

SUPERSYMMETRY PROBLEMS: MONDAY

(1) **Symmetries.**

- (a) Construct an orthonormal basis of 3-dimensional space, using for two of your vectors: $\vec{v} = (1/\sqrt{2}, 1/\sqrt{2}, 0)$ and $\vec{w} = (1/\sqrt{3}, -1/\sqrt{3}, 1/\sqrt{3})$. Note: you will need to find your third vector. Check that it's an orthonormal basis.

Now describe a change of basis taking (x, y, z) in the standard basis to (x', y', z') in your new basis. Verify that the matrix that describes this satisfies the condition $A^T A = I$.

- (b) Consider the set of symmetries on (x, y, z) , for each t :

$$\begin{aligned}x' &= x + 2ty + 3t^2z \\y' &= y + 3tz \\z' &= z\end{aligned}$$

Show that these symmetries form a group under composition (that is, doing one such symmetry using t followed by another such symmetry using s).

- (c) For the symmetry in (b), find the infinitesimal description.
 (d) For $O(n)$ prove that the condition that $A^T A = I$ turns into the condition that $A^T = -A$ for the Lie algebra.
 (e) For $O(3)$ use the generators of the Lie algebra

$$J_x = \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & -1 \\ 0 & 1 & 0 \end{bmatrix}, \quad J_y = \begin{bmatrix} 0 & 0 & 1 \\ 0 & 0 & 0 \\ -1 & 0 & 0 \end{bmatrix}, \quad J_z = \begin{bmatrix} 0 & -1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}$$

Show that

$$[J_x, J_y] = J_z$$

$$[J_y, J_z] = J_x$$

$$[J_z, J_x] = J_y$$

- (f) Derive Equation (21) in Jim Gates' notes: gates-Lect1.pdf.
 (g) Suppose we have the complex matrices

$$s_x = \frac{1}{2} \begin{bmatrix} 0 & i \\ i & 0 \end{bmatrix}, \quad s_y = \frac{1}{2} \begin{bmatrix} 0 & -1 \\ 1 & 0 \end{bmatrix}, \quad s_z = \frac{1}{2} \begin{bmatrix} i & 0 \\ 0 & -i \end{bmatrix}$$

and show that these satisfy the same commutation relations as for $J_x, J_y,$ and J_z in $O(3)$. Characterize the matrices spanned by $s_x, s_y,$ and s_z . Note that this means that at least infinitesimally, $O(3)$ can be a symmetry of 2-dimensional complex ordered pairs (w, z) using these relations. These (w, z) are called *spinors*.

(2) **Electromagnetism.**

Let $\vec{E} = (E^1, E^2, E^3)$ be the electric field and $\vec{B} = (B^1, B^2, B^3)$ be the magnetic field. We will work in Lorentz-Heaviside units, where the speed of light c , the magnetic constant μ_0 , and the electric constant ϵ_0 are all 1.

Let ϕ be the *electromagnetic scalar potential* and let $\vec{A} = (A^1, A^2, A^3)$ be the *electromagnetic vector potential*. Recall that \vec{E} and \vec{B} are given by

$$\begin{aligned}\vec{E} &= -\frac{\partial \vec{A}}{\partial t} - \nabla \phi \\ \vec{B} &= \nabla \times \vec{A},\end{aligned}$$

where ∇ is the operator $\left(\frac{\partial}{\partial x}, \frac{\partial}{\partial y}, \frac{\partial}{\partial z}\right)$.

This problem will use the Minkowski metric $\eta_{\mu\nu} = \text{diag}(-1, 1, 1, 1)$ for raising and lowering of indices.

- (a) Define the *current 4-vector* J^μ by (ρ, J^1, J^2, J^3) , and define the *electromagnetic four-potential* by $A_\mu = (\phi, \vec{A})$. Expand $A_\mu J^\mu$. What is the *current one-form* J_μ ?
- (b) Define the *field strength tensor* $F_{\mu\nu}$ by

$$F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu.$$

Write $F_{\mu\nu}$ as a 4×4 matrix in terms of the components of \vec{E} and \vec{B} .

- (c) Two of Maxwell's equations for the electromagnetic field are:

$$\begin{aligned}\nabla \cdot \vec{B} &= 0 \\ \nabla \times \vec{E} &= -\frac{\partial \vec{B}}{\partial t}\end{aligned}$$

Show that these equations are equivalent to

$$\partial_\xi F_{\mu\nu} + \partial_\mu F_{\nu\xi} + \partial_\nu F_{\xi\mu} = 0.$$

- (d) Write $F_{\mu\nu} F^{\mu\nu}$ in terms of E and B , where $E = \|\vec{E}\|$ and $B = \|\vec{B}\|$.
- (e) Let $\vec{J} = (J^1, J^2, J^3)$. The rest of Maxwell's equations for the electromagnetic field are:

$$\begin{aligned}\nabla \cdot \vec{E} &= \rho \\ \frac{\partial \vec{E}}{\partial t} - \nabla \times \vec{B} &= -\vec{J}\end{aligned}$$

Show that these are equivalent to:

$$\partial_\nu F^{\mu\nu} = J^\mu.$$

(3) Topologies and Chromotologies.

In *Adinkras for mathematicians*, Y. Zhang writes:

An n -dimensional *adinkra topology*, or topology for short, is a finite connected simple graph A such that A is bipartite and n -regular (every vertex has exactly n incident edges). We call the two sets in the bipartition of $V(A)$ bosons and fermions, though the actual choice is mostly arbitrary and we do not consider it part of the data. A *chromotology* of dimension n is a topology A such that the following holds.

- The elements of $E(A)$ are colored by n colors ... such that every vertex is incident to exactly one edge of each color.

- For any distinct i and j , the edges in $E(A)$ with colors i and j form a disjoint union of 4-cycles.
- (a) Draw two different chromotopologies.
 - (b) Construct one of your chromotopologies as a directed graph in Sage.
You may wish to consult the documentation at:
<http://www.sagemath.org/doc/prep/Quickstarts/Graphs-and-Discrete.html>
or
<http://www.sagemath.org/doc/reference/graphs/sage/graphs/digraph.html>
 - (c) Can you give an example of an adinkra topology that does not admit a chromotopology?
 - (d) When defining chromotopologies in her senior thesis at Bard College, S. Naples wrote:
Finally, every pair of edge colors $\{a, b\}$ incident to a single vertex is part of a 4-cycle of alternating edge colors, such that the edge colors alternate $abab$.
Does Naples' condition follow from Zhang's definition of a chromotopology? Why or why not?
 - (e) How many chromotopologies can you construct on at most 6 vertices, and with $n \leq 4$?
 - (f) What do you think it should mean for two chromotopologies to be isomorphic?