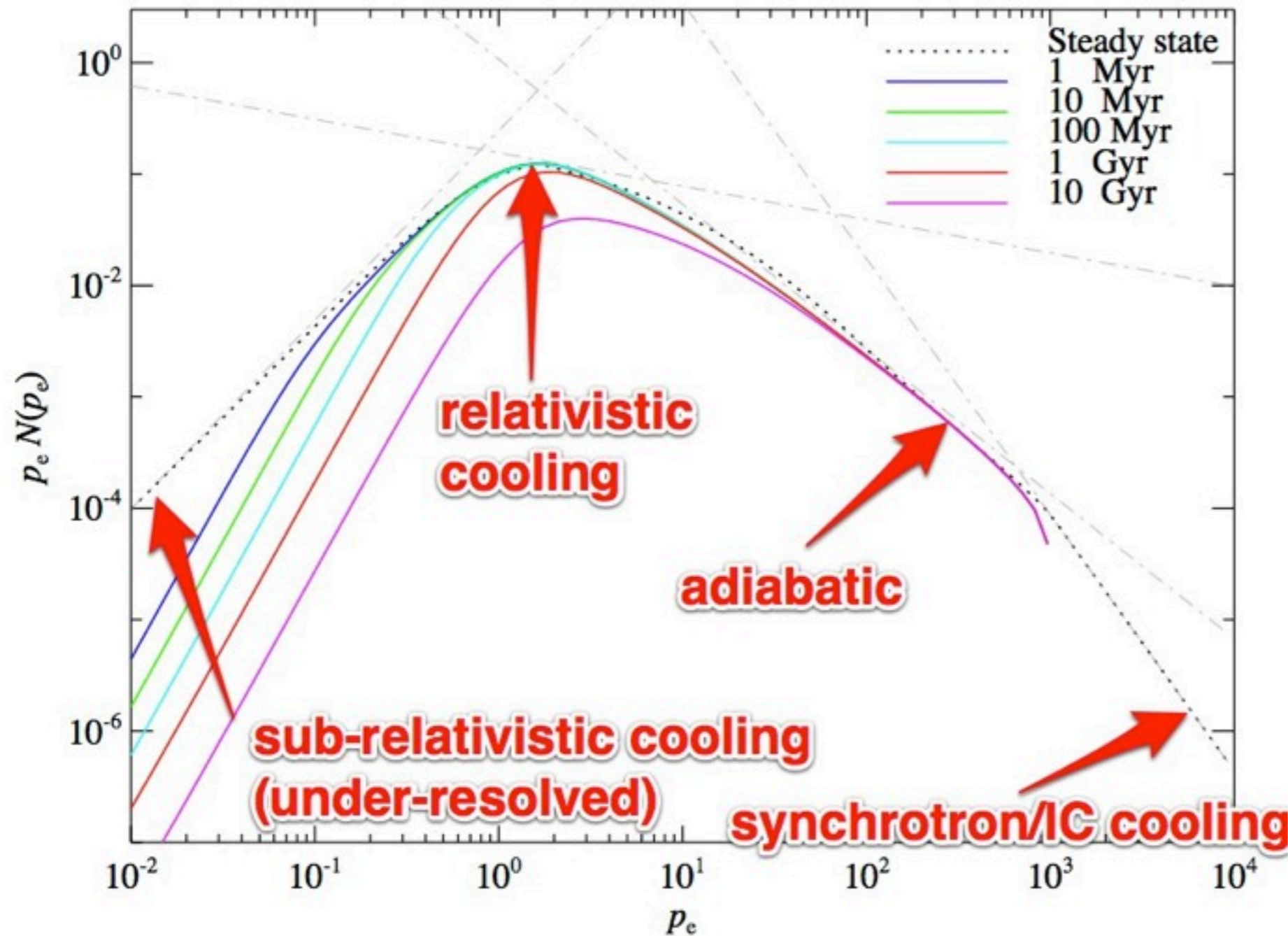


# Time Resolution



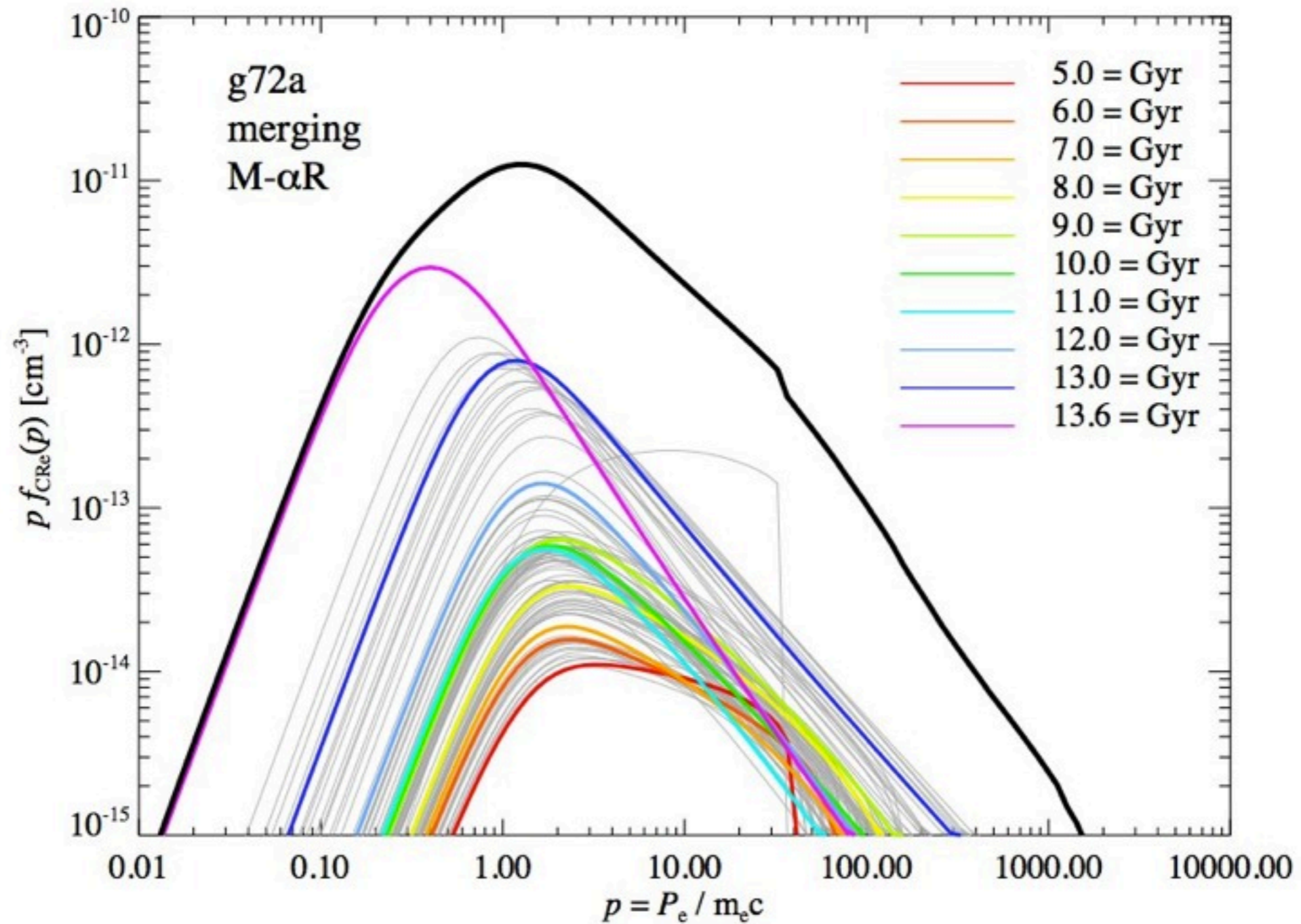
We only resolve  
portion where

$$\Delta t < t_{\text{cool}}$$



But we resolve the crucial peak

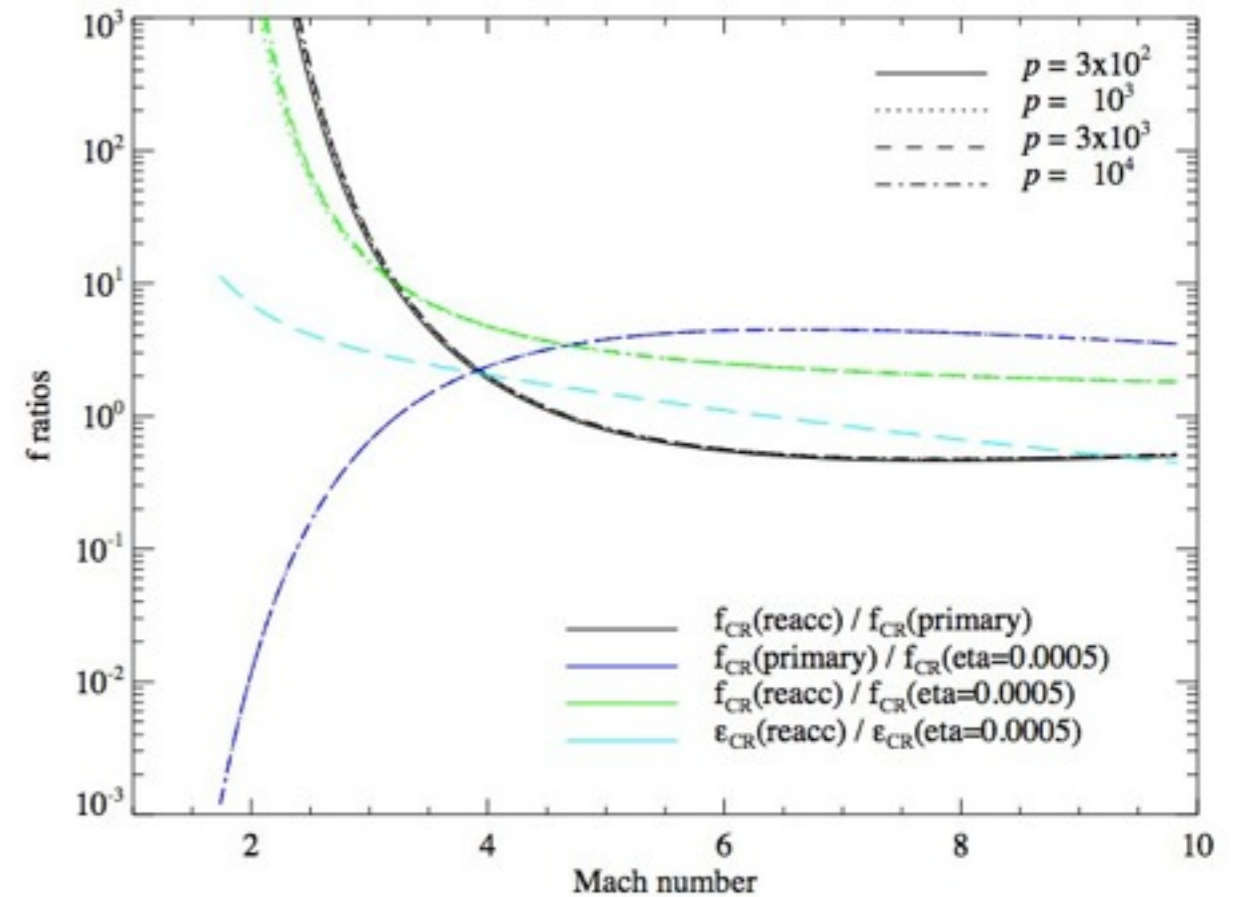
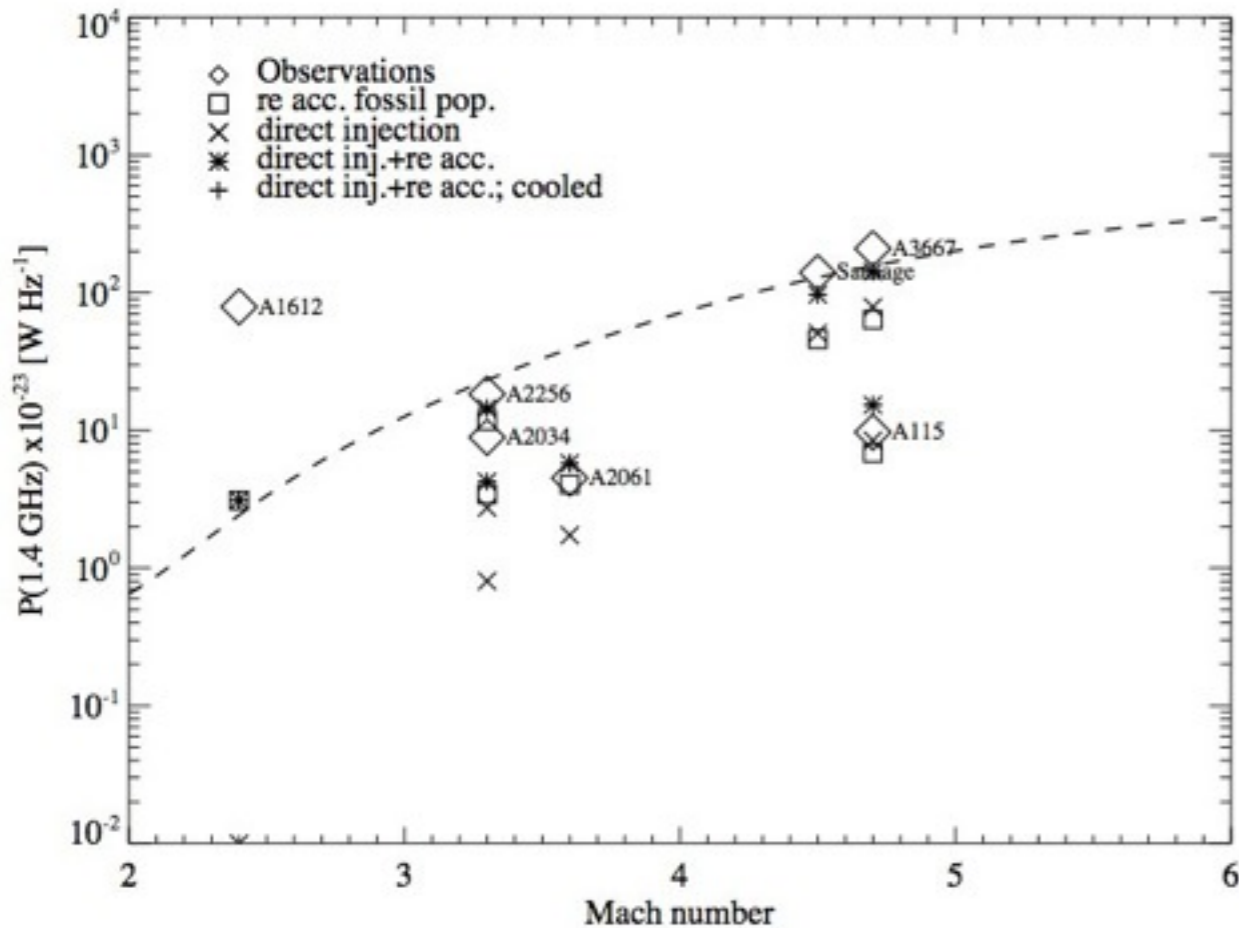
# Stochasticity



Contribution heavily weighted toward late times:  
recent merger history matters...



# The Bottom Line

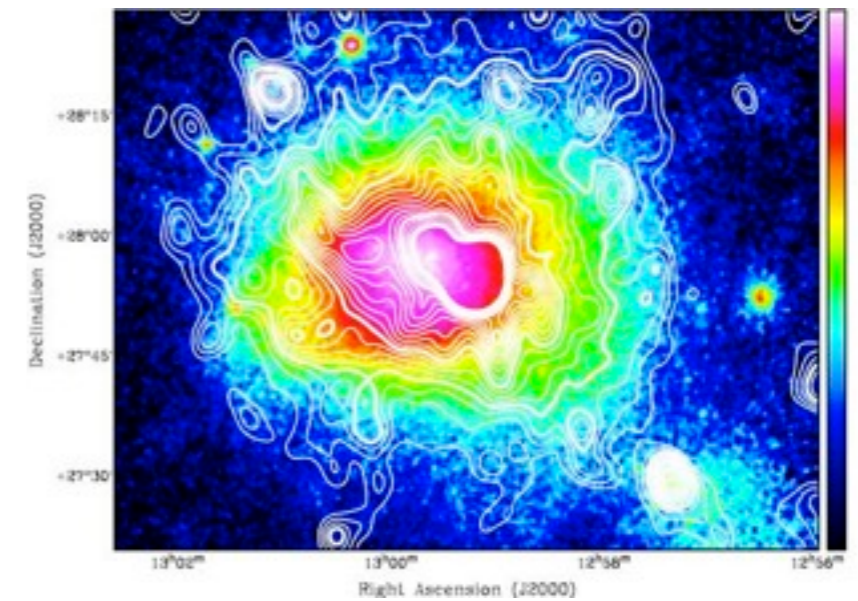


Fossil contribution comparable to direct injection at high  $\mathcal{M}$

**Dominates at low  $\mathcal{M}$**

N.B. Fossils needed even more for radio halos....

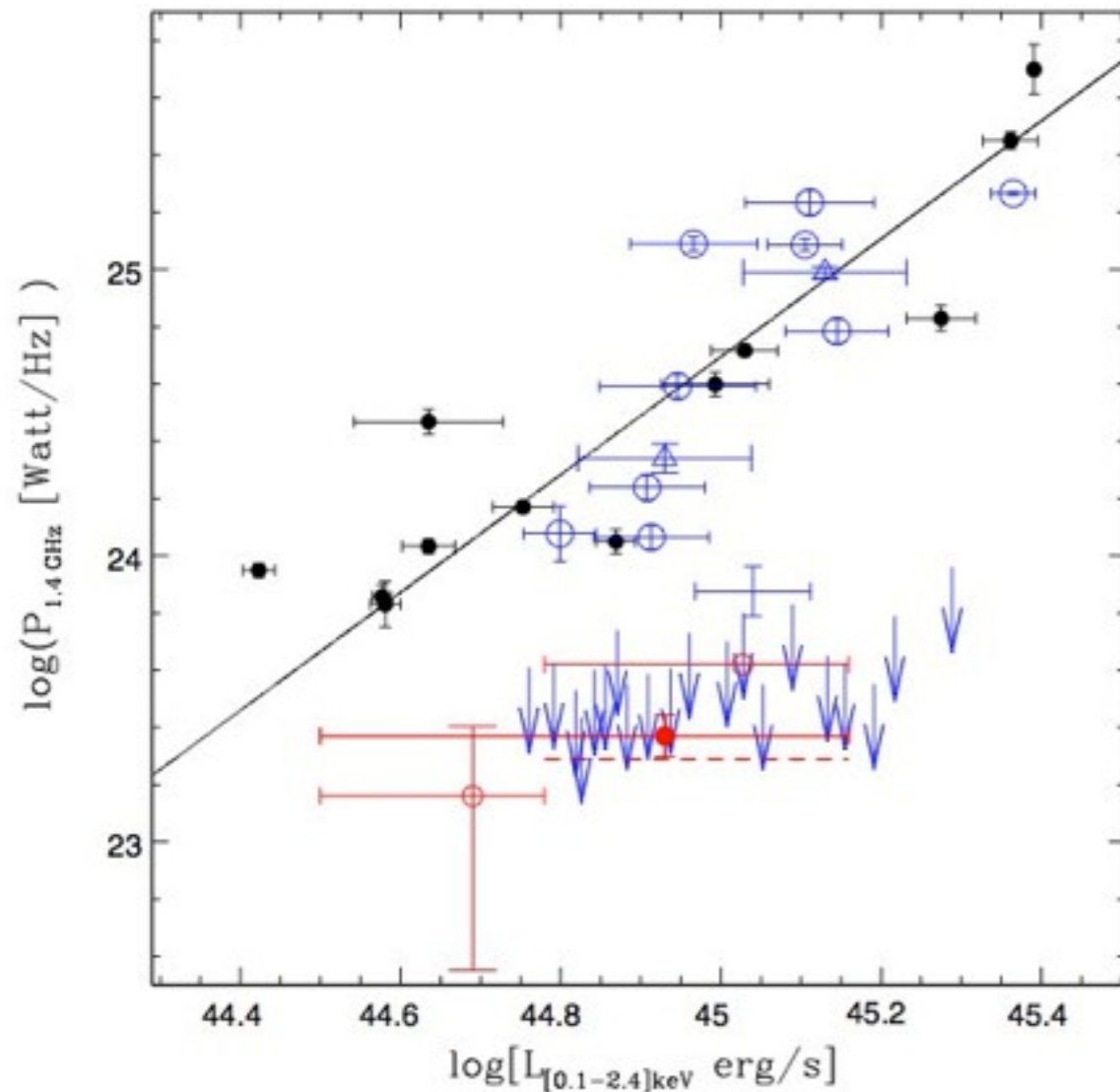




# Turning Radio Halos Off with Cosmic Ray Streaming

Wiener, Oh & Guo, 2012, in prep

# Radio Halos: Two classes



Brown et al 2011  
Brunnetti et al 2009

They are either **ON**  
or **OFF**

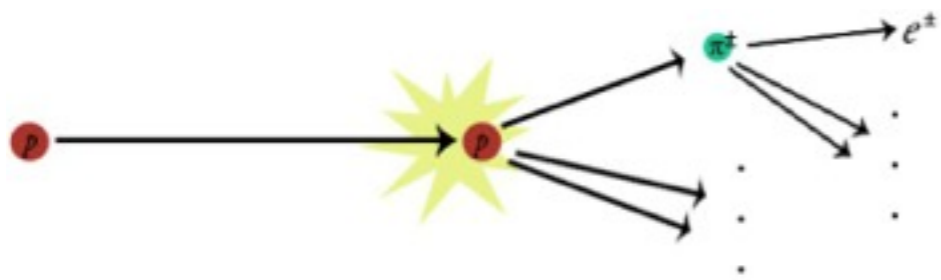
**ON**: recent or ongoing  
merger

**OFF**: Stacked images show  
faint emission (Brown et al 2011)

# Radio Halos: Two models



## Hadronic Model



### The good:

Well-known physics  
Produces nice smooth halos

### The bad:

Can't produce spectral curvature  
Can't produce bimodality

## Turbulent Reacceleration



### The good:

Bimodality is natural  
Produces spectral curvature

### The bad:

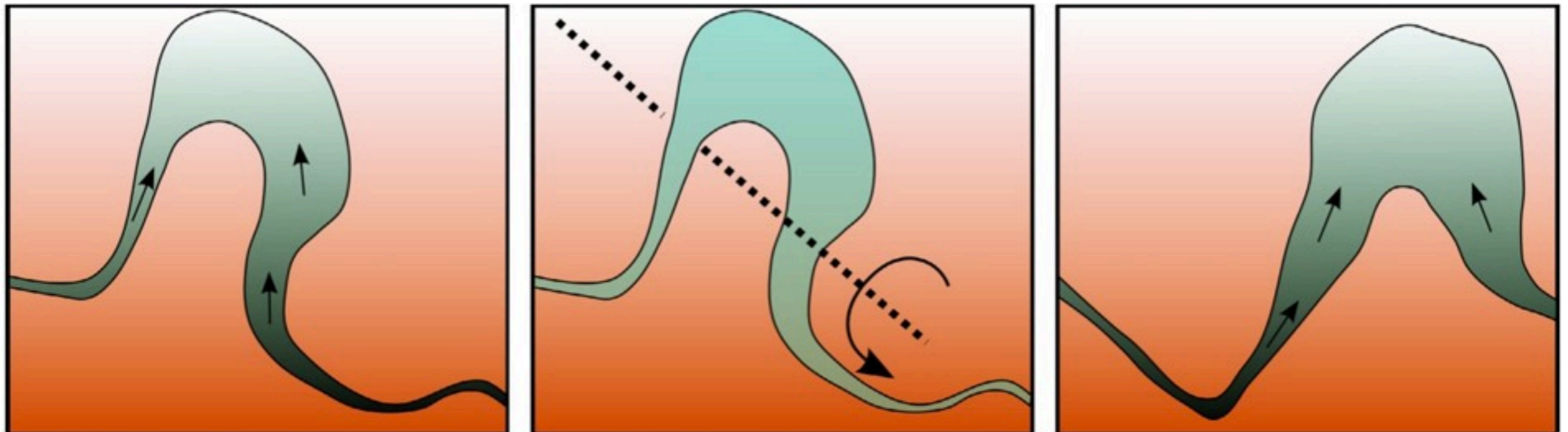
Acceleration efficiency hard to predict  
Need fossil electrons



# Super-Alfvenic streaming



Could it turn off radio halos?



Ensslin, Pfrommer, Miniati, Subramaniam 2011

For this to work, need  $v_D \sim c_S$

Check this!

# Cosmic Ray Speed Limits

In our Galaxy,  $v_D \sim v_A \ll c$



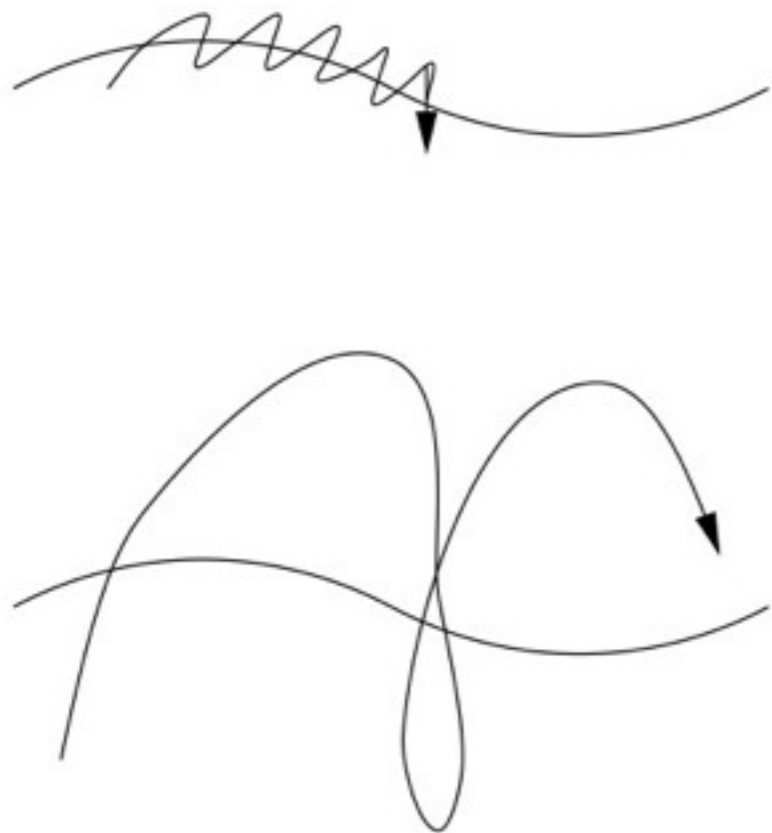
Consistent with observed isotropy  $\delta \sim 10^{-4} \sim \frac{v_D}{c}$

and confinement times of  $\sim 3 \times 10^6$  yr from spallation

Can be explained from CR  
self-confinement



# Resonant wave growth

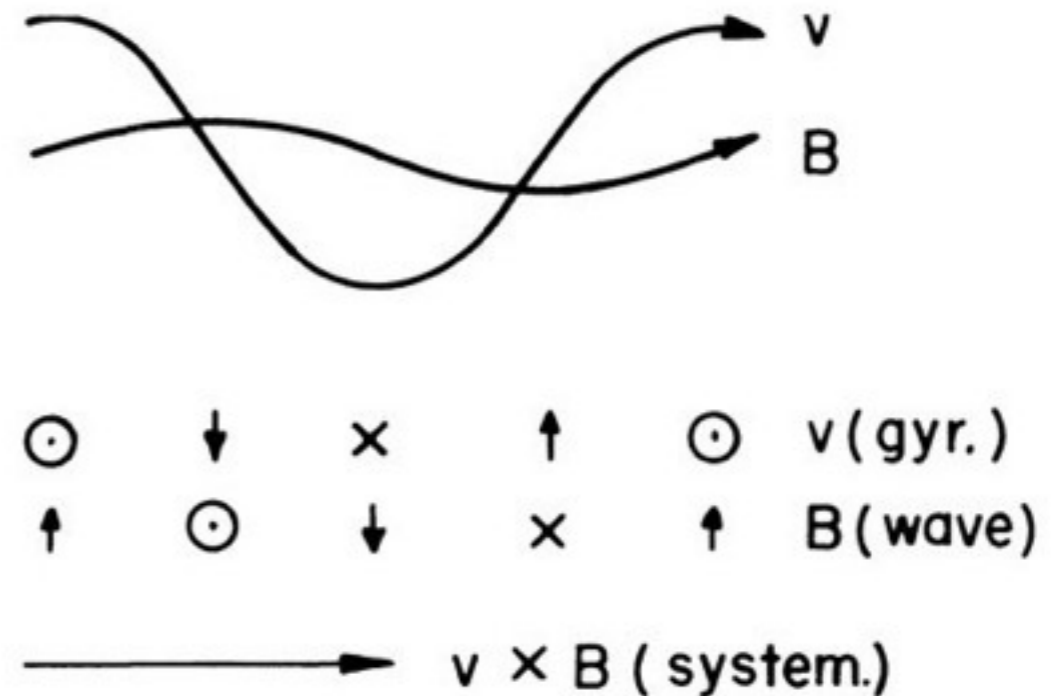


Kulsrud 2006

Resonance condition

$$k_{\parallel} \sim \frac{1}{\mu r_L}$$

Wave growth rate



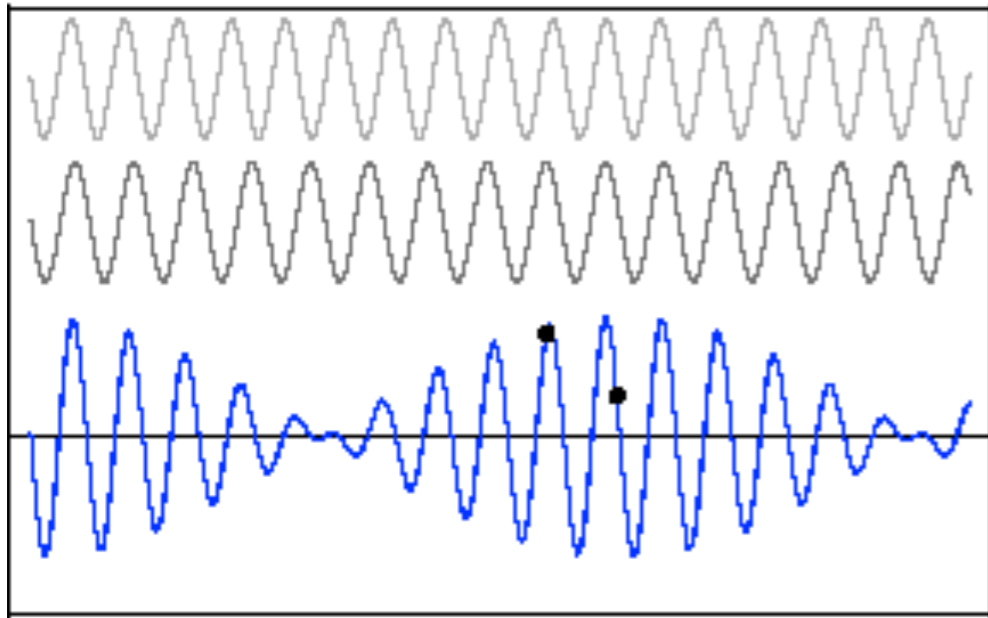
Wentzel 1972

Steady Lorentz force

$$\Gamma_{\text{CR}}(k_{\parallel}) \sim \Omega_0 \frac{n_{\text{CR}}(> \gamma)}{n_i} \left( \frac{v_s}{v_A} - 1 \right)$$



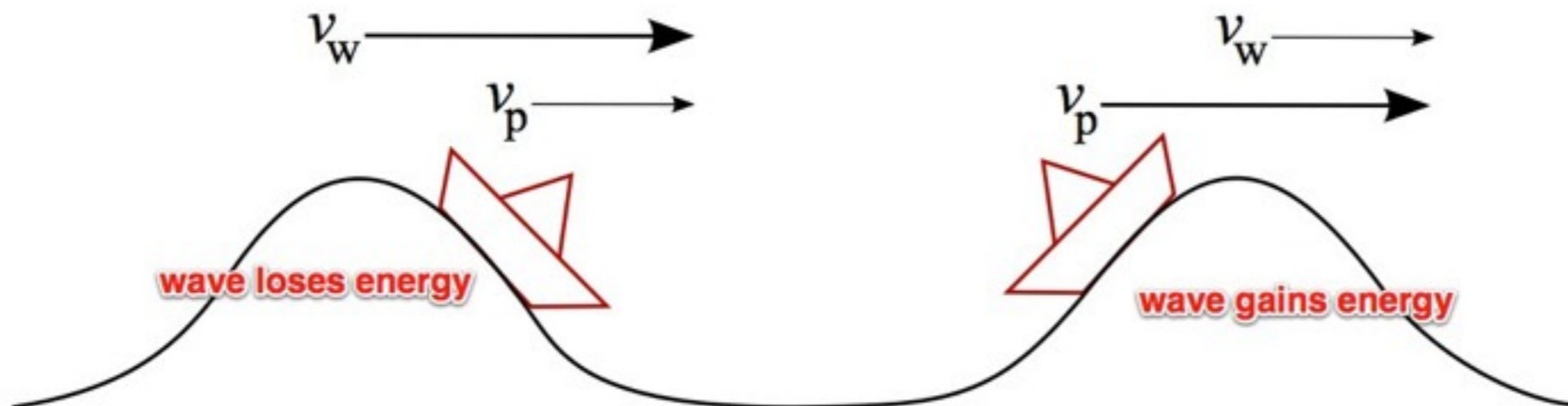
# Non-linear Landau damping



Damping rate:

$$\Gamma \sim \frac{v_i}{v_A} \omega \left( \frac{\delta B}{B} \right)^2$$

Alfvén waves make a beat wave



Thermal particles surf and extract energy

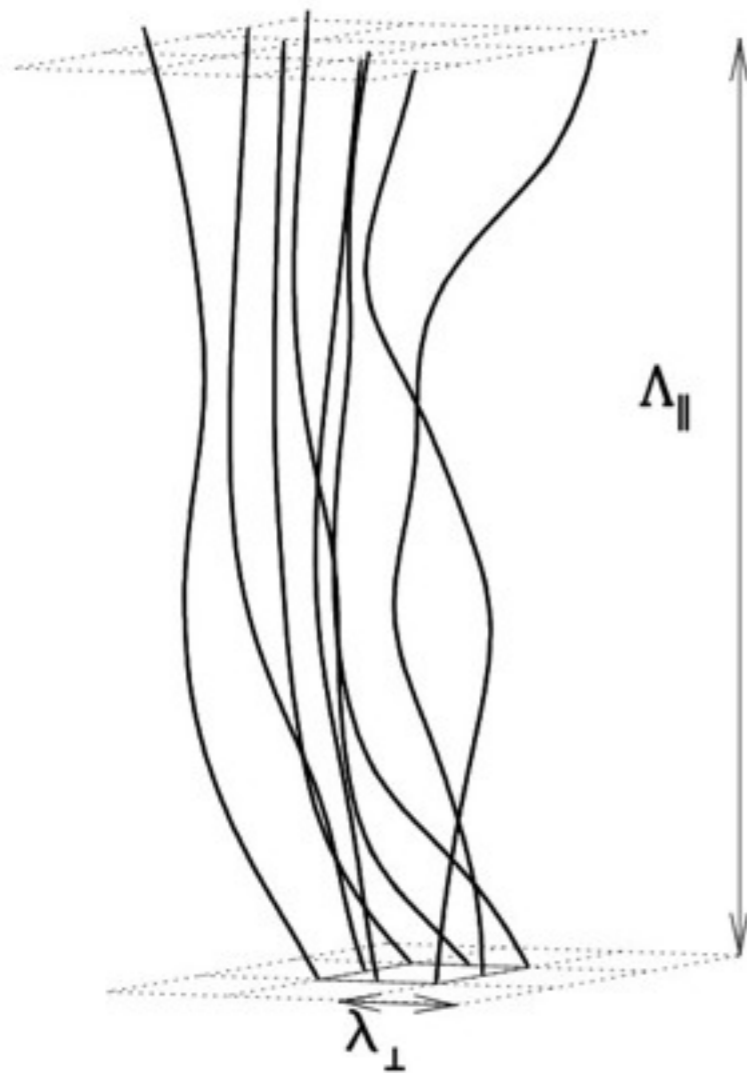
# Turbulent Damping

MHD turbulence is anisotropic

Cascade is primarily transverse to field

Eddy destroyed by cascading to smaller scales

Lithwick & Goldreich 2001



Damping rate

$$\Gamma_{\text{turb,min}} \sim \left( \frac{\epsilon}{r_L v_A} \right)^{1/2}$$

Farmer & Goldreich 2004

where

$$\epsilon \sim \frac{v_A^3}{L}$$

# Streaming speeds

Set  $\Gamma_{\text{grow}} = \Gamma_{\text{damp}}$

Get (for 100 GeV protons):

$$v_D = v_A \left( 1 + 2.4 \frac{T_{4\text{keV}}^{1/4}}{B_{\mu\text{G}} L_{z,100}^{1/2} (n_{-10}^{\text{CR}})^{1/2} (n_{-3}^i)^{3/4}} \gamma_{100}^{(n-3)/2} 100^{(n-3)/2-0.8} \right)$$

Non-linear Landau damping

$$v_D = v_A \left( 1 + 1.6 \frac{B_{\mu\text{G}}^{1/2} n_{i,-3}^{1/2}}{L_{\text{MHD},50}^{1/2} n_{\text{CR},-10}} \gamma_{100}^{n-3.5} 100^{\frac{n-3.5}{1.1}} \right)$$

Turbulent damping

Pretty much still Alfvénic! But note different functional forms... (streaming speed evolves)



# Simulate this!



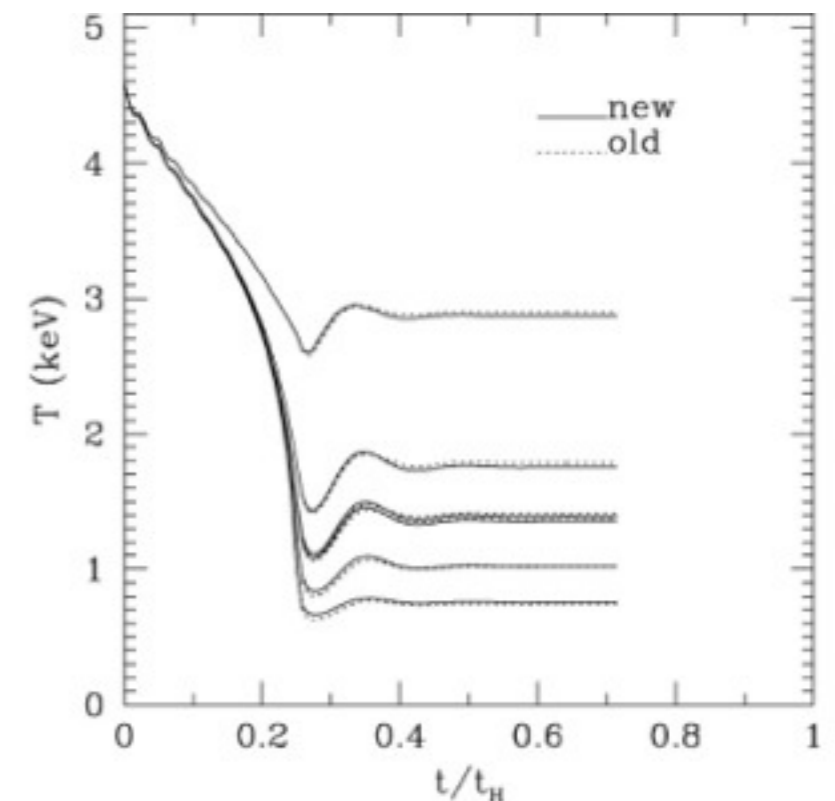
Solve CR equation in spherical symmetry

$$\frac{\partial \bar{f}_p}{\partial t} + (\mathbf{u} + \mathbf{v}_A) \cdot \nabla \bar{f}_p = \nabla \cdot (\kappa \mathbf{nn} \cdot \nabla \bar{f}_p) + \frac{1}{3} p \frac{\partial \bar{f}_p}{\partial p} \nabla \cdot (\mathbf{u} + \mathbf{v}_A) + \bar{Q}$$

where diffusion is

$$\kappa = c^2 \left\langle \frac{1 - \mu^2}{\nu(\mu, \gamma)} \right\rangle \approx L_z (v_D - v_A)$$

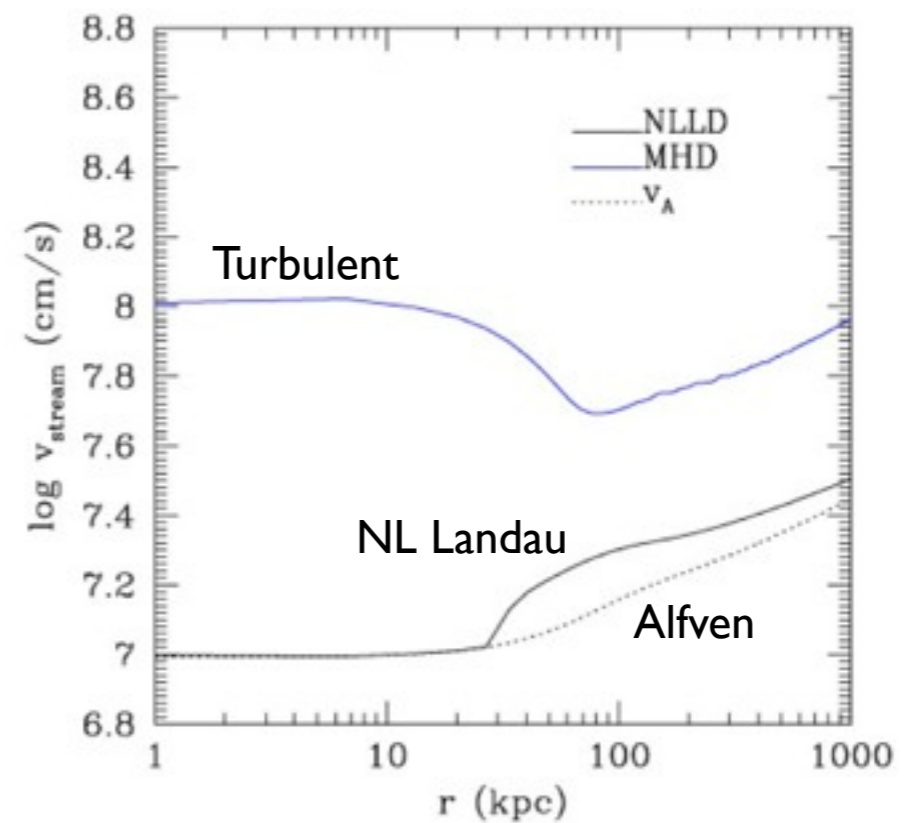
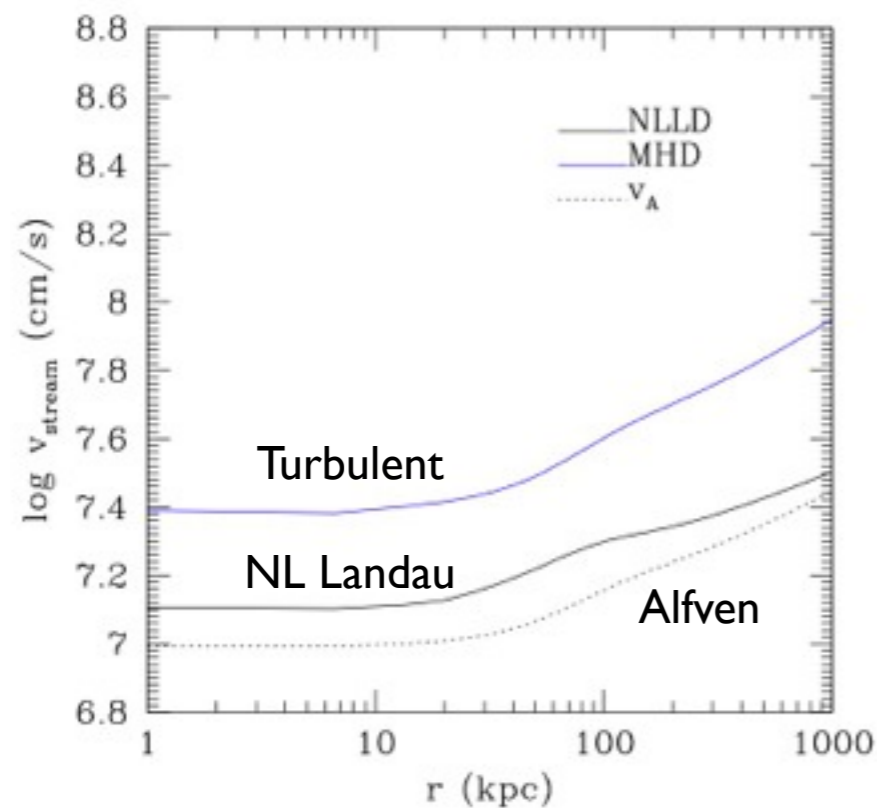
Compare against Guo & Oh  
(2008) when solve for CRs in  
fluid approximation: good  
agreement!



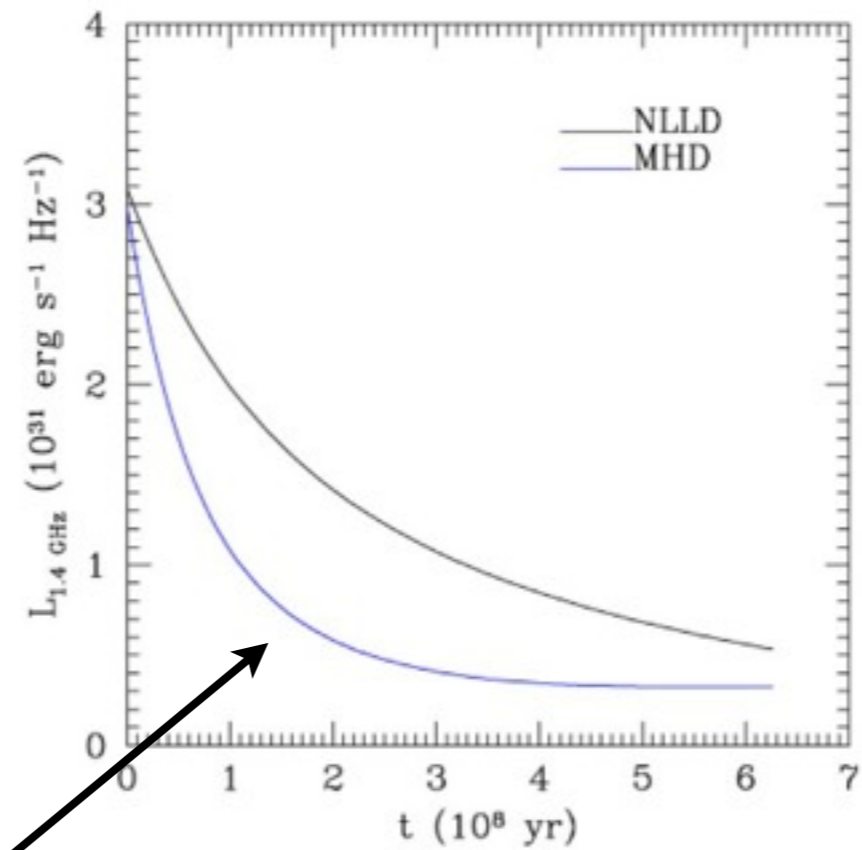
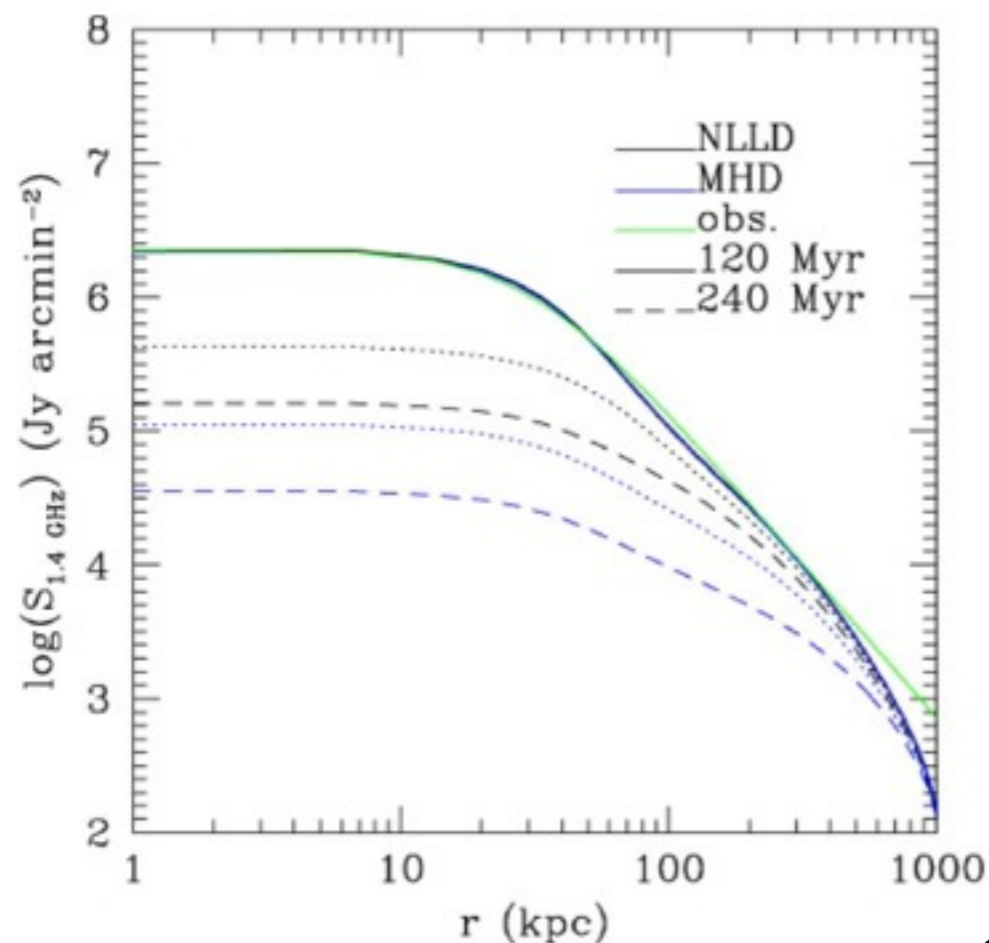


# Results for Perseus

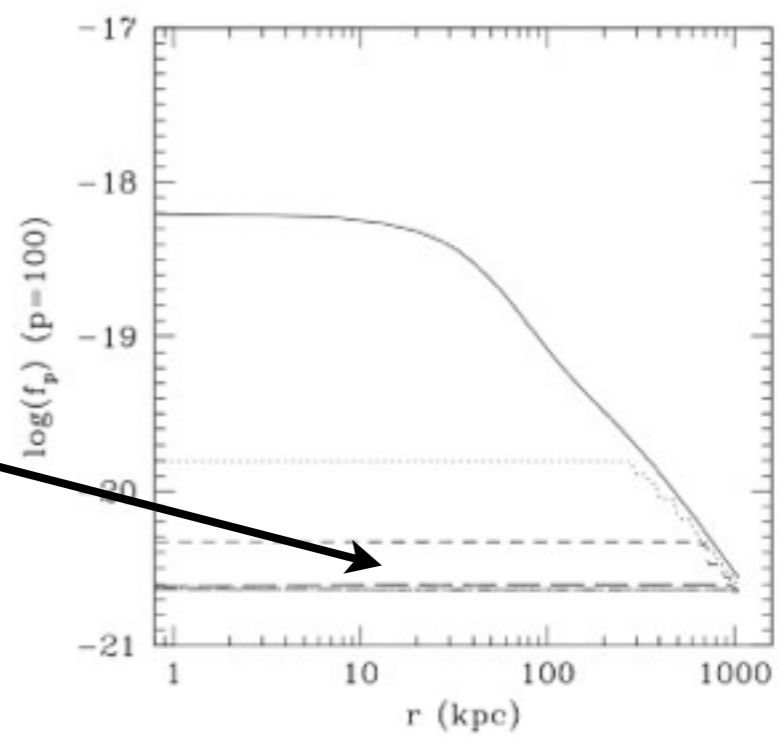
Initialize halo with CR profile needed to reproduce observations (Zandanel et al 2012). Then sit and wait...



Very different evolution in streaming speed for turbulent and non-linear Landau damping



Streaming slows down as  
CR profile flattens

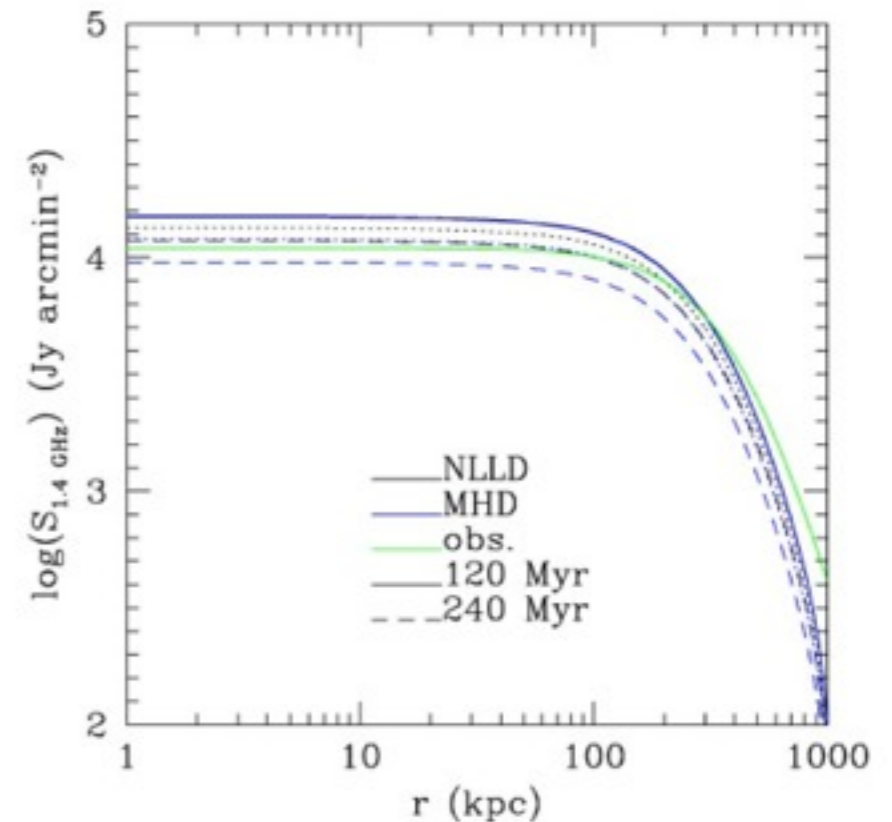
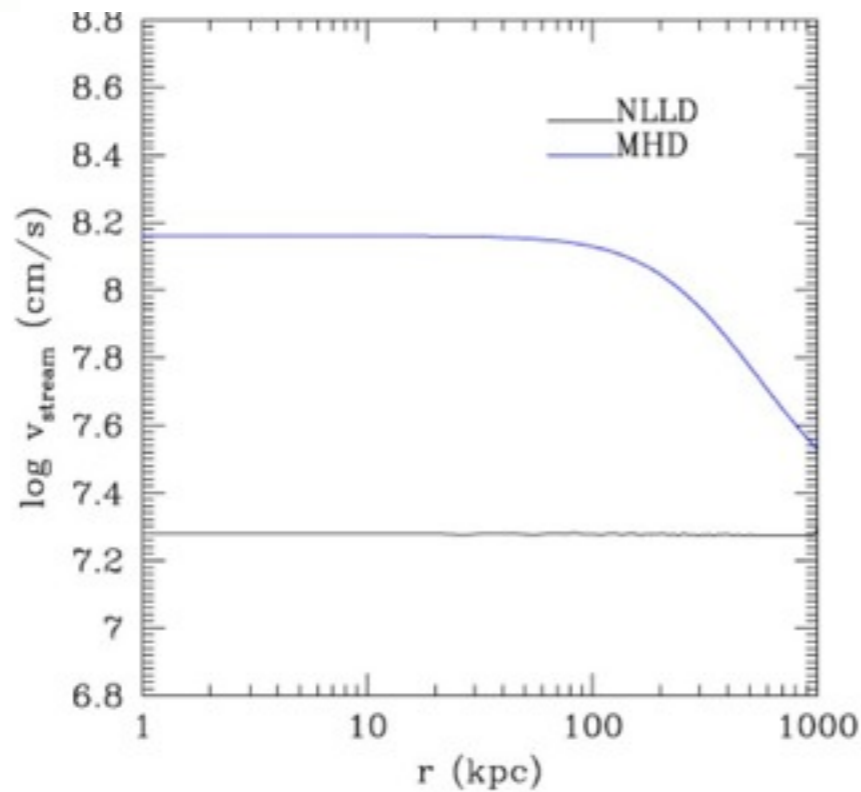


Factor of  $\sim 6$  change in radio luminosity might be possible with turbulent damping



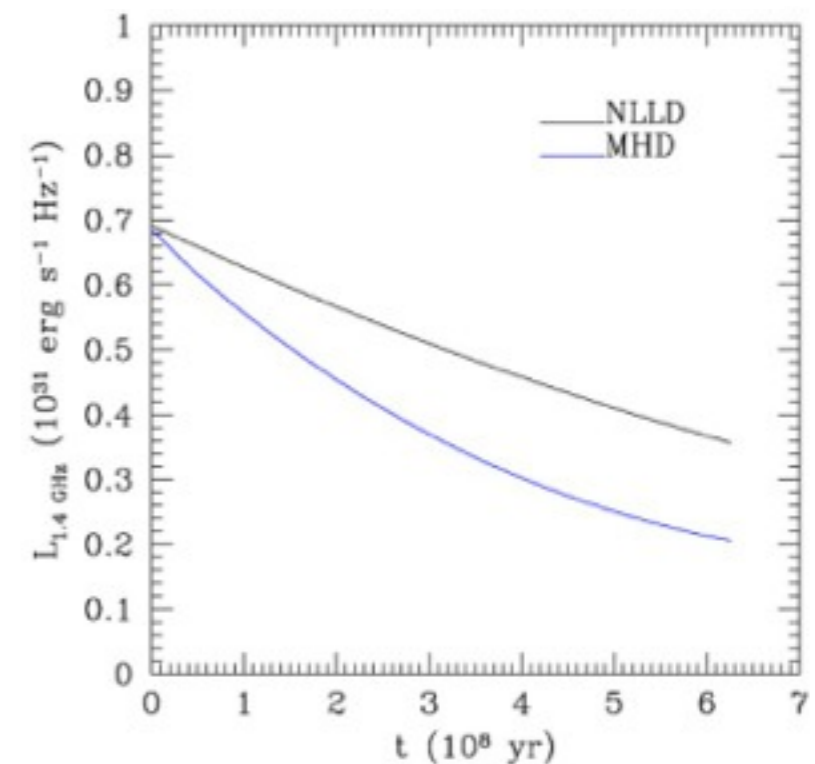


# Results for Coma



Coma has a much larger extent and a flat CR profile as inferred from observations

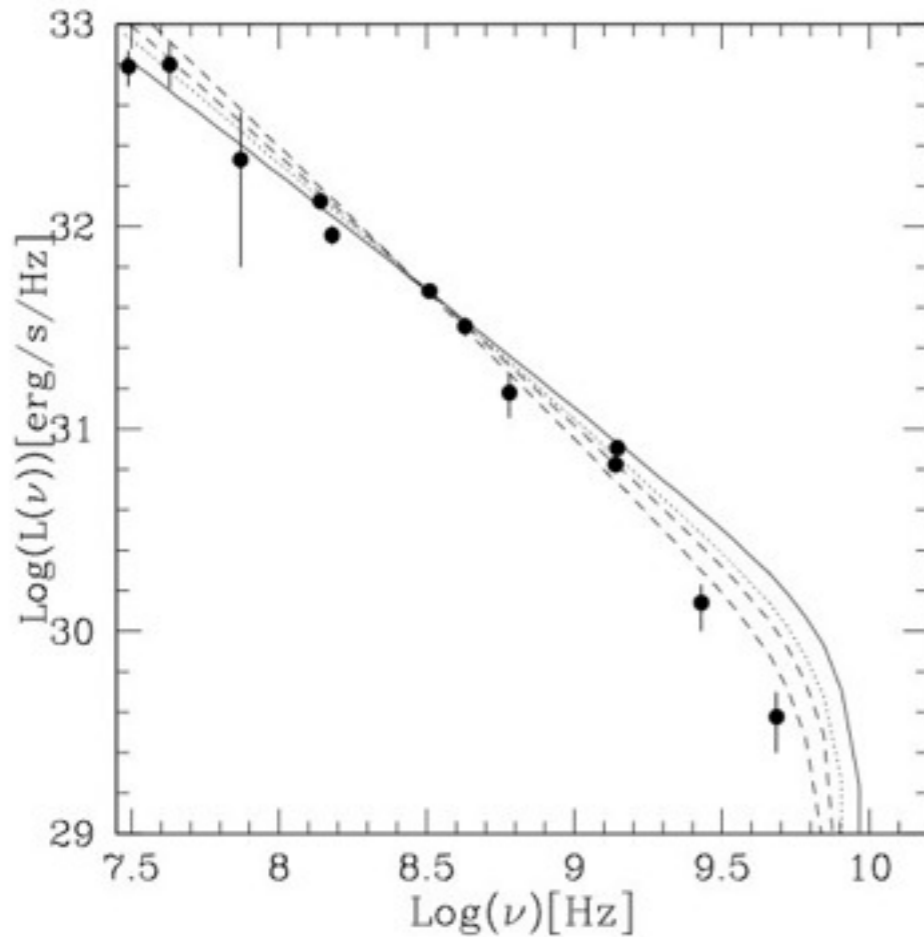
Diffusion time is long, streaming has a small effect.



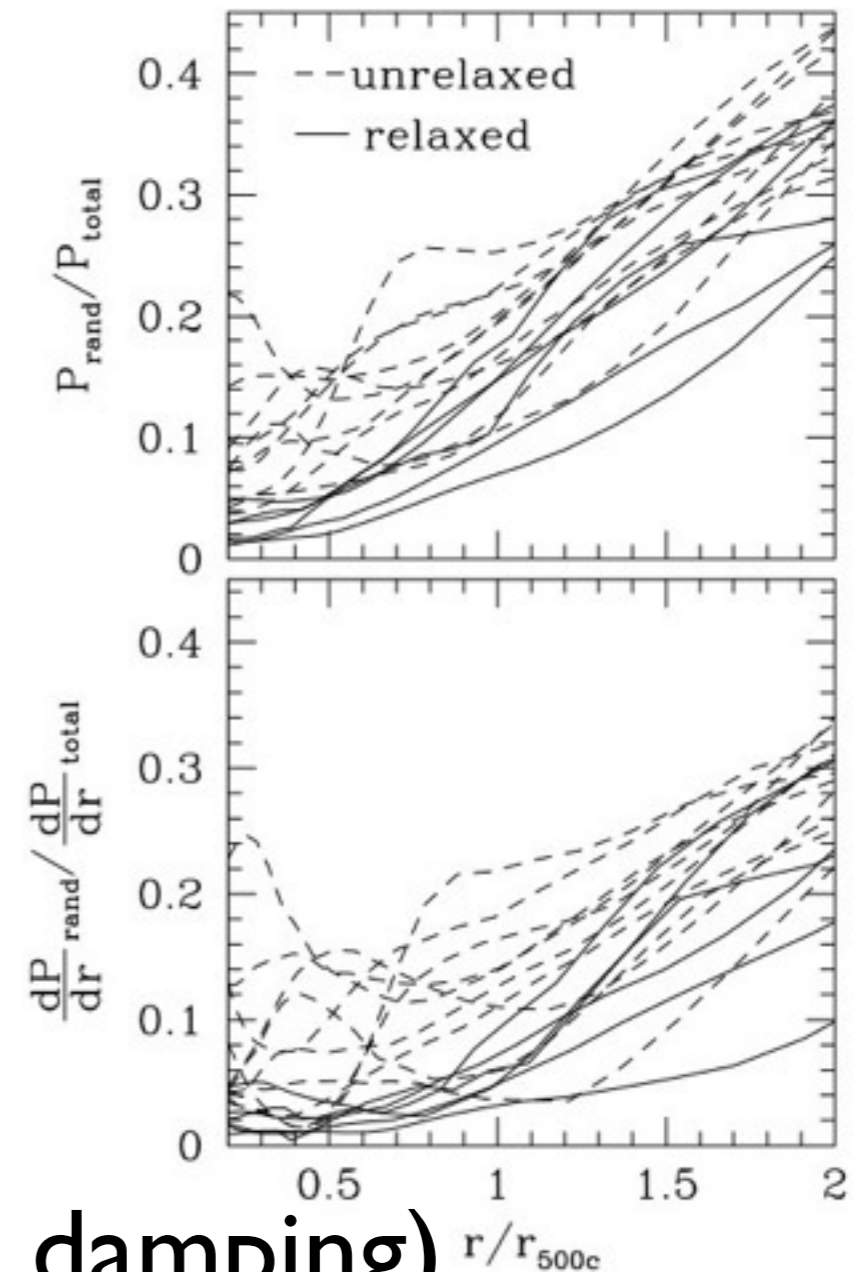


# In Progress

## Spectral curvature



Brunetti et al 2012



Lau et al 2009

Turbulent diffusion

(should use same velocity profile as for damping)

# Take Home

Fossil electrons could allow radio relics to be seen at low Mach numbers

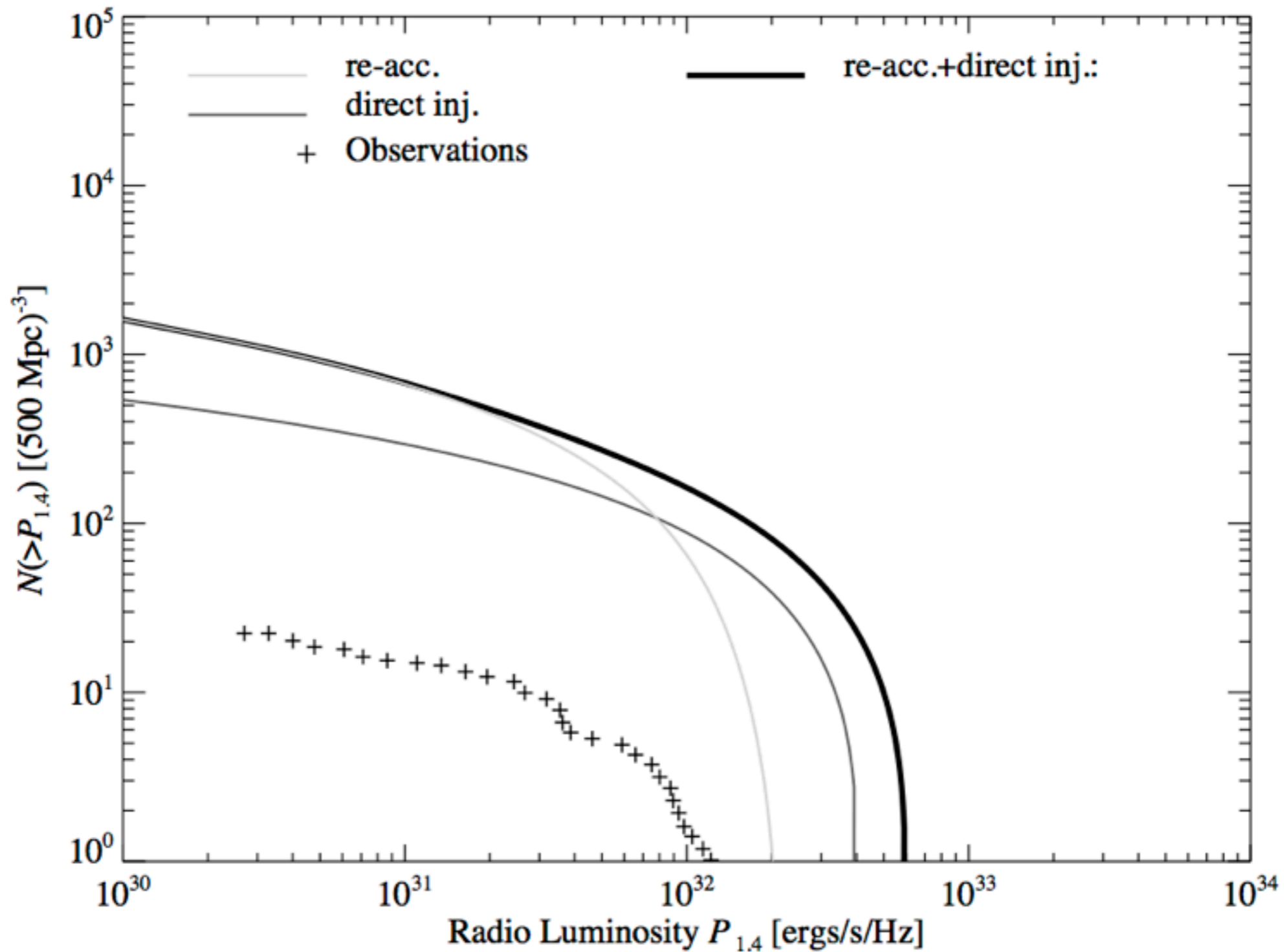
Streaming of cosmic rays might play a role in radio mini-halos, but unlikely in giant radio halos



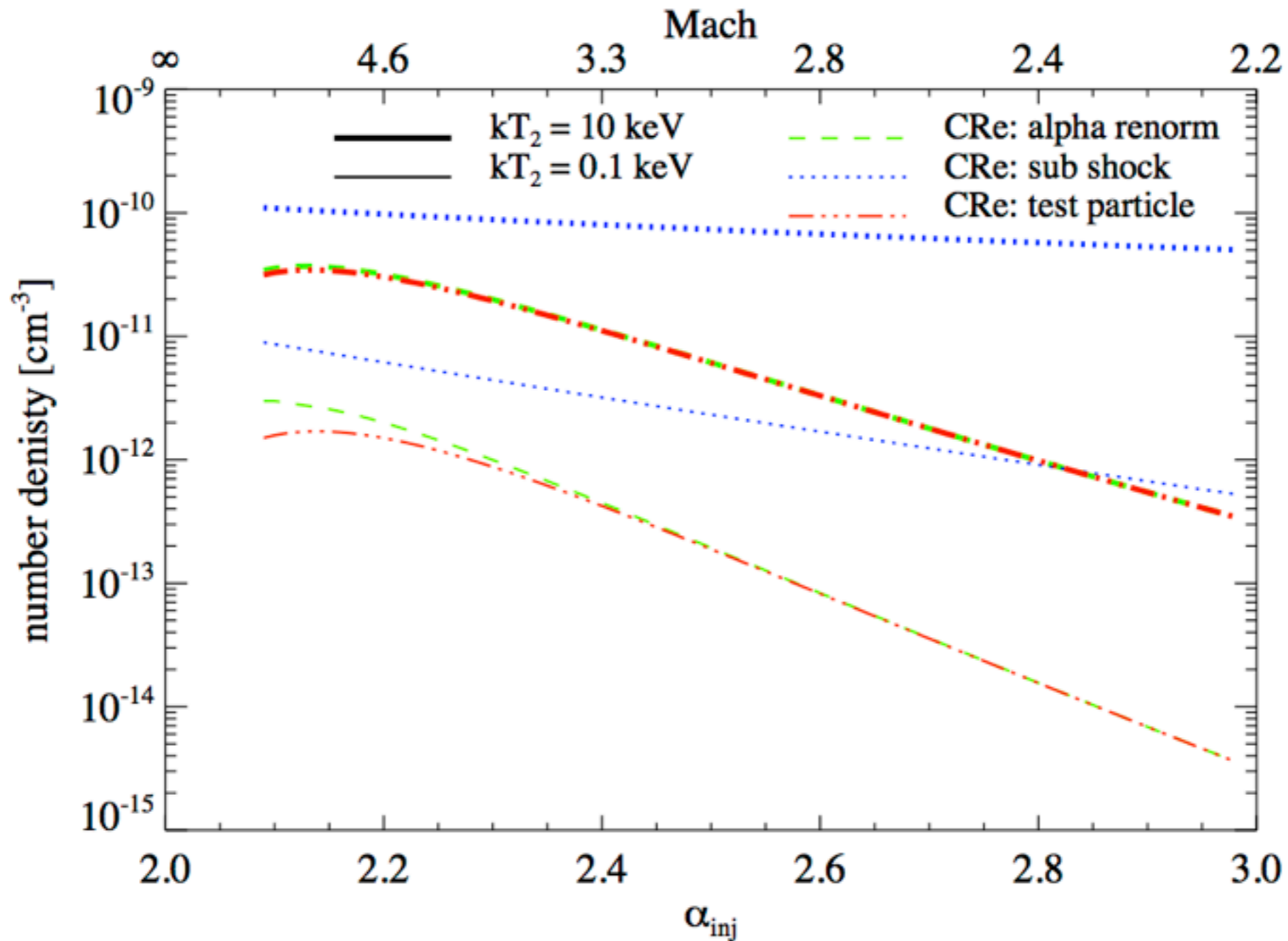




# Relic Counts Predictions



# Model Dependence



Different shock acceleration models can be characterized by the injected number density